



**Weierstrass Institute for
Applied Analysis and Stochastics**



Eigenvalue order statistics and mass concentration in the parabolic Anderson model

Based on joint works with Marek Biskup (UC Los Angeles) and Renato dos Santos (NYU Shanghai)

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Definition

Let Δ be the standard Laplace operator on \mathbb{Z}^d , and let $\xi = (\xi(z))_{z \in \mathbb{Z}^d}$ be a random potential, then $\Delta + \xi$ is called a *random Schrödinger operator*.

Explicitly, $(\Delta + \xi)f(z) = \Delta f(z) + \xi(z)f(z)$ for L^2 -functions f .

Great mathematical interest in the eigenfunctions stems from the famous

Prediction by P.W. ANDERSON (1958):

Anderson localisation: In a large part of the spectrum of $\Delta + \xi$, all values should be eigenvalues, and the corresponding eigenfunction should be exponentially localised around some (randomly distributed) site.

- Anderson localisation has been confirmed for many random potentials ξ for spectral values close to the boundary of the spectrum, or for $\Delta + \beta\xi$ if $|\beta|$ is large enough.
- Two proof methods (1990s, early 2000s): *Fractional moment method* and *Multiscale analysis*.

We are interested in the **upper edge** of the spectrum of $\Delta + \xi$ in a large box B , i.e., in the principal part, including the **principal eigenvalue** (with zero Dirichlet boundary condition),

$$\begin{aligned}\lambda_1(B) &= \sup \{ \langle g, (\Delta + \xi)g \rangle : g \in \ell^2(\mathbb{Z}^d), \text{supp}(g) \subset B, \|g\|_2 = 1 \} \\ &= - \inf \left\{ \|\nabla g\|_2^2 - \sum_z \xi(z)g^2(z) : g \in \ell^2(\mathbb{Z}^d), \text{supp}(g) \subset B, \|g\|_2 = 1 \right\}.\end{aligned}$$

Introduce all the eigenvalues, $\lambda_1(B) > \lambda_2(B) \geq \lambda_3(B) \geq \dots \geq \lambda_{|B|}(B)$.

Our questions:

- What is the upper-tail behaviour of $\lambda_1(B)$, in particular when coupled with $|B| \rightarrow \infty$?
- Is there an extreme-value order statistics for the top eigenvalues in this limit?
- What is the domain of attraction, what are the scaling parameters?
- Does the point process $\sum_{k=1}^{|B|} \delta_{\lambda_k(B)}$ converge, after normalisation?
- Are the corresponding eigenfunctions exponentially localised? If yes, where?

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Recall: If $M_N = \max\{X_1, \dots, X_N\}$ is the maximum of N i.i.d. random variables, then, if some a_N, b_N exist such that $(M_N - a_N)/b_N$ converges towards a non-degenerate variable in law, then this is either Gumbel, or Fréchet or Weibull. Also $\sum_{k=1}^N \delta_{(X_k - a_N)/b_N}$ converges.

Heat equation with random potential; parabolic Anderson model (PAM):

$$\frac{\partial}{\partial t} u(t, z) = \Delta u(t, z) + \xi(z)u(t, z), \quad \text{for } (t, z) \in (0, \infty) \times \mathbb{Z}^d, \quad (1)$$

$$u(0, z) = \delta_0(z), \quad \text{for } z \in \mathbb{Z}^d. \quad (2)$$

Interpretations / Motivations:

- **Random mass transport** through a **random field** of sinks and sources.
- Expected particle number in a **branching random walk** model in a field of **random branching and killing rates**.

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Eigenvalue expansion

$$u(t, z) \sim u_{B_t}(t, z) = \sum_k e^{t\lambda_k(B_t)} \varphi_k(0) \varphi_k(z),$$

where $\varphi_1, \varphi_2, \varphi_3 \dots$ are the corresponding orthonormal eigenfunctions in

$$B_t = t \times \left[-\frac{1}{2}, \frac{1}{2}\right]^d.$$

- In the limit $t \rightarrow \infty$, it is by far not automatic that only $\lambda_1(B_t)$ survives. Rather, the maximum of $e^{t\lambda_k(B_t)} \varphi_k(0)$ over k will be decisive. Therefore, we must know the joint behaviour of *all* top eigenvalues.

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- All the eigenfunctions φ_k decay exponentially fast away from some random site x_k . Then the distance $|x_k - 0|$ determines the value of $\varphi_k(0) \approx e^{-c|x_k|}$.
- The random sites x_k form (after rescaling) a homogeneous Poisson point process, and the eigenvalues $\lambda_k(B_t)$ behave like i.i.d.

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- Earlier work on the asymptotics of the expectation of $e^{t\lambda(B_t)}$ showed [GÄRTNER, K., MOLCHANOV 2007 that the typical shape of the eigenfunction approaches a deterministic shape, which is given as the maximizer of $\varphi \mapsto \lambda(\varphi) - I(\varphi)$, where $I(\varphi)$ describes the rate function for the probability that the potential looks like φ on a large scale.

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- All the details of the above heavily depend on the upper-tail behaviour of $\xi(0)$, i.e., on the asymptotics of $\text{Prob}(\xi(0) > r)$ as $r \rightarrow \text{esssup}\xi(0)$.

- Surveys: [K. 16] and [ASTRAUSKAS 16]
- Eigenvalue order statistics and Poisson point process convergence at spectral top (single-site eigenfunctions): [ASTRAUSKAS 08, 12, 13]
- Eigenvalue order statistics and Poisson point process convergence at spectral top (non-degenerate eigenfunctions, double-exp. dist) [BISKUP, K. 16], see below
- Poisson point process convergence in Anderson localisation regime: [MOLCHANOV 81], [MINAMI 96], [KILIP, NAKANO 07], [GERMINET, KLOPP 13, 14]
- Concentration of the PAM in one single site: [K., LACON, MÖRTERS, SIDOROVA 09], [LACON, MÖRTERS 12], [FIDOROV, MUIRHEAD 14], [SIDOROVA, TWAROWSKI 16]
- Concentration of the PAM in one non-degenerate island: [BISKUP, K., DOS SANTOS 18], see below, related partial result [DING, XU 18] for bounded potential
- first steps for white noise in $d = 2$ by [ALLEZ/CHOUK], [CHOUK/VAN ZUIJLEN (2019)] and [PERKOWSKI/K./VAN ZUIJLEN (2019+)].

Open:

Eigenvalue order statistics, Poisson point process convergence at spectral top for bounded potentials, and concentration of the PAM in one island for other potential distributions, e.g. Gaussian fields in \mathbb{R}^d (smooth or white noise).

We are working here for ξ **double-exponentially distributed**, i.e., for some $\varrho \in (0, \infty)$,

$$\text{Prob}(\xi(0) > r) = \exp \left\{ -e^{r/\varrho} \right\}, \quad r \in \mathbb{R}.$$

Theorem 1 [BISKUP/K., CMP 16]

There is a number $\chi = \chi_\varrho \in (0, 2d)$ and a sequence $(a_L)_{L \in \mathbb{N}}$ with $a_L = \varrho \log \log |B_L| - \chi + o(1)$ as $L \rightarrow \infty$ and, for any $L \in \mathbb{N}$, a sequence $(X_k^{(L)})_k$ in B_L such that, in probability,

$$\lim_{L \rightarrow \infty} \sum_{z: |z - X_k^{(L)}| \leq \log L} \varphi_k(z)^2 = 1, \quad k \in \mathbb{N},$$

and the law of

$$\sum_{k \in \mathbb{N}} \delta_{\left(\frac{1}{L} X_k^{(L)}, (\lambda_k(B_L) - a_L) \log L\right)}$$

converges weakly to a Poisson process on $B_1 \times \mathbb{R}$ with intensity measure $dx \otimes e^{-\lambda} d\lambda$.

- Hence, the **top eigenvalues** in B_L are of **order** $\log \log L$ and leave **gaps of order** $1/\log L$ (rather than $1/|B_L|$ as in the bulk of the spectrum).
- The **localisation centres** are separated by $\asymp L$ and are **homogeneously distributed**.
- We have an assertion reminding on **Anderson localisation at the edge** of the spectrum.

Theorem 2 [BISKUP, K., DOS SANTOS, PTRF 18]

Put $r_L = L \log L \log \log L$ and

$$\Psi_{L,t}(z, \lambda) = \frac{t}{r_L} (\lambda - a_L) \log L - \frac{|z|}{L},$$

and pick k such that $\Psi_{L,t}(X_k^{(L)}, \lambda_k(B_L))$ is maximal. Put $Z_t = X_k^{(L)}$.

Then, with L_t defined by $r_{L_t} = t$, for any $R_t \gg \log t$,

$$\lim_{R \rightarrow \infty} \lim_{t \rightarrow \infty} \frac{1}{U(t)} \sum_{z: |z - Z_t| \leq R} u(t, z) = 1 \quad \text{in probability.}$$

- Hence, the total mass essentially comes from a single $\gg \log t$ -island in the centred box with radius $L_t \asymp t / (\log t \log \log t)$.
- The [Poisson process convergence](#) holds only in the box B_{L_t} , and the contribution from $B_{t \log^2 t}^c$ is easily seen to be negligible. The intermediate region is delicate.
- The two terms in $\Psi_{L,t}$ come from the [eigenvalue and the probabilistic cost](#) for the random walk in the Feynman-Kac formula. The choice of r_L comes from an optimisation of $\mathbb{P}_0(|X_s| \asymp L) e^{(t-s)\lambda_1} \approx e^{-L \log(L/s)} e^{(t-s)a_L}$ over $s \in [0, t]$.

A system is said to **age** if its significant changes come after longer and longer (or shorter and shorter) time lags.

Hence, one can see from the **frequency of changes** how much time has elapsed.

Ageing properties of the PAM can now be studied in terms of the time lags between **jumps of the concentration site**.

These ones, in turn, may be described as follows.

Theorem 3: Scaling limit of concentration location [BKDS 18]

As $t \rightarrow \infty$, the process $(Z_{\theta t}/L_t)_{\theta \in [1, \infty)}$ converges in distribution to a process $(\bar{Z}(\theta))_{\theta \in [1, \infty)}$, whose marginals $\bar{Z}(\theta)$ have d independent components, which are centered and Laplace-distributed (i.e., with density $z \mapsto e^{-|z|/\theta}$).

Furthermore, $(Z_t)_{t \in [0, \infty)}$ is aging in the sense that, for any $s > 0$,

$$\lim_{t \rightarrow \infty} \text{Prob}\left(Z_t = Z_{t+\theta t} \text{ for every } \theta \in [0, s]\right)$$

exists and is a non-trivial function of s .

$$\text{Put } \varepsilon_R = 2d \left(1 + \frac{A}{2d}\right)^{1-2R}.$$

The top eigenvalues in $B = B_L$ remain the top eigenvalues after discarding potential values significantly less than the eigenvalues.

Fix $A > 0$ and $R \in \mathbb{N}$ and put $U = \bigcup_{z \in B: \xi(z) \geq \lambda_1(B) - 2A} B_R(z)$. Then

$$\lambda_k(B) \geq \lambda_1(B) - A/2 \quad \implies \quad |\lambda_k(B) - \lambda_k(U)| \leq \varepsilon_R.$$

- The corresponding ℓ^2 -normalized eigenvector $\varphi = \varphi_k$ decays rapidly away from U .
- Proof uses the martingale $(\varphi(Y_n) \prod_{l=0}^{n-1} \frac{2d}{2d + \lambda - \xi(Y_l)})_{n \in \mathbb{N}}$ (with $(Y_n)_n$ an SRW).

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- Furthermore, we use that $\partial_{\xi(z)} \lambda_k(B) = \varphi(z)^2$.
- Introduce $\xi_s = \xi - s \mathbb{1}_{B \setminus U}$ for $s \in [0, \infty]$. Then

$$|\partial_s \lambda_k(\xi_s, B)| = \sum_{z \in B \setminus U} \varphi_{k, \xi_s}(z)^2,$$

which is very small. Integrating over $s \in [0, \infty]$ gives the estimate.

The top eigenvalues are the principal eigenvalues in local regions, and the corresponding eigenfunctions are exponentially localised.

A bit more precisely, with the help of the variational characterisation of the asymptotics of the PAM [GÄRTNER/K./MOLCHANOV 07], one proves the following.

- U consists of connected components of bounded size, which are far away from each other.

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- U consists of connected components of bounded size, which are far away from each other.
- For any component C , if $\lambda_1(C)$ is close to $a_L \approx \rho \log \log L$, then $\lambda_1(C)$ is bounded away from $\lambda_2(C)$.

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- If λ is an eigenvalue of $\Delta + \xi$ larger than $\lambda_1(B_L) - A/2$ and φ a corresponding ℓ^2 -normalised eigenfunction such that the distance of λ to the nearest eigenvalue (spectral gap) is larger than $3\varepsilon_R$, then φ decays exponentially away from one of the components of U .

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- The proof uses that the path $[0, \infty] \ni s \mapsto \lambda_k(\xi_s, B_L)$ (with $\xi_s = \xi - s\mathbb{1}_{B_L \setminus U}$) does not cross other eigenvalues and therefore admits a continuous choice of corresponding eigenfunctions. The one for $s = \infty$ puts all its mass in one component, and the one for $s = 0$ is uniformly close.

The scale a_L satisfies $\text{Prob}(\lambda_1(B_R) > a_L) = 1/|B_L|$.

Hence we may expect finitely many sites in B_L where the local eigenvalue is $\approx a_L$.

$\lambda_1(B_R)$ lies in the max-domain of a Gumbel random variable

As $L \rightarrow \infty$, for any $s \in \mathbb{R}$,

$$\text{Prob}(\lambda_1(B_R) > a_L + s/\log L) = e^{-s} \frac{1}{|B_L|} (1 + o(1)).$$

- The event $\{\lambda_1(B_R) > a\}$ is more or less the same as the event that some shift of the potential $\xi(\cdot)$ is larger than $a + \chi + \psi(\cdot)$ for some well-chosen function ψ .

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- Shifting ξ by an amount of $s/\log L$ yields an additional factor of e^{-s} , using properties of ψ and of the distribution of ξ and some information from the variational characterisation.