

Mathematics for Semiconductor Heterostructures

Modeling, Analysis, and Numerics

**Weierstrass Institute for
Applied Analysis and Stochastics
September 24 – 28, 2012**

www.wias-berlin.de/workshops/msh2012



Weierstrass Institute for
Applied Analysis and Stochastics

www.wias-berlin.de



Deutsche
Forschungsgemeinschaft

www.dfg.de

Cover figure: Computed 3D temperature distribution in an organic semiconductor device containing a thin layer sequence of n-doped / intrinsic / n-doped C₆₀ between crossbar metal electrodes (*Organic Electronics*, Vol. 13, No. 11, 2012).

Weierstrass Institute for Applied Analysis and Stochastics (WIAS)
Leibniz Institute in Forschungsverbund Berlin e.V.
Mohrenstraße 39 · 10117 Berlin · Germany

Fax: +49 30 20372 303

E-Mail: contact@wias-berlin.de

WWW: <http://www.wias-berlin.de>

**International Workshop on
Mathematics for Semiconductor
Heterostructures
— Modeling, Analysis, and Numerics —**

Organizers

Klaus Gärtner	Weierstrass Institute
Annegret Glitzky	Weierstrass Institute
Hans-Christoph Kaiser	Weierstrass Institute
Francis Nier	Université de Rennes 1

Local Organizers

Olga Kuphal
Marion Lawrenz
Matthias Liero

Support

German Research Foundation (DFG) & Weierstrass Institute, Berlin (WIAS)

Scope

The aim of this workshop is to bring together applied mathematicians and scientists from semiconductor physics and to give them the opportunity to discuss analytical and numerical aspects of relevant semiconductor models including the design of efficient numerical algorithms.

Topics:

- Modeling of micro-, nano-, and optoelectronic devices
- Classical drift-diffusion and energy models
- Semi-classical and quantum transport models
- Asymptotic models
- Organic semiconductors
- Spintronics
- Quantum dots, wires, wells, and waveguides
- Simulation of real-world devices

Monday, 24.09.2012, 08:00 - 17:10

08:00 - 09:00	REGISTRATION
09:00 - 09:30	OPENING
09:30	Virtual Experiments with High Power Semiconductor Devices along and beyond the Rim of the Safe Operating Area Gerhard Wachutka
10:20 - 10:50	COFFEE BREAK
10:50	Using gradient structures for modeling semiconductors Alexander Mielke
11:30	Asymptotic analysis of a 1D KP system Francis Nier
12:10 - 14:00	LUNCH BREAK
14:00	The Oskar3 project Klaus Gärtner
14:40	Semi-classical modeling of quantum dot lasers with microscopic treatment of Coulomb scattering Thomas Koprucki
15:20 - 15:50	COFFEE BREAK
15:50	Electron transport in strongly confined nanostructures: Modeling and simulation Paola Pietra
16:30	Asymptotic models for the transport of an electron gas strongly confined in a slab Fanny Delebecque

Tuesday, 25.09.2012, 09:00 - 17:10

09:00	WKB-schemes for Schrödinger-type equations Anton Arnold
09:40	On the moment problem for quantum hydrodynamics Olivier Pinaud
10:20 - 10:50	COFFEE BREAK and POSTER PRESENTATION Carolin Kreisbeck and Leo Bonato
10:50	A plane-wave based, generalised formulation of continuum elasticity and multiband kp-models Oliver Marquardt
11:30	Optical Recombination in Semiconductor Nanostructures Eoin O'Reilly
12:10 - 14:00	LUNCH BREAK
14:00	The Gateable Feature of DEPFET Sensor Arrays Rainer Richter
14:40	Physics and modeling of strain effects in SiGe heterojunction bipolar transistors Holger Rucker
15:20 - 15:50	COFFEE BREAK
15:50	WIAS-TeSCA simulations in photovoltaics for a point contact concept of heterojunction thin film solar cells Reiner Nürnberg
16:30	Application of WIAS-TeSCA to the analysis of the power limits of semiconductor lasers Hans Wenzel

Wednesday, 26.09.2012, 09:00 - 21:00

09:00	Artificial interface conditions in quantum transport models Andrea Mantile
09:40	Modeling boundary conditions for solving stationary Schrödinger equations Matthias Ehrhardt
10:20 - 10:50	COFFEE BREAK and POSTER PRESENTATION Carolin Kreisbeck and Leo Bonato
10:50	Discretization of Time Dependent Quantum Systems: The Evolution Operators Joseph Jerome
11:30	Linear and non-linear eigensolver-based techniques for quantum modeling and simulations: Applications to carbon-based materials Eric Polizzi
12:10 - 14:00	LUNCH BREAK
14:00	Modeling Spintronic Effects in Silicon Siegfried Selberherr
14:50 - 15:15	COFFEE BREAK
15:45 - 18:00	GUIDED TOUR AT NEUES MUSEUM
18:30 - 21:00	DINNER AT "CUM LAUDE"

Program

Thursday, 27.09.2012, 09:00 - 17:10

09:00	Development and experimental validation of a predictive electrical model of organic light-emitting diodes Reinder Coehoorn
09:40	Problems with the Gaussian density of states in organics: Definite solution, ambiguous, and hardly any solution Gernot Paasch
10:20 - 10:50	COFFEE BREAK and POSTER PRESENTATION Carolin Kreisbeck and Leo Bonato
10:50	Drift-Diffusion-Recombination Models for Excitonic Organic Photovoltaic Devices Klemens Fellner
11:30	Multiscale Modeling of Heterojunction Organic Photovoltaic Devices Matteo Porro
12:10 - 14:00	CONFERENCE PHOTO and LUNCH BREAK
14:00	An electronic model for solar cells taking into account active interfaces Annegret Glitzky
14:40	On a model of magnetization switching driven by a spin current: Modeling and numerical simulations Clément Jourdana
15:20 - 15:50	COFFEE BREAK
15:50	Bilinear control of Schrödinger equations Karine Beauchard
16:30	Some approximate controllability results for bilinear systems Sylvain Ervedoza

Friday, 28.09.2012, 09:00 - 16:00

09:00	An application of the Landauer-Büttiker formula to photon emitting and absorbing systems Hagen Neidhardt
09:40	A new decomposition for the Schrödinger equation in the semiclassical asymptotics and an associated AP numerical scheme Florian Méhats
10:20 - 10:50	COFFEE BREAK
10:50	Plane waveguides with corners in the small angle limit Nicolas Raymond
11:30	Quantum transport in semiconductor nano-heterostructures Paul Racec
12:10 - 14:00	LUNCH BREAK
14:00	Utilizing Wavelets for the Solution of High-Dimensional Semiconductor Transport Equation Vincent Peikert
14:40	Numerical Simulation of Hetero Nanowire Tunnel Field Effect Transistors Andreas Schenk
15:20 - 15:30	CLOSING
15:30 - 16:00	COFFEE BREAK

WKB-schemes for Schrödinger-type equations

Anton Arnold

Vienna University of Technology, Institute of Analysis & Scientific Computing
Wiedner Hauptstr. 8, 1040 Vienna, Austria

e-mail: Anton.Arnold@tuwien.ac.at

We are concerned with the numerical integration of ODEs of the form $\epsilon^2 \psi_{xx} + a(x)\psi = 0$ for given $a(x) \geq \alpha > 0$ in the highly oscillatory regime $0 < \epsilon \ll 1$ (appearing as a stationary Schrödinger equation or 1D Helmholtz equation, e.g.). In two steps we derive an accurate finite difference scheme that does not need to resolve each oscillation: 1) With a WKB-ansatz the dominant oscillations are “transformed out”, yielding a much smoother ODE. 2) For the resulting oscillatory integrals we devise an asymptotic expansion both in ϵ and \hbar . In contrast to existing strategies, the presented method has (even for a large spatial step size h) the same weak limit (in the classical limit $\epsilon \rightarrow 0$) as the continuous solution. Moreover, it has an error bound of the order $O(\epsilon^3 h^2)$. We shall give applications to k-p-Schrödinger systems and to the simulation of semiconductor-nanostructures.

REFERENCES

- [1] A. Arnold, N. Ben Abdallah and C. Negulescu: WKB-based schemes for the oscillatory 1D Schrödinger equation in the semi-classical limit, *SIAM J. Numer. Anal.* 49, No. 4 (2011) 1436-1460.
- [2] J. Geier, A. Arnold: WKB-based schemes for two-band Schrödinger equations in the highly oscillatory regime, to appear in *Nanosystems: Physics, Chemistry, Mathematics* 2, No. 3 (2011).

Bilinear control of Schrödinger equations

Karine Beauchard

CNRS, CMLS École Polytechnique, Centre de Mathématiques Laurent Schwartz,
Route de Saclay, 91120 Palaiseau, France

e-mail: `Karine.Beauchard@math.polytechnique.fr`

We consider a quantum particle in a potential and an electric field, represented by a Schrödinger equation. It is a control system in which the state is the wave function of the particle and the control is the electric field: one wants to steer the quantum particle from a given initial state, to a desired target state, by applying a suitable control. In the dipolar moment approximation, the control acts bilinearly on the state in the equation. The goal of this talk is to propose an overview of results and technics concerning this problem. We address several questions: exact and approximate controllability, feedback stabilization, influence of the dimension, existence of a positive minimal time required for controllability. First, we recall methods introduced in finite dimension (for ODEs), relying on iterated Lie brackets and discuss their adaptation in infinite dimension (for PDEs). Concerning exact controllability, we propose negative and positive results, depending on the functional frame, the dimension of the space and the assumption on the dipolar moment. Then, we discuss their generalizations. Concerning feedback stabilization, first, we present powerful methods in finite dimension (for ODEs), relying on Lyapunov functions and LaSalle invariance principle; then, we propose adaptations in infinite dimension (for PDEs).

Deep-Level Transient Spectroscopy for QD Localization Energy

Leo Bonato, Tobias Nowozin, Andreas Marent, and Dieter Bimberg

Technische Universität Berlin Fakultät II - Institut für Festkörperphysik
Hardenbergstr. 36, 10623 Berlin, Germany

e-mail: Leo.Bonato@tu-berlin.de

Nowadays, the main limitation of so-called “nonvolatile memory devices” is the long time needed to access data. For example, Flash memory can store data for longer than 10 years, but writing data requires about $1\mu\text{s}$. On the other hand, “volatile memory devices” like DRAM have data access times shorter than 15ns, but can only retain information for microseconds [1]. A device able to combine the virtues of volatile and nonvolatile memory would be the “Holy Grail of memory devices”.

The QD-Flash project aims to realize such a device by using quantum dots (QDs) as storage element, embedding them in a MODFET (Modulation-Doping Field Effect Transistor) in order to read and write data. The viability of this concept has been already demonstrated on a prototype showing write times of 80ns, erase times of 350ns and 1.6s storage time at room temperature [2]. The current goal is increasing the storage time, which depends on the localization energy of the QDs [3]. The localization energy depends on the material combination used to manufacture the heterostructure and can be determined by Deep Level Transient Spectroscopy, a technique originally developed for traps in semiconductors [4]. However, a simulation method able to take into account all the complex effects generated in the heterostructures and to provide a reliable estimate of the localization energy is still missing.

The poster presents the results obtained for the localization energy and storage time of GaSb QDs in a GaAs matrix, GaSb QDs in GaAs with one or two AlGaAs barriers and InGaAs QDs in GaP, relating them to the expected values obtained from the theoretical dependence of the storage time on the localization energy. The poster also describes the advantages that could be obtained from mathematical simulations of the devices and the related challenges to be overcome.

REFERENCES

- [1] J. E. Brewer and M. Gill, *Nonvolatile Memory Technologies with Emphasis on Flash*, John Wiley & Sons, Inc. Hoboken, New Jersey, 2008.
- [2] A. Marent et al., “The QD-Flash: a quantum dot-based memory device”, *Semic. Sci. Technol.* 25 000000 (2010)
- [3] A. Marent et al., “106 years extrapolated hole storage time in GaSb/AIAs quantum dots”, *App. Phys. Lett.* 91, 242109 (2007)
- [4] D. V. Lang, “Deep-level transient spectroscopy: a new method to characterize traps in semiconductors”, *J. Appl. Phys.* 45 3023

Development and experimental validation of a predictive electrical model of organic light-emitting diodes

Reinder Coehoorn and Peter A. Bobbert

Philips Research Laboratories Eindhoven and Eindhoven University of Technology,
High Tech Campus 4, Box WAG-12, 5656 AE Eindhoven, The Netherlands

e-mail: Reinder.Coehoorn@philips.com

OLEDs are large-area light sources with potentially unique applications, e.g. as very thin light-emitting foils with every desired color. The light-emitting layer is only approximately 100 nm thick, and consists of a large number of organic semiconductor sub-layers, each with a different function. The physics which describes the opto-electronic processes is radically different from that in inorganic LEDs, due to the presence of strong disorder. Three-dimensional Monte-Carlo simulations show that the current density is therefore filamentary. In the first part of the talk, it is shown that such 3D calculations can now be used to obtain the 3D current density and recombination profile, without the need to develop for each layer and each charge carrier one-dimensional expressions for the mobility function [1]. Although such calculations are computer-time-intensive, they provide information about the 3D spatial distribution of the charge-carrier density and the exciton density, and at a 3D-level about the charge-exciton and exciton-exciton quenching processes which determine the efficiency and lifetime at high current densities. Whereas this is still work-in-progress, it is shown how such studies can already now be used to evaluate the sensitivity of the OLED properties to materials parameters and OLED stack design changes. In the second part of the talk, it is discussed how the results of 3D calculations can be accurately “translated” to one-dimensional mobility models. The following situations concerning transport and recombination in white multilayer OLEDs are considered:

- A novel view on the effects of Gaussian disorder on the mobility [2].
- Accuracy and limitations of parameter extraction methods [3].
- Transport of in an electric field at interfaces and in the presence of trap states [4].
- Modelling of time-dependent transport, as in the case of impedance or dark-injection experiments, including a realistic description of the process of charge-carrier relaxation [5].
- Modelling of transport in systems with very thin layers, and with emissive dyes.
- Modelling of charge-carrier recombination in disordered systems [6].
- Experimental validation using emission profile reconstruction [7].

In an outlook, several open issues, where a 3D-to-1D translation has yet to be realized, are discussed.

REFERENCES

- [1] M. Mesta et al. (unpublished).
- [2] J. Cottaar, L. J. A. Koster, R. Coehoorn, and P. A. Bobbert, *Phys. Rev. Lett.* 107, 136607 (2011).
- [3] R.J. de Vries et al. (unpublished).
- [4] J. Cottaar, R. Coehoorn, and P. A. Bobbert, *Phys. Rev. B* 82, 205203 (2010); J. Cottaar, R. Coehoorn, P.A. Bobbert, *Organ. Electr.* (accepted).
- [5] W. Chr. Germs, J. J. M. van der Holst, S. L. M. van Mensfoort, P. A. Bobbert, and R. Coehoorn, *Phys. Rev. B* 84, 165201 (2011).
- [6] J. J. M. van der Holst, F. W. A. van Oost, R. Coehoorn, and P. A. Bobbert, *Phys. Rev. B* 80 235202 (2009).
- [7] S. L. M. van Mensfoort, M. Carvelli, M. Megens, D. Wehenkel, M. Bartyzel, H. Greiner, R. A. J. Janssen and R. Coehoorn, *Nature Phot.* 4, 329 (2010); M. Carvelli, R.A.J. Janssen and R. Coehoorn, *J. Appl. Phys.* 110, 084512 (2011).

Asymptotic models for the transport of an electron gas strongly confined in a slab

Fanny Delebecque

Université Paul Sabatier, Institut de Mathématiques de Toulouse, équipe MIP,
118 route de Narbonne, 31062 Toulouse CEDEX 9, France

e-mail: `Fanny.Delebecque@math.univ-toulouse.fr`

The first model is a 3D Schrödinger Poisson system in the whole 3D space, singularly perturbed by a strong confinement potential acting on one direction. The obtained asymptotic model is a simple 2D Schrödinger equation coupled with a Poisson equation on whole density, concentrated on the transport plane. In a second model, the confinement is ensured by homogenous boundary conditions on both the wave function and Poisson potential. I will present time averaging technics used to get rid of the strong time oscillations due to the confinement.

Modeling boundary conditions for solving stationary Schrödinger equations

Xavier Antoine ⁽¹⁾, Christophe Besse ⁽²⁾, Matthias Ehrhardt ⁽³⁾, and Pauline Klein ⁽¹⁾

⁽¹⁾ Institut Elie Cartan Nancy, Nancy-Universität, CNRS UMR 7502, INRIA CORIDA Team,

⁽²⁾ Equipe Projet Simpaf - Inria CR Lille Nord Europe, Laboratoire Paul Painlevé, Unité Mixte de Recherche CNRS (UMR 8524), UFR de Mathématiques Pures et Appliquées, Université des Sciences et Technologies de Lille, Cité Scientifique,

⁽³⁾ Lehrstuhl für Angewandte Mathematik und Numerische Analysis, Fachbereich C – Mathematik und Naturwissenschaften, Bergische Universität Wuppertal, Gaußstr. 20, 42119 Wuppertal, Germany

e-mail: Ehrhardt@math.uni-wuppertal.de

In this talk we present some novel absorbing boundary conditions (ABCs) for modeling the solution of linear and nonlinear variable potentials one-dimensional stationary Schrödinger equations. Using pseudodifferential calculus and factorization theorems we construct a hierarchy of novel ABCs and generalize the well-known quantum transmitting boundary condition of Kirk and Lentner to the case of space-dependent potential. Moreover, we propose a rapidly converging iterative method based on finite elements suitable for computing scattering solutions and bound states. The accuracy of our new absorbing boundary conditions is investigated numerically for two different situations. The first problem is related to the computation of linear scattering problems. The second application concerns the computation of energies and ground-states for linear and nonlinear Schrödinger equations. It turns out that these absorbing boundary conditions and their variants lead to a higher accuracy than the usual Dirichlet boundary condition. Finally, our approach also offers the possibility to construct ABCs for higher dimensional problems.

REFERENCES

- [1] X. Antoine, A. Arnold, C. Besse, M. Ehrhardt, A. Schädle, *A review of transparent and artificial boundary conditions techniques for linear and nonlinear Schrödinger equations*, Commun. Comput. Phys. 4 (2008), 729–796. (open-access article)
- [2] X. Antoine, C. Besse and P. Klein, *Absorbing boundary conditions for the one-dimensional Schrödinger equation with an exterior repulsive potential*, J. Comput. Phys. 228 (2009), 312–335.
- [3] X. Antoine, C. Besse, M. Ehrhardt and P. Klein, *Modeling boundary conditions for solving stationary Schrödinger equations*, Preprint 10/04 of the chairs of Applied Mathematics & Numerical Analysis and Optimization & Approximation, University of Wuppertal, February 2010.
- [4] P. Klein, X. Antoine, C. Besse and M. Ehrhardt, *Absorbing boundary conditions for solving stationary Schrödinger equations*, Proceedings of the 16th European Conference on Mathematics for Industry, July 26-30, 2010, Wuppertal, Germany, Springer, Berlin, Heidelberg, 2011.
- [5] P. Klein, X. Antoine, C. Besse and M. Ehrhardt, *Absorbing boundary conditions for solving N -dimensional stationary Schrödinger equations with unbounded potentials and nonlinearities*, Commun. Comput. Phys. 10 (2011), 1280–1304.

Some approximate controllability results for bilinear systems

Sylvain Ervedoza

CNRS & Université Paul Sabatier, Institut de Mathématiques de Toulouse, Équipe MIP
118 route de Narbonne, 31062 Toulouse CEDEX 9, France

e-mail: `Sylvain.Ervedoza@math.univ-toulouse.fr`

This talk is complementary to the one of Karine Beauchard. Here, I will present a brief overview on approximate controllability results for bilinear systems. I will mainly focus on three approaches:

- a Lyapunov approach, developed by Nersesyan and Beauchard-Nersesyan;
- an approach based on the control properties of Galerkin approximations of the system, developed by Boscain, Caponigro, Chambrion, Mason, Sigalotti;
- an approach based on the control properties of an easy approximate system deduced by using the scales of the equation, based on the works by S.E., Puel.

Drift-Diffusion-Recombination Models for Excitonic Organic Photovoltaic Devices

Klemens Fellner

University of Graz, Institute for Mathematics and Scientific Computing,
Heinrichstr. 36, 8010 Graz, Austria

e-mail: `Klemens.Fellner@uni-graz.at`

Joint work with Dan Brinkman, Peter Markowich (DAMTP, University of Cambridge) and Marie-Therese Wolfram (University of Vienna).

We present some recent results and ongoing work on drift-diffusion-reaction systems modelling organic photovoltaic devices. While classical semiconductors show recombination typically throughout the whole device feature organic photovoltaic devices significant charge generation only in the very proximity of an interface between two different organic polymers. We discuss basic questions of modelling, existence and stationary states and current voltage characteristics. Moreover, we present some interesting asymptotic approximations and discuss the use of entropy.

The Oskar3 project

Klaus Gärtner

Weierstrass Institute,
Mohrenstr. 39, 10117 Berlin, Germany

e-mail: Klaus.Gaertner@wias-berlin.de

The main goal of the project is to improve algorithms for 3d semiconductor device simulation by exploiting the nice properties of the analytic problem. Main areas of interest from the algorithmic point of view are:

- a) discretizations, especially cases requiring non Boltzmann statistics;
- b) solving linear systems;
- c) construction of anisotropic Delaunay grids (Gajewski: "The best convergence acceleration is a good grid").

The goal is understanding and solving selected application problems of research interest. Collaborative, true experts in the special field of application and extreme requirements with respect to at least one of the points a) to c) are essential ingredients of a good obstacle course training brains. We use

- A) silicon sensors for high energy and astro physics (high precision, huge volumes, every design is very different and time scales allow discussions - related to b, c);
- B) VCSEL hetero structures (representatives of complex heterostructures with exotic materials, mainly related to a);
- C) organic semiconductors (the hope for a black paint producing electricity over five years with 10 % efficiency at nearly no cost—for the presently used models no limit of practical interest results in the Boltzmann case, hence a)

to adjust the focus.

A discussion of the used and unused mathematical properties will be given, some examples will help to judge the status reached at WIAS.

An electronic model for solar cells taking into account active interfaces

Annegret Glitzky

Weierstrass Institute,

Mohrenstr. 39, 10117 Berlin, Germany

e-mail: `Annegret.Glitzky@wias-berlin.de`

In the talk we discuss an electronic model for solar cells taking into account heterostructures with active interfaces and energy resolved volume and interface trap densities. The model introduced by [1] consists of continuity equations for electrons and holes with thermionic emission transfer conditions at the interface and of ODEs for the trap densities with energy level and spatial position as parameters, where the right hand sides contain generation-recombination as well as ionization reactions. This system is coupled with a Poisson equation for the electrostatic potential.

We study the thermodynamic correctness of the model and motivate a priori estimates for the solutions to the evolution system. Moreover, existence and uniqueness results of weak solutions of the problem are presented. For the existence proof regularized problems have to be solved and bounds of the corresponding solution not depending on the regularization level have to be verified, for details see [2].

REFERENCES

- [1] R. Stangl, C. Leendertz, and J. Haschke, *Numerical simulation of solar cells and solar cell characterization methods: the open-source on demand program AFORS-HET*, Solar Energy (R. D. Rugeanu, ed.), INTECH, Croatia, 2010, pp. 319–352.
- [2] A. Glitzky, *An electronic model for solar cells including active interfaces and energy resolved defect densities*, SIAM J. Math. Anal., to appear, WIAS Preprint 1663.

Discretization of Time Dependent Quantum Systems: The Evolution Operators

Joseph Jerome

Northwestern University,
2033 Sheridan Road, IL 60208-2730 Evanston, USA

e-mail: jwj@math.northwestern.edu

We discuss time dependent quantum systems. Our emphasis is on the linear problem, but our work may be viewed as a framework for several models, including linear iterations involved in time dependent density functional theory (TDDFT), the Hartree-Fock model, or other quantum models. A key aspect of the analysis of the algorithms is the use of time-ordered evolution operators, which allow for both a well-posed problem and its approximation. We discuss the available theory at the outset, and proceed to apply the theory systematically, via approximations. We discuss rigorous use of the rectangular and Gauss quadrature rules. Our work is consistent with first-principle, real-time propagation of electronic states, aimed at finding the non-linear responses of quantum systems. Given the increasing importance of quantum modeling in applications, we are attempting to identify some of the basic issues to be addressed analytically.

On a model of magnetization switching driven by a spin current: modeling and numerical simulations

Clément Jourdana

Université Joseph Fourier - Grenoble 1
51 rue des Mathématiques, BP 53, 38041 Grenoble CEDEX 9, France
e-mail: `Clement.Jourdana@imag.fr`

This is a joint work with Naoufel Ben Abdallah, Elise Fouassier and David Sanchez.

We study a model, introduced in [1], of magnetization switching induced by a spin polarized current without applying an external magnetic field. It consists of a coupled system for the local magnetization (that satisfies a Landau-Lifshitz equation with an additional spin torque) and for the spin density (that satisfies a diffusion equation with a term for the procession phenomenon around the magnetization).

We first write the one dimensional model in an adimensionalized form, using a small parameter ϵ . We then explain the various time and space scales involved in the studied phenomena. Taking into account these scales, we construct an appropriate numerical scheme, that allows us to recover numerically various results of physical experiments. Finally, we perform a formal asymptotic study as ϵ tends to 0. We thus obtain an approximate limit model that we compare with the original model via numerical simulations.

REFERENCES

- [1] S. Zhang, P. M. Levy and A. Fert, *Mechanisms of spin-polarized current-driven magnetization switching*, Phys. Rev. Lett. **88**(23), 2002.

Semi-classical modeling of quantum dot lasers with microscopic treatment of Coulomb scattering

Thomas Koprucki

Weierstrass Institute,
Mohrenstr. 39, 10117 Berlin, Germany

e-mail: `Thomas.Koprucki@wias-berlin.de`

Drift-diffusion models provide a semi-classical description of the carrier transport in semiconductor devices. Self-consistently coupled to equations for the optical field they are established models for the simulation of semiconductor lasers. We consider quantum dot lasers with optically active regions consisting of quantum dots grown on a wetting layer. We present a drift-diffusion based modeling approach for the simulation of quantum dot lasers which uses a multi-species description of the carriers along the quantum dot active region and includes microscopically determined scattering rates describing the capture of the carriers into the quantum dots by Coulomb scattering. The presented results are joint work with U. Bandelow, K. Gärtner, A. Wilms, and A. Knorr (TU Berlin).

REFERENCES

- [1] U. Bandelow, H. Gajewski, R. Hünlich, *Fabry–Perot Lasers: Thermodynamics-Based Modeling*. Appeared in: *Optoelectronic Devices — Advanced Simulation and Analysis*, J. Piprek, ed., Springer, New York, 2005, pp. 63–85.
- [2] Th. Koprucki, A. Wilms, A. Knorr, U. Bandelow, *Modeling of quantum dot lasers with microscopic treatment of Coulomb effects*. *Opt. Quant. Electron.* **42**, pp. 777–783, 2011.

Asymptotic spectral analysis in quantum waveguides with heterogeneous fibers

Carolin Kreisbeck and Luísa Mascarenhas

Universidade Nova de Lisboa, Faculdade de Ciências e Tecnologia, Departamento de Matemática and Centro de Matemática e Aplicações
Quinta da torre, 2829-516 Caparica, Portugal

e-mail: `Carolin.Kreisbeck@mathematik.uni-regensburg.de`

In recent years nanowires have become an active field of research in the physics community. We discuss two settings in which heterogeneous structures affect propagation through a quantum waveguide, highlighting their interaction with curvature and torsion.

First, we analyze the macroscopic behavior of a wire made of composite fibers with microscopic periodic texture. Mathematically, this amounts to determining the asymptotic behavior of the spectrum of an elliptic Dirichlet eigenvalue problem with finely oscillating coefficients in a three-dimensional tube with shrinking cross section. This two-scale problem is treated in a variational framework performing homogenization and 3d-1d dimension reduction by means of Γ -convergence. We show that the effective one-dimensional limit problem is of Sturm-Liouville type and give an explicit formula for the underlying potential.

Second, for a waveguide with inhomogeneous cross section that does not oscillate as before but shrinks along with the wire in the limit process one observes qualitatively different behavior. Stable propagation through such a modulated wire can be guaranteed by imposing symmetry on the inhomogeneities. In the non-symmetric case, localization is in general the dominant effect. We prove, however, that propagation does happen if a specific relation between the shape of the waveguide and its transversal inhomogeneities holds. Thus, propagation can be enforced by simple geometrical operations, i.e. bending and twisting, which provides a theoretical concept for the building of quantum switching devices.

Artificial interface conditions in quantum transport models

Andrea Mantile

Université de Reims, Laboratoire de Mathématiques, U.F.R. Sciences Exactes et Naturelles,
Moulin de la Housse - BP 1039, 51687 Reims CEDEX 2, France

e-mail: `Andrea.Mantile@univ-reims.fr`

In a recent collaboration with F. Nier and A. Faraj, it has been shown that a simple modification of the Laplacian through artificial interface conditions allows an alternative approach to the adiabatic evolution of quantum resonances. The use of this modified framework, may hopefully provide with effective equations for the non-linear dynamics of Schrödinger-Poisson systems in the regime of quantum wells in a semiclassical island. In this perspective, it is important to control the deformations effects introduced on the spectrum and on the time propagator by such interface conditions. In particular we are interested in uniform-in-time estimates of the perturbed semigroup. The main difficulty is due to the non-selfadjoint character of our class of operators involving a lack of accretivity for the corresponding generator of the quantum dynamics. In this framework, a standard approach would only provide with finite-time estimates for the dynamical system. Our approach consists in constructing intertwining operators leading to a dynamical comparison between the modified non-selfadjoint model and the corresponding “physical” Hamiltonian.

A plane-wave based, generalised formulation of continuum elasticity and multiband k·p-models

Oliver Marquardt ^(1,*), Christoph Freysoldt ⁽²⁾, Sixten Boeck ⁽³⁾, Tilmann Hickel ⁽²⁾,
Stefan Schulz ⁽¹⁾, Jörg Neugebauer ⁽²⁾, and Eoin P. O'Reilly ^(1,4)

⁽¹⁾ Tyndall National Institute, Cork, Ireland

⁽²⁾ Max-Planck-Institut für Eisenforschung, Düsseldorf, Germany

⁽³⁾ Gemmantics IT-Consulting, Erkrath, Germany

⁽⁴⁾ University College Cork, Ireland

^(*) Now at: Paul-Drude-Institut Berlin,
Hausvogteiplatz 5-7, 10117 Berlin, Germany

e-mail: O.Marquardt@yahoo.ie

Multiband k·p-models have been used successfully for the theoretical modelling of semiconductor nanostructures of different shapes, sizes, and for a wide range of materials. These models offer a computationally efficient approach to the electronic properties of quantum dots, wires, and wells on the basis of a continuum representation of the system under consideration. We will present a general formulation of k·p-models, that allows to study arbitrary nanostructure geometries, sizes, and material compositions employing an n-band Hamiltonian. The details of the Hamiltonian as well as the nanostructure geometry and composition are defined in user-generated input files, which adds great flexibility to tailor the calculation to the question of interest. In this manner, simulations using the well-established eight-band k·p-model as well as more sophisticated 14-band approaches or simple one-band effective mass models are straightforwardly possible. To account for elastic and piezoelectric properties, a continuum elasticity model has also been implemented. The resulting strain and piezoelectric potentials can then be considered in the calculation of the electronic properties. Both the continuum elasticity model as well as the multiband k·p-formalism were implemented in a plane-wave framework within the existing S/PHI/nX physics simulation package. This allows to make use of existing, highly efficient minimisation algorithms and to apply a reciprocal space representation for gradient operations, allowing for a significant boost in efficiency. The specific issues that occur in a plane-wave based implementation of these two continuum models will be furthermore discussed in example simulations that have been performed with our approach.

**A new decomposition for the Schrödinger equation in the
semiclassical asymptotics and an associated AP numerical
scheme**

Florian Méhats

IRMAR. University of Rennes 1,
Campus de Beaulieu, 35000 Rennes CEDEX, France
e-mail: Florian.Mehats@univ-rennes1.fr

This is a joint work with Ch. Besse and R. Carles. We propose a new decomposition “à la Grenier” for linear and nonlinear Schrödinger equations, which leads to a locally well-posed problem, independently of the semiclassical parameter. Based on this formulation, we give a numerical scheme which is Asymptotic Preserving in the semiclassical limit (before caustics). The performances of this method will be compared with the ones of the standard splitting method.

Using gradient structures for modeling semiconductors

Alexander Mielke

Weierstrass Institute,
Mohrenstr. 39, 10117 Berlin, Germany

e-mail: Alexander.Mielke@wias-berlin.de

Joint work with Annegret Glitzky and Matthias Liero (WIAS Berlin).

We show that classical energy-drift-diffusion systems used for modeling of semiconductor devices can be written as gradient systems in the form

$$\partial_t \mathbf{u} = -\mathcal{K}(\mathbf{u})D\Phi(\mathbf{u}),$$

where \mathbf{u} is a vector of densities of the different species and possibly the temperature (or the internal energy), cf. [Mie11]. While the potential Φ is the free energy in the isothermal case and the negative entropy in the temperature-dependent case, the Onsager operator $\mathcal{K}(\mathbf{u}) = G(\mathbf{u})^{-1}$ is the inverse of the metric tensor G and has an additive split into the different dissipation mechanics:

$$\mathcal{K}(\mathbf{u})\xi = \mathcal{K}_{\text{diff}}(\mathbf{u})\xi + \mathcal{K}_{\text{react}}(\mathbf{u})\xi = -\text{div}(\mathbb{M}(\mathbf{u})\nabla\xi) + \mathbb{K}(\mathbf{u})\xi.$$

In the case of Fermi-Dirac statistics this modeling lead to new and more consistent modeling of diffusion enhancement and reaction kinetics. In particular, the modeling of semiconductors via gradient system (Φ, \mathcal{K}) yields always thermodynamically consistent models and highlights the different physical aspects of the theory in a clear way.

It is shown in [GIM12, Mie12] that gradient structures can also used effectively to model hybrid models with nontrivial bulk-interface interaction and active interfaces, which are need in photovoltaics. Finally, we discuss how the theory of geodesically λ -convex gradient can be applied in this field, cf. [LiM12].

REFERENCES

- [GIM12] A. GLITZKY and A. MIELKE. A gradient structure for systems coupling reaction-diffusion effects in bulk and interfaces. *Z. angew. Math. Phys. (ZAMP)*, 2012. To appear. WIAS preprint 1603 (April 2011, doi 10.1007/s00033-0012-207-y).
- [LiM12] M. LIERO and A. MIELKE. Gradient structures and geodesic convexity for reaction-diffusion systems. Submitted. WIAS preprint 1701 (April 2012).
- [Mie11] A. MIELKE. A gradient structure for reaction-diffusion systems and for energy-drift-diffusion systems. *Nonlinearity*, 24, 1329–1346, 2011.
- [Mie12] A. MIELKE. Thermomechanical modeling of energy-reaction-diffusion systems, including bulk-interface interactions. *Discr. Cont. Dynam. Systems Ser. S*, 2012. To appear. WIAS preprint 1661 (Nov. 2011).

An application of the Landauer-Büttiker formula to photon emitting and absorbing systems

Hagen Neidhardt

Weierstrass Institute,
Mohrenstr. 39, 10117 Berlin, Germany

e-mail: `Hagen.Neidhardt@wias-berlin.de`

The abstract Landauer-Büttiker formula is applied to a Jaynes-Cummings model coupled to leads. As a result, we obtain formulas for electron and photon currents in terms of the partial cross-sections of a scattering system for which the unperturbed and perturbed Hamiltonians correspond to systems where the leads are either decoupled or coupled to the Jaynes-Cummings model. It is shown that the resolvent difference of the coupled and decoupled Hamiltonians is a trace class operator. The electron and photon currents are analyzed in detail.

The talk is based on a common paper with Horia Cornean, Lukas Wilhelm and Valentin Zagrebnov.

Asymptotic analysis of a 1D KP system

Francis Nier

IRMAR. University of Rennes 1,
Campus de Beaulieu, 35000 Rennes CEDEX, France

e-mail: `Francis.Nier@univ-rennes1.fr`

In this talk we will show how the semiclassical analysis allows to elucidate the complex subband structure of a KP system for bipolar semiconductors.

WIAS-TeSCA simulations in photovoltaics for a point contact concept of heterojunction thin film solar cells

Reiner Nürnberg

Weierstrass Institute,
Mohrenstr. 39, 10117 Berlin, Germany

e-mail: `Reiner.Nuernberg@wias-berlin.de`

For a long time 1D simulation software such as SCAPS were the widely used semiconductor simulation tools in photovoltaics because of the relatively simple structure of solar cells. A novel point contact concept is an example for the necessary and useful application of 2D and 3D simulation in this field.

The photovoltaic performance of the wider-band gap chalcopyrite based thin film solar cells is limited by the recombination at the interface between the absorber and the buffer layer. In principle, the heterojunction should exhibit a high density of defect states at this interface due to lattice mismatch. The concept of a point-contact geometry at the front side of thin film chalcopyrite solar cell devices includes (i) an interfacial defect state passivation and (ii) a current transport only through these point contacts that cover a few percent of the device area, while the rest of the surface is electronically passivated by a dielectric coating. The introduction of a point-contact type structure allows for the relaxation of the requirements at the absorber/buffer heterointerface and could dramatically increase performance of heterojunction thin film solar cells.

The software WIAS-TeSCA for the numerical simulation of charge transfer processes in semiconductor structures, in principle, is suitable for use in thin film photovoltaics. It is based on the drift-diffusion model and allows for heterostructures. Modeling of the structured passivation layer with point contacts can be simplified by taking a structural unit with cylindrical symmetry so that the modeling is performed in 2D but provides results which will be a good approximation to the 3D structure. In the present model study, especially, the effect of the interfacial defect level and density, and of the point contact radius and area ratio have been investigated.

Optical Recombination in Semiconductor Nanostructures

Eoin O'Reilly ^(1,2), Stefan Schulz ⁽¹⁾, Sarah Krug ⁽¹⁾, Oliver Marquardt ^(1,3), and Miguel Caro ^(1,2)

⁽¹⁾ Tyndall National Institute, Cork, Ireland

⁽²⁾ Department of Physics, University College Cork, Ireland

⁽³⁾ Now at Paul-Drude-Institut Berlin, Germany

e-mail: Eoin.OReilly@tyndall.ie

Two of the key current challenges in photonics are the development of entangled photon sources and the development of lasers and efficient light sources across the full visible spectrum. There is considerable interest in the use of semiconductor quantum dots (QDs) to achieve both these aims. Here we discuss the issues involved to investigate the piezoelectric potential and electronic structure of III-N wurtzite and of site-controlled (111)-oriented zinc-blende dots. We investigate the single-particle electronic structure of (111)-oriented InGaAs/GaAs QDs by means of an 8-band $k \cdot p$ model. We then discuss how this model needs to be extended in order to identify how best to engineer the dots to give orthogonally polarized exciton states that are degenerate (zero fine structure splitting), thereby supporting the emission of entangled photons. This requires both to treat dots where there are only a few confined electron but many confined hole states, as well as to incorporate the effects of disorder into the system. We highlight that accurate knowledge of physical parameters remains an issue for such dots [1]. We then turn to consider the influence of the built-in potential in wurtzite structures. We first compare the elastic and the first-order piezoelectric tensor of the (111)-zinc-blende systems with the corresponding quantities in a wurtzite structure and point out similarities and differences [2]. We then turn to consider the analysis of the built-in potential and its consequences for polar and for non-polar III-N QDs [3]. We propose that the growth of InGaN/GaN QD structures can lead to a significant reduction in the built-in potential compared to InGaN/GaN quantum well structures, allowing efficient optical emission to longer wavelengths [4]. We also highlight that the calculated field values and their effects are very sensitive to details of the dot shape and composition in such heterostructures.

REFERENCES

- [1] L. O. Mereni, O. Marquardt, G. Juska, V. Dimastrodonato, E. P. O'Reilly and E. Pelucchi, *Phys. Rev. B* **85**, 155453 (2012)
- [2] S. Schulz, M. A. Caro, E. P. O'Reilly, and O. Marquardt, *Phys. Rev. B* **84**, 125312 (2011)
- [3] S. Schulz, A. Berube, and E. P. O'Reilly, *Phys. Rev. B* **79**, 081401 (2009)
- [4] S. Schulz, and E. P. O'Reilly, *Phys. Rev. B* **82**, 033411 (2010)

Problems with the Gaussian density of states in organics: definite solution, ambiguous, and hardly any solution

Gernot Paasch⁽¹⁾ and **Susanne Scheinert**⁽²⁾

⁽¹⁾ Leibniz Institute for Solid State and Materials Research IFW Dresden,
PF 27 01 16, 01171 Dresden, Germany

⁽²⁾ Ilmenau Technical University, Institute of Solid State Electronics
e-mail: G.Paasch@ifw-dresden.de

Disordered organic materials are the basis of organic electronics with organic light emitting diodes for displays and lighting, organic field-effect transistors (OFET) and circuits made with them, and organic solar cells. Almost twenty years since the pioneering work of Bäessler [1] on hopping transport in disordered organics it is widely accepted that the density of states (DOS) is a Gaussian. From the application of various numerical models for the mobility, a variance of the order of 100 meV seems to be typical. Although a large number of device modelling using the Gaussian DOS has been published, support of design and technology of devices by full two-dimensional simulation is required. Until recently this possibility was not completely implemented in the simulation programs as sDEVICE [2] or ATLAS [3] due to peculiarities of the Gaussian DOS. We discuss here three problems important for implementation and application.

Fast and efficient numerical simulation requires analytical expressions for the carrier density (as the Fermi integral for the square-root DOS). Alternatively, the numerical integration of the product of the DOS and the Fermi distribution is in principle possible and for some cases of non-parabolic bands available. For the Gaussian DOS we developed an analytical approximation for the carrier density (Gauss-Fermi integral) [4] which is meanwhile implemented in sDEVICE [2]. Detailed testing and comparison with experimental OFET current characteristics is in progress but is also connected with further problems.

For the Gaussian DOS an enormous amount of theoretical work on the mobility has been done and published. The extended Gaussian disorder model (EGDM) [5, 6] (describing the dependence of the mobility on temperature, carrier concentration, and field) was to our opinion a first highlight after Bäessler, followed later by the extended correlated disorder model (ECDM). However there exist many other models, also from recent times, with different dependencies. There remain serious ambiguities: Which model is applicable to a given organic material, which parameters should be used or how one can determine them, are the models sufficient for application at high concentrations and high fields?

Most problematic is the question what the band edge is and how to determine it. As a routine method the valence band edge is determined by ultraviolet photoemission (UPS). However, this method probes regions of the DOS which are unoccupied in devices and transport takes place at much lower energies not accessible up to date by UPS. Moreover, the upper part of the UPS spectrum below the gap is rather broad compared with the narrow hopping transport DOS. Thus we do not really know where this DOS is situated. We supposed [7] that the transport DOS is the narrower tail of the broader UPS distribution. With some uncertainty one can then fix the position of the Gaussian transport DOS on the binding energy scale. But there remains a further problem, namely where is the band edge situated relative to the maximum of the Gaussian DOS. The method presented in [7] is not unique and we will discuss several other possibilities, but until now without a definite choice. It might be that the injection into a Gaussian DOS requires a new theoretical treatment and its implementation beyond the traditional one.

REFERENCES

- [1] H. Bäessler, Phys. Status Solidi **B 175**, 15 (1993).

- [2] Sentaurus Device User Guide, Synopsys, www.synopsys.com, 2008.
- [3] P. Henry, Atlas User Manual, Silvaco International, Santa Clara, CA, 2008.
- [4] G. Paasch and S. Scheinert, *J. Appl. Phys.* **107**, 104501 (2010).
- [5] R. Coehoorn, W. F. Pasveer, P. A. Bobbert, and M. A. J. Michels, *Phys. Rev.* **B 72**, 155206 (2005).
- [6] W. F. Pasveer, J. Cottaar, C. Tanase, R. Coehoorn, P. A. Bobbert, P. W. M. Blom, D. M. de Leeuw, and M. A. J. Michels, *Phys. Rev. Lett.* **94**, 206601 (2005).
- [7] S. Scheinert, M. Grobosch, G. Paasch, I. Hörselmann, M. Kupfer, and J. Bartsch, *J. Appl. Phys.* **111**, 064502 (2012).

Utilizing Wavelets for the Solutions of High-dimensional Semiconductor Transport Equations

Vincent Peikert and Andreas Schenk

Integrated Systems Laboratory, ETH Zürich,
 Gloriastr. 35, CH-8092 Zürich, Switzerland

e-mail: VPeikert@iis.ee.ethz.ch

Introduction. Transport equations (Boltzmann, Wigner) on 6-dimensional phase spaces describe the behavior of charge carriers in semiconductors and form the basis for Technology Computer Aided Design (TCAD) on device level. However, the numerical solution of these high-dimensional equations involves huge matrices which practically limits contemporary direct solvers to problems with only one real space dimension. Since the integral operator in the Wigner transport equation (WTE) produces full matrices, 2-dimensional simulations are only feasible for the Boltzmann transport equation (BTE) so far. Even for the BTE the simulation of 2D devices calls for coarse meshes and geometrical simplicity which sets limitations to the accuracy and the practical application for the new TCAD challenges introduced by the rapidly increasing importance of 3D devices. By far the most established numerical BTE approach is the Spherical Harmonics Expansion Method (SHE) which is based on a representation of the momentum space in spherical coordinates and an expansion of the angular variables in SH. The SH Y_l^m build a dense basis in L_2^S on the sphere: $V_S^0 \subset V_S^1 \subset \dots \subset V_S^N \subset L_2^S$, with $V_S^N := \text{span}\{Y_l^m : l = 0, \dots, N; m = -l, \dots, l\}$. Any V_S^k may be decomposed as $V_S^k = W_S^k \oplus V_S^{k-1}$, with $W_S^k \perp V_S^{k-1}$. A special property of SH is that the detail spaces W_S^k are spanned by SH Y_l^m themselves. The SH coefficients only contain the additional detail information to the orders before and decay quickly with the detail order (typically only 7 orders are necessary to simulate silicon transistors). The SH coefficient size has recently been used as an adaptive criterion for a variable order SHE expansion and the feasibility of 3D SHE simulations have been investigated [1].

Multi-Wavelets. All contemporary numerical BTE approaches apply piecewise polynomials (pp) as basis functions. Usually, they apply piecewise constants (pc) in all dimensions (the SHE applies pc in the energy and real space directions). Recently, the Discontinuous Galerkin (DG) method for the BTE has been proposed [2]. The DG method is flux conserving and is (in contrast to SHE) stable with high-order pp in real space and energy directions (hp-stable) and is therefore a promising alternative. Let V_n^k be the space of all one-dimensional pp of the degree less than k on a uniform mesh (to keep the notation simple) with 2^n equidistant intervals within $[0, 1]$. By bisecting the mesh, pp build a dense basis in L_2 and fulfill the nested structure $V_0^k \subset V_1^k \dots \subset V_n^k \dots \subset L_2([0, 1])$. It is possible to decompose V_{n+1}^k into V_n^k and a perpendicular space W_n^k which only contains the additional details of V_{n+1}^k compared to V_n^k : $V_n^k \oplus W_n^k = V_{n+1}^k$, with $V_n^k \perp W_n^k$. Recursively, V_n^k can be decomposed by detail spaces $V_n^k = V_0^k \oplus W_0^k \oplus W_1^k \oplus \dots \oplus W_{n-1}^k$. W_l^k is spanned by $2^{l-1}k$ detail functions $\psi_{l,i,j}$ which are called Multi-Wavelets (MWs) (see [3]) where l indexes the detail order, j indexes the vanishing moment order and i indexes the position.

Multi-Wavelet Discontinuous Galerkin Method. Instead of using $\text{pp} \in V_n^k$ (nodal DG: NDG), the solution of the BTE Φ could be expanded in tensor products of one-dimensional MW bases (Multi-Wavelet Discontinuous Galerkin method: MWDG):

$$\Phi(x, y, z, \mu, \phi, \omega) = \sum_{I,J,K,L} \alpha_{I,J,K,L} \psi_I^{x,y,z} \psi_J^\mu \psi_K^\phi \psi_L^\omega$$

(multi-indexes I are used for $\{l, i, j\}$). x, y, z are the real space coordinates, μ is the cosine of the polar angle, ϕ is the azimuthal angle and ω is the energy. Note that a tensor product is only build within the momentum space and with the real space. Within the real space, unstructured grids are typically necessary and the construction of MWs x, y, z on unstructured grids can be done following the procedure in [4]. A strongly growing community utilizes the detail and vanishing moment properties of wavelets for the adaptive compression for different complex problems. For the solution of PDEs, hierarchical bases build the “driving force” behind modern adaptive solvers and pre-conditioners. However, of particular interest are the properties of wavelet tensor bases. The coefficient decay (as a property of the detail nature of wavelets) is strongly enhanced in multi-dimensional tensor bases so that most wavelets can be canceled. This is referred to as high-dimensional wavelets compression (HWC) here and is called adaptive sparse grids in literature [5]. Preliminary work on adaptive sparse grids concentrated on FEM formulations and on scalar wavelets rather than on the flux conserving formulations that are necessary for the solution of transport equations. Furthermore, in contrast to scalar wavelets, MWs have an additional vanishing moment hierarchy within each detail to represent high order pp and compression leads to some kind of super sparse grid. This is referred to as vanishing moment compression (VMC) here. Preliminary work on DG and MWs (e.g. [6]) use MW coefficients as adaption criterion, but always keep a full grid so that the MWDG and the NDG stay equivalent. MWDG simulations of 1D devices with pc MWs in [7] show that adaptive HWC can compress the basis by about 99%. Furthermore, MWDG simulations of 1D devices with higher-order polynomials in [8] perform significantly better even for uncompressed MWDG (up to 96% saving for only 1.4% error in the current) and HWC and VMC can compress the basis by additional 90%. Future Work. Even larger compression rates are expected for future full hp-adaptive MWDG simulations of 3D devices. A drawback of MWs is the higher coupling of difference operators compared to a conventional pp basis which can reduce the advantages of MWs compression. However, the product structure of the BTE within the momentum space and with the real space enables (in combination with the corresponding product basis above) to decompose the system matrix into a Kronecker product $M = M_{x,y,z} \otimes M_{\mu} \otimes M_{\phi} \otimes M_{\omega}$. Utilizing the distributive property enables to calculate a matrix vector product without building up the matrix M and with a complexity that scales with the number of unknowns only (unidirectional principle [9]). Hence, an iterative solver in combination with a multi-level pre-conditioner make the problem independent of the couplings within the subdimensions. If quantum mechanical effects become important, the WTE has to be solved. Since the solution of the BTE and the WTE stay similar, similar high MW compression rates are expected for the WTE. MWs are advantageous for the discretization of integral operators such as the Wigner potential, since the matrices show a quasi-sparse structure. Additionally, the product structure of the Wigner potential enables the unidirectional principle so that 3D MWDG WTE solvers could become feasible in the future.

REFERENCES

- [1] K. Rupp and T. Grasser, A. Jüngel, IEDM, Washington DC, USA, (2011).
- [2] Y. Cheng, I.M. Gamba, A. Majorana, C. Shu, Comput. Methods Appl. Mech. Engrg. 198, 3130 (2009).
- [3] B. Alpert, SIAM Journ. Math. Anal. 24, 246 (1993).
- [4] W. Sweldens, Applied and Computational Harmonic Analysis 3, 186 (1996).
- [5] H.-J. Bungartz, M. Griebel, Acta Numerica 13, 147, (2004).
- [6] A. Shelton, PhD Thesis, Georgia Institute of Technology, Atlanta (2008).
- [7] V. Peikert, A. Schenk, SISPAD, Osaka, Japan, (2011).
- [8] V. Peikert, A. Schenk, IWCE, Madison (WI), USA, (2012).
- [9] R. Balder PhD Thesis, TU Munich (1994).

Electron transport in strongly confined nanostructures: modeling and simulation

Clément Jourdana ⁽¹⁾ and Paola Pietra ⁽²⁾

⁽¹⁾ Université Joseph Fourier - Grenoble 1

⁽²⁾ Istituto di Matematica Applicata e Tecnologie Informatiche “Enrico Magenes”, CNR,
via Ferrata 1, 27100 Pavia, Italy

e-mail: Pietra@imati.cnr.it

We present an effective mass model, derived in [1], describing the ballistic transport of electrons in ultra-scaled confined nanostructures. Due to the strong confinement, the crystal lattice is considered periodic only in the one dimensional transport direction and an atomistic description of the entire cross-section is given. The model consists of a sequence of 1D device dependent Schrödinger equations, one for each energy band, which retain the effects of the confinement and of the transversal crystal structure. In order to model field effect transistors, self-consistent computations include the resolution, in the whole 3D domain, of a Poisson equation. Simulations of the electron transport in semiconducting single-walled carbon nanotubes are presented, using a classical-quantum hybrid strategy. The ballistic effective mass model is used only in the active region, and it is coupled to a drift-diffusion model (derived for strongly confined nanostructures in [2]) in the source and drain regions. Appropriate current preserving interface conditions (as in [3]) are used.

REFERENCES

- [1] N. Ben Abdallah, C. Jourdana, P. Pietra, *An effective mass model for the simulation of ultra-scaled confined devices*. to appear in Math. Models Methods Appl. Sci.
- [2] C. Jourdana, N. Vauchelet, 2011. *Analysis of a diffusive effective mass model for strongly confined nanowires* Kinet. Relat. Models **4**, 1121-1142.
- [3] M. Baro, N. Ben Abdallah, P. Degond, and A. El Ayyadi, 2005. *A 1d coupled Schrödinger drift-diffusion model including collisions* J. Comput. Phys. **203**, 129-153.

On the moment problem for quantum hydrodynamics

Olivier Pinaud

Stanford University, Department of Mathematics & Universite Lyon 1,
Building 380, office 382V, CA 94305 Stanford, USA

e-mail: `Pinaud@math.univ-lyon1.fr`

We address the following inverse problem in quantum statistical physics: does the quantum free energy (von Neumann entropy + kinetic energy) admit a unique minimizer among the density operators having a given local density $n(x)$? We give a positive answer in various configurations and characterize the minimizer in dimension one. The problem is related to a recent theory developed by P. Degond and C. Ringhofer concerning the derivation of quantum hydrodynamics models. This is joint work with F. Mehats.

Linear and non-linear eigensolver-based techniques for quantum modeling and simulations: applications to carbon-based materials

Eric Polizzi

Department of Electrical and Computer Engineering, University of Massachusetts,
Marcus 201, 100 Natural Resources Road, MA 01003 Amherst, USA

e-mail: Polizzi@ecs.umass.edu

In a recent article [Phys. Rev. B 79, 115112, (2009)], the FEAST algorithm has been presented as a general purpose eigenvalue solver which combines accuracy, robustness, high-performance and (linear) parallel scalability. Here, FEAST will be presented beyond the “black-box” solver as a fundamental modeling framework for the electronic structure problem spanning the research fields of solid-state physics and quantum chemistry.

First, it will be shown that FEAST can naturally address the original numerical complexity of the electronic structure problem as formulated by Slater in 1937 [Phys. Rev. 51, 846-851, (1937)]. The non-linear eigenvalue problem of type $\mathbf{A}(\lambda)\mathbf{x} = \lambda\mathbf{B}\mathbf{x}$ arising from the muffin-tin decomposition of the real-space domain is first derived and then reformulated to be solved exactly within the FEAST framework. This new framework is presented as a fundamental and practical solution for performing both accurate and scalable electronic structure calculations, bypassing the various issues of using traditional approaches such as linearization and pseudopotential techniques.

Second, we present a modification of the FEAST algorithm that solves the full non-linear eigenvector problem of type $\mathbf{A}(\mathbf{x})\mathbf{x} = \lambda\mathbf{B}\mathbf{x}$ (i.e. Schrödinger/Poisson) as an efficient and robust alternative to traditional self-consistent field method (such as Newton-Broyden, Anderson mixing, direct inversion of the iterative subspace and mixed Pulay techniques). The resulting approach is potentially one order of magnitude faster than conventional iterative methods, and it converges to the correct solution regardless of the choice of initial guess.

Finally, we present and discuss FEAST-based time-domain propagation techniques for time dependent quantum problem. Efficient time-dependent simulations have become increasingly important for characterizing the electron dynamics under time dependent external perturbations such as electromagnetic fields, pulsed lasers, and particle scattering. Using innovative spectral decomposition and discretization of time-ordered evolution operators, we aim to address the traditional limitations in term of trade-off between robustness and performances in such calculations.

In order to illustrate the efficiency of the FEAST framework, numerical examples are provided for various molecules and carbon-based materials using our in-house all-electron finite element implementation for the DFT/Kohn-Sham/LDA problem.

Multiscale Modeling of Heterojunction Organic Photovoltaic Devices

Matteo Porro

Dipartimento di Matematica “Francesco Brioschi”, Politecnico di Milano,
Piazza L. da Vinci 32, 20133 Milano, Italy

Center for Nano Science and Technology @PoliMi, Istituto Italiano di Tecnologia,
via Pascoli 70/3, 20133 Milano, Italy

e-mail: Matteo.Porro@mail.polimi.it

In this lecture, we present our most recent results in the modeling and simulation of heterojunction organic solar cells [2]. The proposed mathematical model is an extension of that introduced in [1] for planar heterojunction solar cells to treat arbitrary multi-dimensional morphologies. Our approach consists in describing diffusion and transport phenomena of excited states and charge carriers in a heterogeneous domain Ω of \mathbb{R}^d , $d = 1, 2, 3$, constituted by two non-overlapping subdomains Ω_i , $i = 1, 2$, using mass balance equations supplied with the Drift-Diffusion model for flux densities. Charge dissociation and recombination phenomena occurring in a thin region around the $(d-1)$ -interface are lumped with a scale transition procedure into flux transmission conditions. This approach is twofold beneficial since it allows to account for the local dependence of dissociation on the orientation of the electric field and to reduce the computational effort compared to previous multi-dimensional models [5, 6]. The resulting system of fully coupled nonlinear PDEs/ODEs is faced with Rothe’s method using a quasi-Newton technique for linearization, the Edge-Average Finite Element method studied in [7] for spatial discretization and adaptive BDF schemes for time advancement. Simulations are extensively discussed to validate the model against realistic data and results available in the literature. In the final part of the lecture, we show that the proposed model can be used also for investigating the behavior of different classes of devices, e.g. the photoactivated light harvesting capacitor proposed in [3] or the sensing substrate used for neuronal photostimulation in [4].

REFERENCES

- [1] J.A. Barker, C.M. Ramsdale, and N.C. Greenham. Modeling the current-voltage characteristics of bilayer polymer photovoltaic devices. *Phys. Rev. B*, 67:075205, 2003.
- [2] C. de Falco, M. Porro, R. Sacco, and M. Verri. Multiscale modeling and simulation of organic solar cells. *Computer Methods in Applied Mechanics and Engineering*, 2012. To appear.
- [3] M. Garbugli, M. Porro, V. Roiati, A. Rizzo, G. Gigli, A. Petrozza, and G. Lanzani. Light energy harvesting with nano-dipoles. *Nanoscale*, 4:1728–1733, 2012.
- [4] D. Ghezzi, M.R. Antognazza, M. Dal Maschio, E. Lanzarini, F. Benfenati, and G. Lanzani. A hybrid bioorganic interface for neuronal photoactivation. *Nature Communications*, 2:166, 2011.
- [5] J.W. Jerome, M.A. Ratner, J.D. Servaites, C.-W. Shu, and S. Tan. Simulation of the Buxton-Clarke model for organic photovoltaic cells. In *Computational Electronics (IWCE), 2010 14th International Workshop on*, pages 1–4, Oct. 2010.
- [6] J. Williams and A.B. Walker. Two-dimensional simulations of bulk heterojunction solar cell characteristics. *Nanotechnology*, 19:424011, 2008.
- [7] J. Xu and L. Zikatanov. A monotone finite element scheme for convection-diffusion equations. *Mathematics of Computation*, 68(228):1429–1446, 1999.

Quantum transport in semiconductor nano-heterostructures

Paul N. Racec

Weierstrass Institute,
Mohrenstr. 39, 10117 Berlin, Germany

e-mail: Paul.Racec@wias-berlin.de

In this talk we review some mathematical challenges and our solutions arising from modeling of semiconductor nanodevices. Often, the active part of a nanodevice is a quantum system to which electrical contacts are attached such that there is an open quantum system on hand. The scattering theory is suitable for describing such systems. A practical method for computing the scattering matrix, even for 2D and 3D systems, is the R-matrix formalism. For this, one solves the Wigner-Eisenbud eigenvalue problem, i.e. the single-band effective mass Schrödinger equation on the scattering region with mixed Dirichlet/Neumann boundary conditions. Using the finite volume method one may take into account even a position dependent effective mass inside the heterostructure. The R-matrix formalism also yields the scattering wave functions inside the scattering region. Applications of this formalism to nanowire heterostructures will be provided.

As a consequence of opening the quantum system its discrete spectrum becomes a continuous spectrum with resonances. The practical computation of resonances as poles of the scattering matrix is also possible within R-matrix formalism. Furthermore, it allows for a systematic expansion in case of interacting resonances. We obtain a qualitative description of the peaks and dips in the linear conductance of a single electron transistor by Fano lines, where the complex asymmetry parameter is computed explicitly.

In order to compute the physical observables, one needs to know the density matrix of the system. The Landauer-Büttiker formalism provides an ansatz for the density matrix for open quantum systems. In the last years, a detailed analysis has been done in our group in order to mathematically justify the above ansatz.

Finally we discuss future challenges of a realistic modeling of semiconductor nanodevices, like many-particle description and dissipative interaction.

This is a joint work with Hans-Christoph Kaiser, Hagen Neidhardt and Roxana Racec.

Plane waveguides with corners in the small angle limit

Nicolas Raymond

IRMAR. University of Rennes 1,
Campus de Beaulieu, 35000 Rennes CEDEX, France
e-mail: `Nicolas.Raymond@univ-rennes1.fr`

The plane waveguides with corners considered here are infinite V-shaped strips with constant thickness. They are parametrized by their sole opening angle. We study the eigenpairs of the Dirichlet Laplacian in such domains when their angle tends to 0. We provide multi-scale asymptotics for eigenpairs associated with the lowest eigenvalues. For this, we investigate the eigenpairs of a one-dimensional model which can be viewed as their Born-Oppenheimer approximation. We also investigate the Dirichlet Laplacian on triangles with sharp angles. The eigenvalue asymptotics involve powers of the cube root of the angle, while the eigenvector asymptotics include simultaneously two scales in the triangular part, and one scale in the straight part of the guides.

Particle Detection with DEPFET Arrays in Gated Mode

Reiner Richter

Max-Planck-Institut Halbleiterlabor
Otto-Hahn-Ring 6, 81739 München, Germany

e-mail: rar@hll.mpg.de

Arrays of Depleted Field effect Transistors (DEPFET) are mostly operated in rolling shutter mode. During this CCD like operation scheme one pixel row is read out while the others are collecting signals. For the very compact pixel design used for the Belle II pixel detector we found a very simple procedure to switch off the signal collection for the entire matrix preserving the information already stored in the pixel. In this way we are able to mask out the noise which may occur during and after bunch injection in the accelerator. The talk will cover the device related aspects of this new operation mode. 3D-device-simulations are used to gain an insight into the device and to optimize the pixel design. The feasibility of this very promising approach is demonstrated by laser and tests beam data taken on a Belle II test matrix.

The DEPFET (DEPLETED Field Effect Transistor) is a semiconductor detector concept which combines detection and amplification within one device. A p-channel MOSFET is integrated onto a silicon detector substrate, which becomes fully depleted by a sufficiently high negative voltage applied to a p+ contact on the backside. A potential minimum is formed by sideward depletion, which is shifted directly underneath the transistor channel at a depth of about $1 \mu\text{m}$ by an additional phosphorus implantation underneath the external gate. Incident particles generate electron-hole pairs within the fully depleted bulk. While the holes drift into the back contact, electrons are accumulated in the potential minimum, called the internal Gate. If the transistor is switched on the electrons modulate the channel current. An amplification of more than 500pA per signal electron can be achieved with the technology currently available.

The removal of the signal charge and thermally generated electrons from the internal gate is called Clear. An n+ contact neighboring to the internal Gate is pulsed at a positive voltage providing a reach-through into the internal Gate. There is no any reset noise if the entire charge is completely removed. The probably most important feature of the DEPFET is the very small capacitance of the internal gate resulting in a low noise performance, even at room temperature and high operation speed. DEPFET arrays are usually operated in rolling shutter mode. Since only the addressed row is consuming power the rolling shutter mode is the most power efficient mode to operate a DEPFET array. In an accelerator particle bunches have to be refreshed by additional injections. Since new bunches cannot be injected in the same phase space volume they have to cool by synchrotron radiation which leads to noisy bunches. Due to rolling shutter scheme any junk charge introduced by noisy bunches which hits the detector spoils the entire detector array. According to pessimistic estimations more than 20% of the data could be affected due to this phenomena at Belle II. Noisy bunches may appear every $10 \mu\text{s}$ – the revolution time at Belle - in a well defined time window.

The pixel detector can be protected from noisy bunches if the DEPFET can be 'gated' into an insensitive (blind) mode during that time. Signals already collected in those pixel which were not yet readout have to be preserved during this blind phase. The basic idea of the 'gated mode' operation is to temporary provide an electrode which is much more attractive for electrons than the internal Gate. In the Belle DEPFET design the clear region is very large compared to Gate and the internal Gate beneath (see Fig. 1 and 2). The Figures illustrate the completely different charge collection behavior for $\text{Clear}_{\text{hi}} = 18\text{V}$ and $\text{Clear}_{\text{low}} = 3\text{V}$. During the noisy bunch phase the Clear lines of the entire matrix are raised at Clear_{hi} to protect the internal Gates from being hit. Related system aspects like high peak currents and cross talk are still under investigation.

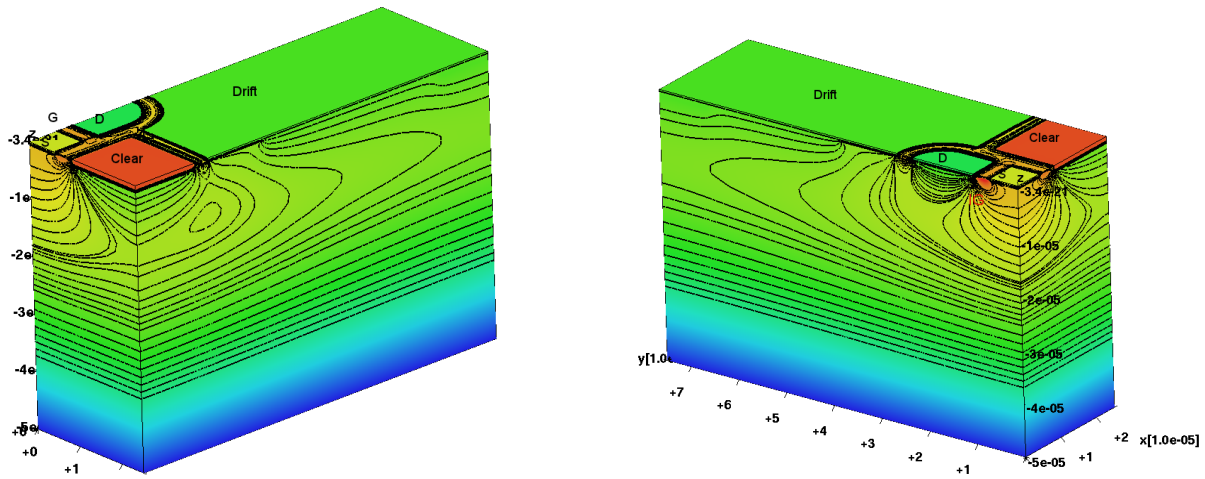


Fig. 1: 3D potential distributions of the 'collection state'. View on Clear and Drift regions (left) and on Drift, Drain (D), Source (S) and Gate (right). Clear is in off state (3V). Generated electrons drift into the internal Gate (IG). The simulation make use of Neumann boundary conditions. Every lateral boundary is a symmetric plane thus one half of a pixel is shown.

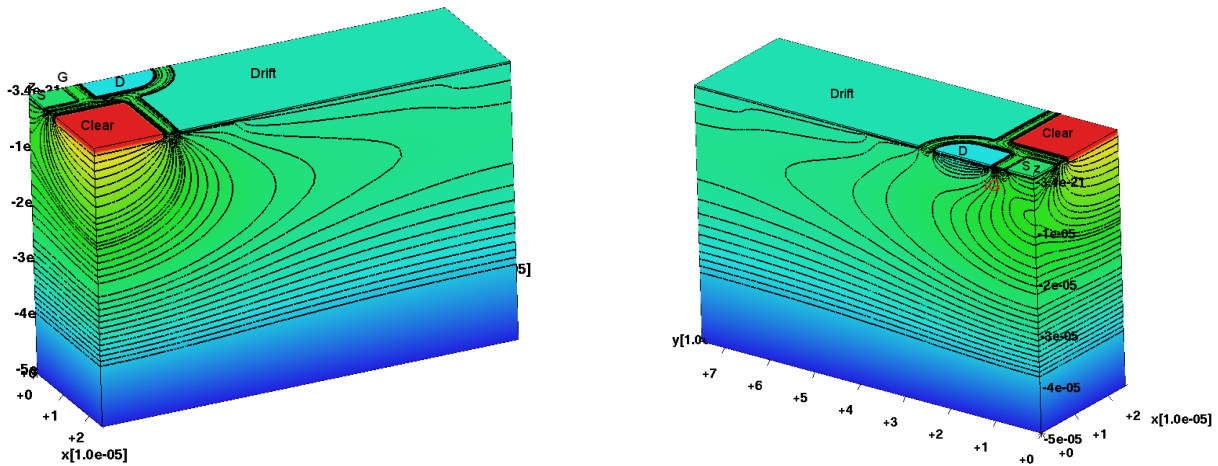


Fig. 2: 3D potential distributions of the 'blind state'. View on Clear and Drift regions (left) and on Drift, Drain (D), Source (S) and Gate (right). Clear is in on state (18V). Generated electrons drift into the Clear. The internal Gate (IG) is separated by potential barriers.

The selectivity to the actual *Clear* process is given by the external Gate. During *Clear* the external Gate has to be switched to the on state (current is flowing). This provides the most possible negative potential of the internal Gate and the biggest potential difference between the *Clear* region and the internal Gate. Electrons can be emitted over the potential barrier in-between. If the external Gate is switched off the hole channel at the surface disappears and the strong capacitive coupling between external and internal Gate leads to a much more positive potential in the latter one. Fig. 3 drafts the resulting potential slopes between internal Gate and *Clear*. Fig. 4. shows a measurement during blind mode. The removed charge is plotted versus the external Gate voltage. The *Clear* suppression works very well.

Abstracts

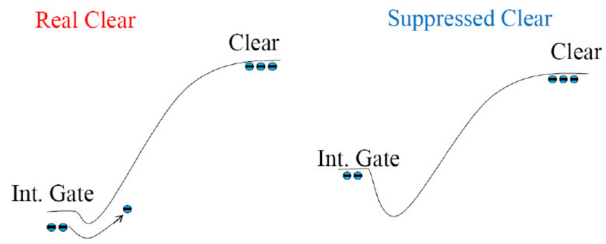


Fig. 3: Principle of the clear selectivity. Maximum potential between internal Gate and Clear during Real Clear and Suppressed Clear (blind mode).

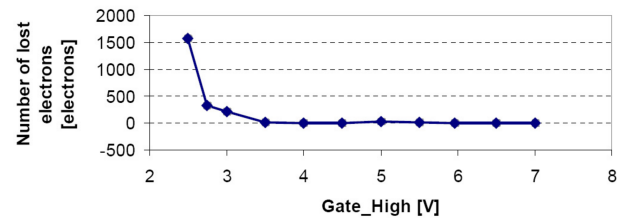


Fig. 4: Number of electrons removed from the internal Gate vs external Gate voltage. 12000 electrons were deposited before (laser illumination).

Physics and Modeling of Strain Effects in SiGe Heterojunction Bipolar Transistors

Holger Rücker and Bernd Heinemann

IHP GmbH, Innovations for High Performance Microelectronics,
Im Technologiepark 25, 15236 Frankfurt (Oder), Germany
e-mail: Ruecker@ihp-microelectronics.com

Strain engineering is a key technological method for performance enhancement of state-of-the-art CMOS technologies. The effect of strain on the semiconductor band structure is exploited to manipulate carrier transport properties and device characteristics in a beneficial way. Strain plays also an essential role for the electrical behavior of SiGe heterojunction bipolar transistors (HBT). The SiGe base layer experiences a biaxial compression due to its lattice mismatch to the Si substrate. Since this stress is close to the stability limit of the SiGe layer reliability concerns have discouraged the use of further stress enhancement techniques in SiGe HBTs so far. Nevertheless, the device fabrication process imposes significant stresses in addition to the lattice mismatch stress which can have significant impact on device characteristics. In this talk, we study the impact of process induced stress on the electrical characteristics of SiGe HBTs. Particularly large stresses are imposed by metal interconnects. We found experimentally that layout variations of aluminum metallization lines in the vicinity of a HBT can change collector currents by as much as 50% [1]. The main cause for this effect is uniaxial stress in the surface-normal direction imposed by the metal layer stack. The stress field and its effect on device characteristics can be calculated based on known mechanical properties and deformation potential theory.

Stress Simulation: The evolution of stress during the fabrication of a five-layer aluminum metallization was calculated taking into account the full thermal cycle of the process and measured intrinsic stress values of the individual layers. While the structured metal lines are essentially stress free at the deposition temperature (380°C) of the interlayer dielectrics they experience a strong tensile stress normal to the silicon surface when they are covered with oxide and cooled down to room temperature due to the large thermal expansion coefficient of aluminum. As a consequence, tensile stress up to 500 MPa is build up in transistor regions below a stack of 5 metal layers while adjacent Si regions without metal coverage are under compressive stress.

Device Simulation: Electrical characteristics of the strained HBTs were simulated within the drift-diffusion model. The effect of strain is taken into account by the splitting of electron and hole valleys according to deformation potential theory and k-p band structure. The mobility model includes modified inter-valley scattering due to strain splitting. The dominating effect of strain on the HBT characteristics is due to the strain dependence of the band gap in the SiGe base. Tensile strain in the surface-normal direction enhances the valley splitting introduced by lattice mismatch stress resulting in a reduced band gap in the base region of the HBT and in enhanced collector currents. The experimentally observed dependence of the collector current on the metal layout is well described by the presented simulation approach.

REFERENCES

- [1] H. Rücker, B. Heinemann, R. Barth, M. Lisker, IEEE IEDM Techn. Dig., pp. 34.5.1-4, 2011.

Numerical Simulation of Hetero Nanowire Tunnel Field Effect Transistors

Andreas Schenk⁽¹⁾, Reto Rhyner⁽¹⁾, Mathieu Luisier⁽¹⁾, and Cedric Bessire⁽²⁾

⁽¹⁾ Integrated Systems Laboratory, ETH Zürich,
Gloriastr. 35, 8092 Zürich, Switzerland

⁽²⁾ IBM Research-Zürich

e-mail: Schenk@iis.ee.ethz.ch

Tunnel Field Effect Transistors (TFETs) are candidates for low-power logic switches with steep (sub-thermal) slope which could enable a strongly reduced supply voltage. Device and circuit simulations of TFETs have shown a 8x smaller delay-energy product at a supply voltage of 0.35V compared to CMOS [1]. To improve the ON/OFF-current ratio, III-V/Si hetero junctions have been proposed [2]. Using nanowires has additional advantages: (i) the possibility of many different material combinations [3], (ii) efficient strain relaxation with shrinking diameter [3], (iii) a good electrostatic control due to the surrounding gate. Simulation results of this presentation are based on experimental data by Tomioka et al. [4, 5, 6] and Björk et al. [7] who advanced the integration of InAs nanowires on Si with nanometerscale hetero epitaxy.

We show how the combined application of a quantum transport solver and a TCAD tool can help to understand the behavior of InAs/Si hetero nanowire Esaki diodes and TFETs leading to design guidelines for the optimization of geometry, doping, gating, and biasing. We used the massively parallel, multi-dimensional, and atomistic quantum transport simulator OMEN [8] which is based on a $sp^3d^5s^*$ tight-binding representation of the band structure. Quantum transport simulation can be done either in the Non-equilibrium Greens Function (NEGF) formalism, if electron-phonon scattering must be included, or using the much faster Wave Function formalism in the ballistic case. OMEN has been applied to direct and phonon-assisted band-to-band tunneling (BTBT) in InAs, Si, and Ge nanowire homo TFETs [9]. The commercial device simulator Sentaurus-Device [10] offers various local and non-local BTBT models. However, neither a theory nor an analytical model for BTBT in a hetero junction between a direct and an indirect semiconductor exist. A practical workaround has to be used with Sentaurus-Device, since a tunnel path across the hetero interface must either belong to a direct (zero-phonon) or to a phonon-assisted tunnel process. Therefore, the “*dynamic nonlocal path BTBT model*” [11], calibrated for InAs, is also used on the silicon side, fitted to experimental data of [12], whereas the calibrated model for Si [13] is also used on the InAs side after proper modifications.

First, the BTBT current of short, unconfined Esaki homo diodes ($\langle 111 \rangle$, 20nm length, abrupt doping) was simulated with OMEN for different materials and doping levels. For the direct materials (InAs, GaSb) and for Ge, where coherent BTBT is dominant [9], the simulation of bulk-like diodes is straightforward. Bulk simulations are needed because of the absence of geometrical confinement in the fabricated nanowire TFETs (diameters in the range 25nm – 100nm). It turns out that InAs has the highest BTBT current density, followed by GaSb and Ge. The upper limit for InAs is $\sim 10\text{MA}/\text{cm}^2$. In the case of Si, due to the demanding electron-phonon coupling, at least one-dimensional confinement has to be applied. The estimated bulk limit of Si remains below $100\text{kA}/\text{cm}^2$, a factor 500 smaller than that of InAs. The doping concentrations at the InAs side of InAs/Si nanowire hetero Esaki diodes produced at IBM Research-Zurich [14, 15] are determined by reverse modeling. Measured InAs/Si nanowire TFET $I_D V_{GS}$ characteristics [6] are then compared with calibrated TCAD simulations. Whereas the measured $I_D V_{GS}$ curves show an almost constant slope over 2-3 orders of magnitude, a very weak ambipolarity, and a strong saturation of the ON-current for each source-drain voltage, simulation yields much higher ON-currents, a strong ambipolarity, curved slopes typical for BTBT, a minimum point slope of

45mV/dec, and no ON-current saturation. The most likely explanation of the measured currents is that they are dominated by defect-assisted tunneling (DAT), either during interface or bulk Shockley-Read-Hall generation. Although multi-phonon coupling parameters of the involved defects in InAs are not known, the shape of the $I_D V_{GS}$ curves can be qualitatively reproduced with a physics-based DAT model [16] in Sentaurus-Device. The apparent absence of BTBT in the measurement could be due to compressive biaxial strain in the highly lattice mismatched system.

REFERENCES

- [1] I. Young, Digest DRC, Penn State University, 3-6 (2012).
- [2] A. S. Verhulst et al., IEEE Electron Device Lett. 29, 1398 (2008).
- [3] M. T. Björk et al., Appl. Phys. Lett. 80, 1058 (2002).
- [4] K. Tomioka et al., Nano Lett. 8, 3475 (2008).
- [5] K. Tomioka et al., Nanotechnology 20, 145302 (2000).
- [6] K. Tomioka et al., Appl. Phys. Lett. 98, 083114 (2011).
- [7] M. T. Björk et al., Appl. Phys. Lett. 97, 163501 (2010).
- [8] M. Luisier et al., Phys. Rev. B 74, 205323 (2006).
- [9] M. Luisier et al., J. Appl. Phys. 107, 084507 (2010).
- [10] Synopsys Inc., Sentaurus-Device User Guide, Version 2011.09, Mountain View, California, (2011).
- [11] E. O. Kane, J. Phys. Chem. Solids 12, 181 (1959).
- [12] P. M. Solomon et al., J. Appl. Phys. 95 (10), 5800 (2004).
- [13] A. Schenk, Solid-State Electronics 36 (1), 19 (1993).
- [14] A. Schenk et al., Proc. SISPAD, 263 (2011).
- [15] C. D. Bessire et al., Nano Lett. 11 (10), 4195 (2011).
- [16] A. Schenk, Solid-State Electronics 35 (11), 1585 (1992).
- [17] C. G. Van de Walle, Phys. Rev. B 39 (3), 1871 (1989).

Modeling Spintronic Effects in Silicon

Dmitry Osintsev, Viktor Sverdlov, and Siegfried Selberherr

Institute for Microelectronics, TU Wien,
Gußhausstraße 27-29, A-1040 Wien, Austria

e-mail: Selberherr@iue.tuwien.ac.at

1. INTRODUCTION

The electron spin can change its orientation to opposite very fast by utilizing an amazingly small amount of energy, which offers a unique opportunity to reduce the power per operation in semiconductor logic devices. The spin field-effect transistor (SpinFET) is a switch which employs the spin properties of electrons. The SpinFET is composed of a semiconductor channel region sandwiched between the two ferromagnetic contacts, source and drain (Fig. 5). The source contact injects spin-polarized electrons in the semiconductor region. The gate-voltage-dependent spin-orbit interaction in the channel is used to modulate the current through the SpinFET [1]. It causes the electron spin to precess during the electron propagation through the channel. Only the electrons with their spin aligned to the drain contact magnetization can leave the channel through the drain contact, thus contributing to the current. The spin-orbit interaction is controlled electrically by applying an external gate voltage.

2. MODEL AND RESULTS

To calculate the transport properties of the ballistic spin field-effect transistor schematically shown in Fig. 5 we consider the Hamiltonian in the ferromagnetic regions in the following form [2], [3]:

$$(1) \quad \begin{aligned} H &= \frac{p_x^2}{2m_f^*} + h_0\sigma_z, & x < 0, \\ H &= \frac{p_x^2}{2m_f^*} \pm h_0\sigma_z, & x > L, \end{aligned}$$

where m_f^* is the effective mass in the contacts, $h_0 = 2PE_F/(P^2+1)$ is the exchange splitting energy with P defined as the spin polarization in the ferromagnetic regions, E_F is the Fermi energy, and σ_z is the Pauli matrix; \pm in (1) stands for the parallel (P) and anti-parallel (AP) configuration of the contact magnetization. The effective Hamiltonian in the channel is [2]

$$(2) \quad H = \frac{p_x^2}{2m_s^*} + \delta E_c + \frac{\beta}{\hbar}(p_x\sigma_x - p_y\sigma_y) + \frac{1}{2}g\mu_B B\sigma^*.$$

Here m_s^* is the subband effective mass, δE_c is the band mismatch between the ferromagnetic and the semiconductor region, g is the Landé factor, μ_B is the Bohr magneton, B is the magnetic field, and $\sigma^* \equiv \sigma_x \cos \gamma + \sigma_y \sin \gamma$ with γ defined as the angle between the magnetic field and the transport direction. As it was recently demonstrated [4], [5], the major contribution to the spin-orbit interaction in thin silicon films is due to the interface-induced inversion asymmetry and is mathematically equivalent to the linear Dresselhaus term. The coefficient β of the spin-orbit interaction in silicon heterostructures is a linear function of the gate voltage, which is the key for the current modulation in the channel (Fig. 5): it causes the electron spin to precess during the propagation through the channel. Fig. 6 shows the modulation of the tunnel magnetoresistance $TMR \equiv \frac{G^P - G^{AP}}{G^{AP}}$ defined by the channel conductances in P/AP configuration of the relative source/drain magnetizations on the strength of the spin-orbit interaction β in silicon channels of different length. As the length of the channel decreases, the period of the TMR modulation

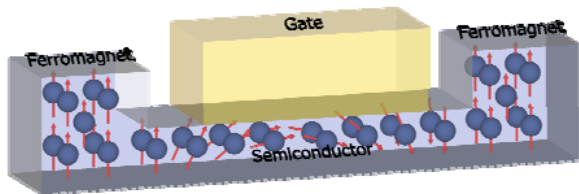


Fig. 5: Illustration of the Datta-Das spin field effect transistor. The spin-orbit interaction with the gate voltage dependent strength alters the spin polarization direction in the channel close to the ferromagnetic drain. Only electrons with the spin direction aligned to the drain magnetization contribute to the current.

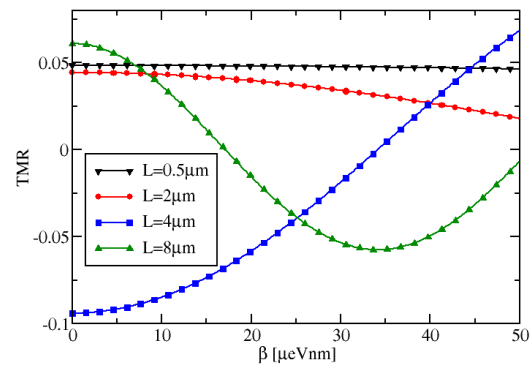


Fig. 6: TMR dependence on the value of the Dresselhaus spin-orbit interaction for $E_F = 2.47\text{eV}$, $\delta E_c = 2.154\text{eV}$, $z = 1$.

increases. $L = 4\mu\text{m}$ is sufficient to observe half of the oscillation and thus to modulate the TMR by adjusting the strength of the spin-orbit interaction. To facilitate the spin injection into silicon delta-function barriers of strength $z = 2m_F U / \hbar^2 k_F$ are introduced. As shown in Fig. 7, increasing the barrier leads to a more pronounced TMR modulation as a function of β . Fig. 8 shows oscillations of the TMR on the value of the bandgap mismatch δE_c . The period of the oscillations is inversely proportional to the length of the channel. Temperature exerts a significant influence on the oscillatory amplitude as shown in Fig. 9. For $L = 0.2\mu\text{m}$ the oscillatory behavior of the TMR completely vanishes for $T = 50\text{K}$. The reason is a relatively short period of the TMR oscillation with respect to δE_c shown in Fig. 8. For longer channels the period becomes even shorter, hence the shorter channels are needed to preserve the TMR modulation at higher temperatures as a function of δE_c . A different option to proceed to room temperature operation is to increase the channel length. Fig. 10 shows a possibility to modulate the TMR by changing the value of β even at room temperature, however, the channel length should be a few micrometers for the parameters specified [6]. The non-zero spin-orbit interaction leads to an increased spin relaxation. In quasi-one-dimensional electron structures, however, a suppression of the spin relaxation was predicted [7].

This work is supported by the European Research Council through the grant #247056 MOSIL-SPIN.

REFERENCES

- [1] S.Datta, B.Das, *Appl. Phys. Lett.* **56**, 665 (1990).
- [2] M.Cahay, S.Bandyopadhyay, *Phys. Rev. B* **69**, 045303 (2004).
- [3] K.Jiang et al., *IEEE Transactions on Electron Devices* **57**, 2005 (2010).
- [4] M.O.Nestoklon et al., *Phys. Rev. B* **77**, 155328 (2008).
- [5] M.Prada et al., *New J. Phys.* **13**, 013009 (2011).
- [6] D.Osintsev et al., *Solid-State Electronics* **71**, 25 (2012).
- [7] A.Bournel et al., *European Phys. J. Appl. Phys.* **4**, 1 (1998).

Abstracts

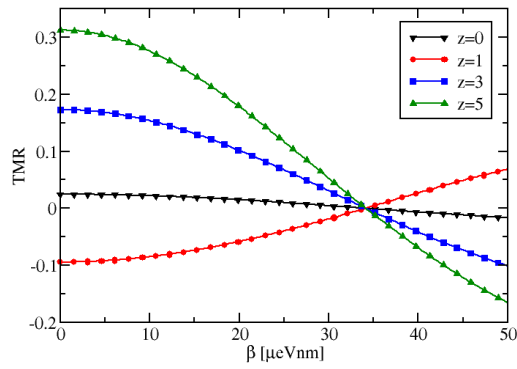


Fig. 7: TMR dependence on the spin-orbit interaction strength for $L = 4 \mu\text{m}$.

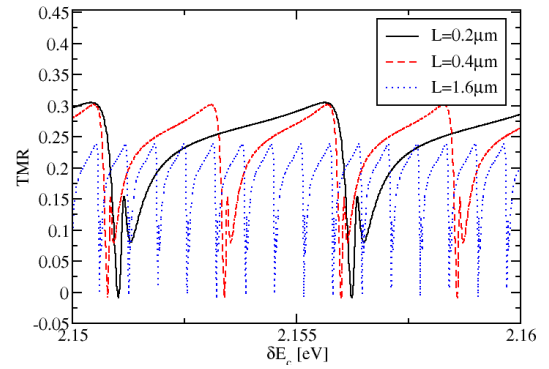


Fig. 8: TMR dependence on the conduction band mismatch for $z = 3$, $\beta = 42.3 \mu\text{eVnm}$.

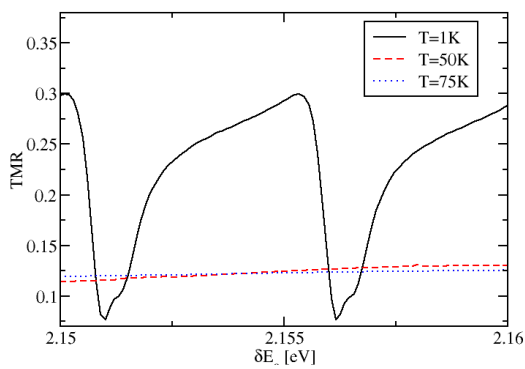


Fig. 9: TMR for different temperatures for the parameters from Fig. 8, and $L = 0.2 \mu\text{m}$, $V = 1 \text{meV}$.

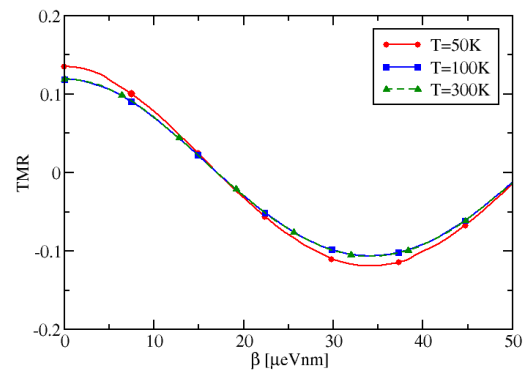


Fig. 10: TMR dependence on β for $L = 8 \mu\text{m}$. The other parameters are the same as in Fig. 8.

Virtual Experiments with High Power Semiconductor Devices Along and Beyond the Rim of the Safe Operating Area

Gerhard Wachutka

Munich University of Technology
Arcisstraße 21, 80333 Munich, Germany
e-mail: Wachutka@tep.ei.tum.de

Reliability and robustness against harsh operating conditions such as high temperature, hard unclamped switching, irradiation by cosmic particles, or mechanical vibrations are of utmost importance in the design of modern high power semiconductor devices.

Today's challenge is to make virtual experiments and tests on the computer, which are qualitatively correct and quantitatively accurate even for devices that have never been built before, and under operational conditions that very rarely occur as long as the device is kept within the so-called "safe operating area (SOA)". What we are interested in is to explore the rim of the SOA and even to go beyond it in order to study failure and, eventually, destruction mechanisms with a view to improving robustness and reliability of the devices. To this end, predictive high-fidelity computer simulations of "virtual destruction" have to be mastered.

While the physical models (typically energy-domain coupling continuous field models) which underlie the description of the inner life of electronic devices have become more and more elaborate and highly detailed, even under extreme and exceptional operating conditions, the capabilities and the quality of the simulation tools in which they are implemented still needs quite a lot of improvement. We will address these problems with reference to selected real-life examples, leading to a proclamation of improvements to be made towards a more judicious interpretation of numerical simulations w.r.t. accuracy, reliability and, hence, predictiveness.

Application of WIAS-TeSCA to the analysis of the power limits of semiconductor lasers

Hans Wenzel

Ferdinand-Braun-Institut, Leibniz-Institut für Höchstfrequenztechnik
Gustav-Kirchhoff-Straße 4, 12489 Berlin, Germany

e-mail: `Hans.Wenzel@fbh-berlin.de`

The factors that limit both the continuous wave and the pulsed output power of broad-area laser diodes driven at very high currents are investigated theoretically and experimentally. The decrease in the gain due to self-heating under CW operation and spectral holeburning under pulsed operation as well as heterobarrier carrier leakage and longitudinal spatial holeburning limit the maximum achievable output power.

Arnold, Anton	Vienna University of Technology, Institute of Analysis & Scientific Computing Wiedner Hauptstr. 8, 1040 Vienna, AUSTRIA anton.arnold@tuwien.ac.at
Beauchard, Karine	CNRS, CMLS & École Polytechnique, Centre de Mathématiques Laurent Schwartz Route de Saclay, 91120 Palaiseau, FRANCE karine.beauchard@math.polytechnique.fr
Bonato, Leo	Technische Universität Berlin, Institut für Festkörperphysik Hardenbergstr. 36, 10623 Berlin, GERMANY leo.bonato@tu-berlin.de
Coehoorn, Reinder	Philips Research Laboratories, Department of Photonic Materials and Devices High Tech Campus 4, Box WAG-12, 5656 AE Eindhoven, THE NETHERLANDS reinder.coehoorn@philips.com
Cruz, Valene Brenda	Karlsruhe Institute of Technology, Helmholtz Institute Ulm Ulm campus, 89069 Ulm, GERMANY valene.cruz@uni-ulm.de
Delebecque, Fanny	Université Paul Sabatier, Institut de Mathématiques de Toulouse, équipe MIP 118 route de Narbonne, 31062 Toulouse CEDEX 9, FRANCE fanny.delebecque@math.univ-toulouse.fr
Ehrhardt, Matthias	Bergische Universität Wuppertal, Fachbereich C - Mathematik und Naturwissenschaften, Lehrstuhl für Angewandte Mathematik und Numerische Analysis Gaußstrasse 20, 42119 Wuppertal, GERMANY ehrhardt@math.uni-wuppertal.de
Ervedoza, Sylvain	CNRS & Université Paul Sabatier, Institut de Mathématiques de Toulouse, équipe MIP 118 route de Narbonne, 31062 Toulouse CEDEX 9, FRANCE sylvain.ervedoza@math.univ-toulouse.fr
Fellner, Klemens	University of Graz, Institute for Mathematics and Scientific Computing Heinrichstr. 36, 8010 Graz, AUSTRIA klemens.fellner@uni-graz.at
Fiebach, André	Weierstrass Institute, Numerical Mathematics and Scientific Computing Mohrenstr. 39, 10117 Berlin, GERMANY andre.fiebach@wias-berlin.de
Fuhrmann, Jürgen	Weierstrass Institute, Numerical Mathematics and Scientific Computing Mohrenstr. 39, 10117 Berlin, GERMANY juergen.fuhrmann@wias-berlin.de
Gajewski, Herbert	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY herbert.gajewski@wias-berlin.de

Participants

Gärtner, Klaus	Weierstrass Institute, Numerical Mathematics and Scientific Computing Mohrenstr. 39, 10117 Berlin, GERMANY klaus.gaertner@wias-berlin.de
Glitzky, Annegret	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY annegret.glitzky@wias-berlin.de
Griepentrog, Jens A.	Weierstrass Institute, Modeling of Damage Processes Mohrenstr. 39, 10117 Berlin, GERMANY jens.griepentrog@wias-berlin.de
Gröger, Konrad	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY konrad.groeger@wias-berlin.de
Heinemann, Bernd	IHP GmbH, Innovations for High Performance Microelectronics Im Technologiepark 25, 15236 Frankfurt (Oder), GERMANY heinemann@ihp-microelectronics.com
Jerome, Joseph	Northwestern University 2033 Sheridan Road, IL 60208-2730 Evanston, USA jwj@math.northwestern.edu
Jourdana, Clément	Université Joseph Fourier - Grenoble 1 51 rue des Mathématiques, BP 53, 38041 Grenoble CEDEX 9, FRANCE clement.jourdana@imag.fr
Kaiser, Hans-Christoph	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY hans-christoph.kaiser@wias-berlin.de
Koprucki, Thomas	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY thomas.koprucki@wias-berlin.de
Korte, Lars	Helmholtz-Zentrum Berlin für Materialien und Energie GmbH, Institut für Silizium-Photovoltaik Kekuléstr. 5, 12489 Berlin, GERMANY korte@helmholtz-berlin.de
Kreisbeck, Carolin	Universidade Nova de Lisboa, Faculdade de Ciências e Tecnologia, Departa- mento de Matemática and Centro de Matemática e Aplicações Quinta da torre, 2829-516 Caparica, PORTUGAL Carolin.Kreisbeck@mathematik.uni-regensburg.de
Leendertz, Caspar	Helmholtz-Zentrum Berlin für Materialien und Energie GmbH, Institut für Silizium-Photovoltaik Kekuléstr. 5, 12489 Berlin, GERMANY caspar-leendertz@helmholtz-berlin.de

Liero, Matthias	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY matthias.liero@wias-berlin.de
Mantile, Andrea	Université de Reims, Laboratoire de Mathématiques U.F.R. Sciences Exactes et Naturelles, Moulin de la Housse - BP 1039, 51687 Reims CEDEX 2, FRANCE andrea.mantile@univ-reims.fr
Marquardt, Oliver	Paul-Drude-Institut für Festkörperelektronik Hausvogteiplatz 5-7, 10117 Berlin, GERMANY o.marquardt@yahoo.ie
Méhats, Florian	IRMAR, University of Rennes 1 Campus de Beaulieu, 35000 Rennes CEDEX, FRANCE florian.mehats@univ-rennes1.fr
Mielke, Alexander	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY alexander.mielke@wias-berlin.de
Neidhardt, Hagen	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY hagen.neidhardt@wias-berlin.de
Nier, Francis	IRMAR. University of Rennes 1 Campus de Beaulieu, 35000 Rennes CEDEX, FRANCE francis.nier@univ-rennes1.fr
Nürnberg, Reiner	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY reiner.nuernberg@wias-berlin.de
O'Reilly, Eoin	University College Cork, Tyndall National Institute, Photonics Theory Group 7-9 Dyke Parade, Cork, IRELAND eoin.oreilly@tyndall.ie
Paasch, Gernot	IFW Dresden (Leibniz-Institut für Festkörper- und Werkstoffforschung) PF 27 01 16, 01171 Dresden, GERMANY G.Paasch@ifw-dresden.de
Peikert, Vincent	ETH Zurich, Integrated Systems Lab Gloriastrasse 35, 8092 Zurich, SWITZERLAND vpeikert@iis.ee.ethz.ch
Pietra, Paola	Istituto di Matematica Applicata e Tecnologie Informatiche (IMATI) CNR via Ferrata 1, 27100 Pavia, ITALY pietra@imati.cnr.it
Pinaud, Olivier	Stanford University, Department of Mathematics & Université Lyon 1 Building 380, office 382V, CA 94305 Stanford, USA pinaud@math.univ-lyon1.fr

Participants

Platzek, Florian	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY florian.platzek@wias-berlin.de
Polizzi, Eric	University of Massachusetts Marcus 201, 100 Natural Resources Road, MA 01003 Amherst, USA polizzi@ecs.umass.edu
Porro, Matteo	Department of Mathematics "Francesco Brioschi", Politecnico di Milano, Center for Nano Science and Technology @PoliMi, Istituto Italiano di Tecnologia Piazza Leonardo da Vinci, 32, 20133 Milan, ITALY matteo.porro@mail.polimi.it
Racec, Paul	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY paul.racec@wias-berlin.de
Raymond, Nicolas	IRMAR. University of Rennes 1 Campus de Beaulieu, 35000 Rennes CEDEX, FRANCE nicolas.raymond@univ-rennes1.fr
Rehberg, Joachim	Weierstrass Institute, Partial Differential Equations Mohrenstr. 39, 10117 Berlin, GERMANY joachim.rehberg@wias-berlin.de
Richter, Rainer	Max-Planck-Institut Halbleiterlabor Otto-Hahn-Ring 6, 81739 Munich, GERMANY rar@hll.mpg.de
Rücker, Holger	IHP GmbH, Innovations for High Performance Microelectronics Im Technologiepark 25, 15236 Frankfurt (Oder), GERMANY ruecker@ihp-microelectronics.com
Schenk, Andreas	ETH Zurich, Integrated Systems Lab Gloriastrasse 35, 8092 Zurich, SWITZERLAND schenk@iis.ee.ethz.ch
Selberherr, Siegfried	Technische Universität Wien, Institut für Mikroelektronik Gusshausstrasse 27-29/360, A-1040 Vienna, AUSTRIA Selberherr@TUWien.ac.at
Stephan, Holger	Weierstrass Institute, Numerical Mathematics and Scientific Computing Mohrenstr. 39, 10117 Berlin, GERMANY holger.stephan@wias-berlin.de
Teodoreanu, Ana-Maria	Technische Universität Berlin, Fachgebiet Halbleiterbauelemente Einsteinufer 19, 10587 Berlin, GERMANY ana-maria.teodoreanu@tu-berlin.de
Wachutka, Gerhard	TU München, Lehrstuhl für Technische Elektrophysik Arcisstraße 21, 80333 Munich, GERMANY wachutka@tep.ei.tum.de

Wenzel, Hans	Ferdinand-Braun-Institut, Leibniz-Institut für Höchstfrequenztechnik Gustav-Kirchhoff-Straße 4, 12489 Berlin, GERMANY hans.wenzel@fbh-berlin.de
Wölz, Martin	Paul-Drude-Institut für Festkörperelektronik Hausvogteiplatz 5-7, 10117 Berlin, GERMANY martin.woelz@pdi-berlin.de
Wünsche, Hans-Jürgen	Humboldt-Universität zu Berlin, Institut für Physik Newtonstraße 15, 12489 Berlin, GERMANY ede@physik.hu-berlin.de

