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**Energy estimates for continuous and discretized
reaction-diffusion systems in heterostructures**

Outline

- ▷ Model equations
- ▷ Weak formulation and energy functionals
- ▷ Energy estimates for the continuous problem
- ▷ Results for space and time discrete problem
- ▷ Concluding remarks

The model

X_i	species, $i = 1, \dots, m$	\bar{u}_i	reference densities
v_i	chemical potentials	$u_i = \bar{u}_i e^{v_i}$	densities

- reversible reactions of mass action type

$$\alpha_1 X_1 + \dots + \alpha_m X_m \rightleftharpoons \beta_1 X_1 + \dots + \beta_m X_m, \quad (\alpha, \beta) \in \mathcal{R} \subset \mathbb{Z}_+^m \times \mathbb{Z}_+^m$$

net rate $k_{\alpha\beta}(\mathbf{e}^{v \cdot \alpha} - \mathbf{e}^{v \cdot \beta}), \quad v = (v_1, \dots, v_m)$

net production rate of species X_i $R_i := \sum_{(\alpha, \beta) \in \mathcal{R}} k_{\alpha\beta}(\mathbf{e}^{v \cdot \alpha} - \mathbf{e}^{v \cdot \beta})(\beta_i - \alpha_i)$

- mass fluxes $j_i = -u_i \mathbf{S}_i(\cdot) \nabla v_i, \quad i = 1, \dots, m,$
- anisotropies $\mathbf{S}_i(x) = Q_i^T(x) \text{diag}(\mu_i^k(x)) Q_i(x)$

continuity equations

$$\frac{\partial u_i}{\partial t} + \nabla \cdot j_i = R_i \text{ in } \mathbb{R}_+ \times \Omega, \quad \nu \cdot j_i = 0 \text{ on } \mathbb{R}_+ \times \Gamma,$$

$$u_i(0) = U_i \text{ in } \Omega, \quad i = 1, \dots, m.$$

Assumptions (A)

$\Omega \subset \mathbb{R}^N$ bounded Lipschitzian domain, $N \leq 3$, $\Gamma = \partial\Omega$;

$\mu_i^k \in L_+^\infty(\Omega)$, $\text{ess inf}_\Omega \mu_i^k \geq \delta$, $k = 1, \dots, N$, $i = 1, \dots, m$;

$\bar{u}_i \in L_+^\infty(\Omega)$, $\bar{u}_i \geq \delta$, $U_i \in L_+^\infty(\Omega)$, $q_i \in \mathbb{Z}$, $i = 1, \dots, m$;

$\mathcal{R} \subset \mathbb{Z}_+^m \times \mathbb{Z}_+^m$ finite subset, $k_{\alpha\beta} \in L_+^\infty(\Omega)$, $\int_\Omega k_{\alpha\beta} \, dx > 0$ for $(\alpha, \beta) \in \mathcal{R}$,

$\max \left\{ \sum_{i=1}^m \alpha_i, \sum_{i=1}^m \beta_i \right\} \leq 3 \quad \forall (\alpha, \beta) \in \mathcal{R} \quad \text{if } N > 2$,

there are no “false” equilibria in the sense of Prigogine & Defay '54,

$\sum_{i=1}^m \int_\Omega U_i \kappa_i \, dx > 0 \quad \forall \kappa \in \mathcal{S}^\perp$, $\kappa \geq 0$, $\kappa \neq 0$.

stoichiometric subspace

$$\mathcal{S} := \text{span}\{\alpha - \beta : (\alpha, \beta) \in \mathcal{R}\}$$

Weak formulation

$$v = (v_1, \dots, v_m) \in V = H^1(\Omega; \mathbb{R}^m), \quad u = (u_1, \dots, u_m) \in V^*$$

operators $A : V \cap L^\infty(\Omega, \mathbb{R}^m) \rightarrow V^*, E = (E_1, \dots, E_m) : V \rightarrow V^*$

$$\langle Av, \bar{v} \rangle_V := \int_{\Omega} \left\{ \sum_{i=1}^m \bar{u}_i \mathbf{e}^{v_i} \mathbf{S}_i \nabla v_i \cdot \nabla \bar{v}_i + \sum_{(\alpha, \beta) \in \mathcal{R}} k_{\alpha\beta} \left(\mathbf{e}^{\alpha \cdot v} - \mathbf{e}^{\beta \cdot v} \right) (\alpha - \beta) \cdot \bar{v} \right\} dx$$

$$\langle Ev, \bar{v} \rangle_V := \int_{\Omega} \sum_{i=1}^m \bar{u}_i \mathbf{e}^{v_i} \bar{v}_i dx, \quad \bar{v} \in V$$

Problem (P)

$$\left. \begin{aligned} u'(t) + Av(t) &= 0, \quad u(t) = Ev(t) \text{ f.a.a. } t \in \mathbb{R}_+, \quad u(0) = U, \\ u &\in H_{\text{loc}}^1(\mathbb{R}_+; V^*), \quad v \in L_{\text{loc}}^2(\mathbb{R}_+; V) \cap L_{\text{loc}}^\infty(\mathbb{R}_+; L^\infty(\Omega)^m). \end{aligned} \right\} \quad (\text{P})$$

Energy functionals

dissipation rate

$$D(v) = \int_{\Omega} \sum_{i=1}^m \bar{u}_i \mathbf{e}^{v_i} \mathbf{S}_i \nabla v_i \cdot \nabla v_i \, dx$$

$$+ \int_{\Omega} \sum_{(\alpha, \beta) \in \mathcal{R}} k_{\alpha\beta} (e^{v \cdot \alpha} - e^{v \cdot \beta}) (\alpha - \beta) \cdot v \, dx \geq 0 \quad \forall v \in V \cap L^{\infty}(\Omega, \mathbb{R}^m)$$

free energy

$$F(u) = \int_{\Omega} \sum_{i=1}^m \left(u_i \left(\ln \frac{u_i}{\bar{u}_i} - 1 \right) + \bar{u}_i \right) dx$$

Results for the continuous problem (G./Hünlich '97, G./Gärtner '07)

- Invariants:** (u, v) solution to (P) \implies

$$u(t) - U \in \mathcal{U} := \left\{ u \in V^* : (\langle u_1, 1 \rangle, \dots, \langle u_m, 1 \rangle) \in \mathcal{S} \right\} \quad \forall t > 0.$$

- Thermodynamic equilibrium:** There exists a unique solution (u^*, v^*) to

$$Av^* = 0, \quad u^* = Ev^*, \quad u^* - U \in \mathcal{U}. \quad (\text{S})$$

It holds $\nabla v^* = 0$ and $v^* \in \mathcal{S}^\perp$.

- Monotone decay of the free energy:** Let (u, v) be a solution to Problem (P). Then

$$F(u(t_2)) \leq F(u(t_1)) \leq F(U) \quad \text{for } t_2 \geq t_1 \geq 0.$$

- Exponential decay of the free energy:** Let (u, v) be a solution to Problem (P), and let (u^*, v^*) be the thermodynamic equilibrium. Then there exists a constant $\lambda > 0$ such that

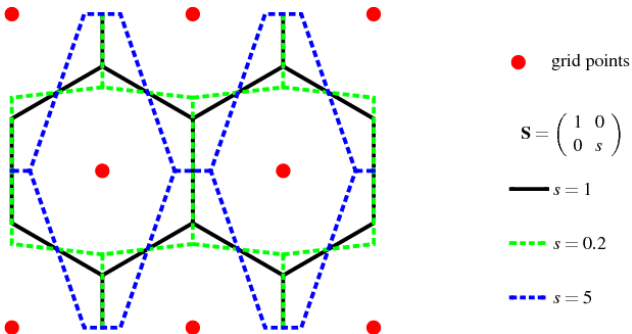
$$F(u(t)) - F(u^*) \leq e^{-\lambda t} (F(U) - F(u^*)) \quad \forall t \geq 0.$$

Space and time discretized problems

fixed set of grid points $x^k, k \in K$, for each species anisotropic Voronoi boxes

$$V_i^k = \{x \in \bar{\Omega} : d_i(x, x^k) \leq d_i(x, x^l) \quad \forall l \in K\}, \quad i = 1, \dots, m, \quad k \in K$$

$$d_i(x, y)^2 := (x - y)^T \mathbf{S}_i^{-1} (x - y)$$



Space and time discretized problems

(B) $\bar{u}_i = \text{const}$, $i = 1, \dots, m$, $k_{\alpha\beta} = \text{const}$ $(\alpha, \beta) \in \mathcal{R}$, $\tau = \text{const}$,
 S_i constant, symmetric, positive definite matrices.

- u_i^k mass of species X_i in V_i^k , $U_i^k = \int_{V_i^k} U_i \, dx$, $k \in K, i = 1, \dots, m$
- potentials v_i^k associated to grid points x^k
- discrete state equations $u_i^k = \bar{u}_i e^{v_i^k} |V_i^k|$, $k \in K, i = 1, \dots, m$
- notation $\vec{u} = (u_i^k)_{k \in K, i=1, \dots, m}$, $\vec{v} = (v_i^k)_{k \in K, i=1, \dots, m}$

(C) $\mathcal{Z} = \{0, t_1, \dots, t_n, \dots\}$ partition of \mathbb{R}_+ , $t_n \in \mathbb{R}_+$, $t_{n-1} < t_n$, $n \in \mathbb{N}$,
 $t_n \rightarrow \infty$ as $n \rightarrow \infty$, $h_n = t_n - t_{n-1}$, $\sup_{n \in \mathbb{N}} h_n < \infty$.

Space and time discretized problems

Discretization scheme

$$\left. \begin{aligned} \frac{u_i^k(t_n) - u_i^k(t_{n-1})}{h_n} &= - \sum_{l \in K} J_i^{kl}(t_n) |\partial V_i^k \cap \partial V_i^l| + R_i^k(t_n), \\ k \in K, n \geq 1, i &= 1, \dots, m, \\ \vec{u}(0) &= \vec{U}. \end{aligned} \right\} \quad (\text{PD})$$

discretized fluxes

$$J_i^{kl} = -\bar{u}_i Z_i^{kl} \frac{v_i^l - v_i^k}{|x^l - x^k|} |\mathbf{S}_i \nu_i^{kl}|, \quad \nu_i^{kl} \text{ outer unit normal of } V_i^k, \quad Z_i^{kl} = \frac{1}{2}(\mathbf{e}^{v_i^k} + \mathbf{e}^{v_i^l})$$

source terms from reactions

$$R_i^k = \sum_{\alpha, \beta \in \mathcal{R}} (\beta_i - \alpha_i) \sum_{k_1 \in K} \cdots \sum_{k_{i-1} \in K} \sum_{k_{i+1} \in K} \cdots \sum_{k_m \in K} R_{\alpha\beta} [v_1^{k_1}, \dots, v_{i-1}^{k_{i-1}}, v_i^k, v_{i+1}^{k_{i+1}}, \dots, v_m^{k_m}] \\ \times |V_1^{k_1} \cap \cdots \cap V_{i-1}^{k_{i-1}} \cap V_i^k \cap V_{i+1}^{k_{i+1}} \cap \cdots \cap V_m^{k_m}|,$$

$$R_{\alpha\beta} [v_1^{k_1}, \dots, v_m^{k_m}] = k_{\alpha\beta} (\mathbf{e}^{\sum_{i=1}^m \alpha_i v_i^{k_i}} - \mathbf{e}^{\sum_{i=1}^m \beta_i v_i^{k_i}})$$

Discrete energy functionals

discrete version of E $\hat{E}: \mathbb{R}^{Mm} \rightarrow \mathbb{R}^{Mm}, \quad (M = \#K)$

$$\hat{E}\vec{v} = \left((\bar{u}_i \mathbf{e}^{v_i^k} |V_i^k|)_{k \in K} \right)_{i=1, \dots, m}$$

discrete free energy

$$\hat{F}(\vec{u}) = \sum_{i=1}^m \sum_{k \in K} \left(u_i^k v_i^k - u_i^k + \bar{u}_i |V_i^k| \right) \quad \text{for } \vec{u} = \hat{E}\vec{v}$$

invariants: (\vec{u}, \vec{v}) solution to the discretized problem (PD) \implies

$$\vec{u}(t_n) - \vec{U} \in \hat{\mathcal{U}} = \left\{ \vec{u} \in \mathbb{R}^{Mm} : \left(\sum_{k \in K} u_1^k, \dots, \sum_{k \in K} u_m^k \right) \in \mathcal{S} \right\} \quad \forall n \in \mathbb{N}$$

Steady states of the discretized problem

Theorem 1. [Thermodynamic equilibrium]

We assume (A) and (B). Then there is a unique solution (\vec{u}^*, \vec{v}^*) to Problem

$$\left. \begin{aligned} \sum_{l \in K} J_i^{kl} |\partial V_i^k \cap \partial V_i^l| - R_i^k &= 0, \quad k \in K, \quad i = 1, \dots, m, \\ \vec{u} &= \hat{E} \vec{v}, \quad \vec{u} - \vec{U} \in \hat{\mathcal{U}}. \end{aligned} \right\} \quad (\text{SD})$$

This solution satisfies $\vec{v}^* \in \hat{\mathcal{U}}^\perp$.

Idea of the proof

- introduce $G : \mathbb{R}^{Mm} \rightarrow \overline{\mathbb{R}}$,

$$G(\vec{v}) := \hat{F}^*(\vec{v}) + I_{\hat{\mathcal{U}}^\perp}(\vec{v}) - \langle \vec{U}, \vec{v} \rangle, \quad \vec{v} \in \mathbb{R}^{Mm},$$

- $(I_{\hat{\mathcal{U}}^\perp}$ characteristic function of $\hat{\mathcal{U}}^\perp$). G is proper, lsc., strictly convex
- If (\vec{u}, \vec{v}) is a solution to (SD) then \vec{v} is the unique minimizer of G . If \vec{v} is a minimizer of G then $(\hat{E}\vec{v}, \vec{v})$ is a solution to (SD).
- suffices to show that $G(\vec{v}) \rightarrow \infty$ if $\|\vec{v}\| \rightarrow \infty$ (indirect proof).

Discrete Poincaré like inequality (Glitzky '08)

Theorem 2. [Estimate of the free energy by the dissipation rate]

Let (A) and (B) be fulfilled. Moreover, let (\vec{u}^*, \vec{v}^*) be the thermodynamic equilibrium according to Theorem 1. Then for every $\rho > 0$ there exists a constant $c_\rho > 0$ such that

$$\widehat{F}(\vec{u}) - \widehat{F}(\vec{u}^*) \leq c_\rho \widehat{D}(\vec{v})$$

for all $\vec{u} \in \mathcal{N}_\rho := \left\{ \vec{u} = \widehat{E}\vec{v} \in \mathbb{R}^{Mm} : \widehat{F}(\vec{u}) - \widehat{F}(\vec{u}^*) \leq \rho, \vec{u} - \vec{U} \in \widehat{\mathcal{U}} \right\}$.

$$\begin{aligned} \widehat{D}(\vec{v}) = & \sum_{i=1}^m \sum_{k,l \in K, l < k} \bar{u}_i Z_i^{kl} \frac{(v_i^l - v_i^k)^2}{|x^l - x^k|} |\mathbf{S}_i \nu_i^{kl}| |\partial V_i^k \cap \partial V_i^l| \\ & + \sum_{(\alpha, \beta) \in \mathcal{R}} \sum_{k_1 \in K} \cdots \sum_{k_m \in K} R_{\alpha\beta} [v_1^{k_1}, \dots, v_m^{k_m}] \sum_{i=1}^m (\alpha_i - \beta_i) v_i^{k_i} |V_1^{k_1} \cap \cdots \cap V_m^{k_m}| \geq 0. \end{aligned}$$

Energy estimates for (PD)

Theorem 3. [Monotone end exponential decay of the free energy]

We assume (A), (B) and (C). Then the (fully implicit in time) discretization scheme (PD) is dissipative, i.e. solutions (\vec{u}, \vec{v}) to (PD) fulfil

$$\hat{F}(\vec{u}(t_{n_2})) \leq \hat{F}(\vec{u}(t_{n_1})) \leq \hat{F}(\vec{U}) \quad \text{for all } t_{n_1} < t_{n_2}.$$

Moreover, there exists a $\lambda > 0$ such that

$$\hat{F}(\vec{u}(t_n)) - \hat{F}(\vec{u}^*) \leq e^{-\lambda t_n} (\hat{F}(\vec{U}) - \hat{F}(\vec{u}^*)) \quad \forall n \geq 1.$$

Proof

- $\hat{F} : \mathbb{R}^{Mm} \rightarrow \overline{\mathbb{R}}$ is convex, lsc., differentiable in arguments $\vec{u} > 0$.

$$\vec{u} = \hat{E}\vec{v} \implies \vec{v} = \hat{F}'(\vec{u}), \quad \hat{F}(\vec{u}) - \hat{F}(\vec{w}) \leq \hat{F}'(\vec{u}) \cdot (\vec{u} - \vec{w}) \quad \forall \vec{w} \in \mathbb{R}^{Mm}$$

- $\hat{F}(\vec{u}(t_l)) - \hat{F}(\vec{u}(t_{l-1})) \leq (\vec{u}(t_l) - \vec{u}(t_{l-1})) \cdot \vec{v}(t_l) = -\hat{D}(\vec{v}(t_l))$

Proof of Theorem 3

- let $\lambda \geq 0, n_2 > n_1 \geq 0$

$$\begin{aligned} & e^{\lambda t_{n_2}} \left(\widehat{F}(\vec{u}(t_{n_2})) - \widehat{F}(\vec{u}^*) \right) - e^{\lambda t_{n_1}} \left(\widehat{F}(\vec{u}(t_{n_1})) - \widehat{F}(\vec{u}^*) \right) \\ & \leq \sum_{l=n_1+1}^{n_2} h_l e^{\lambda t_{l-1}} \left\{ e^{\lambda h_l} \lambda \left(\widehat{F}(\vec{u}(t_l)) - \widehat{F}(\vec{u}^*) \right) - \widehat{D}(\vec{v}(t_l)) \right\} \end{aligned}$$

- $\widehat{D}(\vec{v}(t_l)) \geq 0 \quad \forall l \geq 1, \quad \text{setting } \lambda = 0 \implies$

$$\widehat{F}(\vec{u}(t_{n_2})) \leq \widehat{F}(\vec{u}(t_{n_1})) \leq \widehat{F}(\vec{U}) \quad \forall t_{n_2} > t_{n_1} \geq 0$$

- $\widehat{F}(\vec{u}(t_l)) - \widehat{F}(\vec{u}^*) \leq \widehat{F}(\vec{U}) - \widehat{F}(\vec{u}^*) =: \rho, \quad \vec{u}(t_l) \in \vec{U} + \widehat{\mathcal{U}} \implies \vec{u}(t_l) \in N_\rho, l \geq 1,$
Theorem 2 supplies $c_\rho > 0$ s.t.

$$\widehat{F}(\vec{u}(t_l)) - \widehat{F}(\vec{u}^*) \leq c_\rho \widehat{D}(\vec{v}(t_l)) \quad \forall l$$

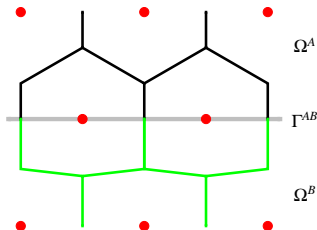
- choose $\lambda > 0$ s.t. $\lambda e^{\lambda \bar{h}} c_\rho < 1, n_1 = 0$

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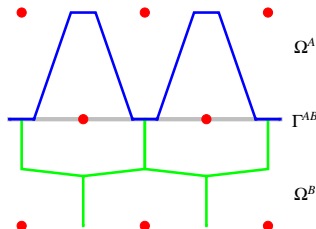
Numerical treatment of heterostructures

$$\overline{\Omega} = \cup_{I \in \mathcal{I}} \overline{\Omega^I}, \quad \Gamma^{AB} = \overline{\Omega^A} \cap \overline{\Omega^B}, \quad \text{on } \Omega^I \text{ constant coefficients}$$

take grid $\{x^k : x^k \in \overline{\Omega}, k = 1, \dots, \widetilde{M}\}$ respecting interfaces Γ^{AB} , $A, B \in \mathcal{I}$, $|\Gamma^{AB}| > 0$
 especially endpoints of Γ^{AB} are grid points



desired



undesired

Numerical treatment of heterostructures

on Ω^I
$$\mathbf{S}_i^I = Q_i^{IT} \text{diag}(\mu_i^{kI}) Q_i^I$$

$$\gamma := \max_{I \in \mathcal{I}} \max_{i=1, \dots, m} \arccos \frac{\min_k(\mu_i^{kI})}{\max_k(\mu_i^{kI})}$$

for $|\Gamma^{AB}| > 0$ let

$$\kappa^{AB} := \frac{\max. \text{ Euclidean dist. of directly neighboring grid points on het. interface } \Gamma^{AB}}{\min. \text{ Euclidean dist. of inner grid points to the het. interface } \Gamma^{AB}}$$

The restriction on the position of vertices on and close to interfaces and boundaries

$$\max_{A, B \in \mathcal{I}, |\Gamma^{AB}| > 0} \kappa^{AB} \leq \sqrt{2 - 2 \sin \gamma}$$

allows to handle general heterostructures and boundary conditions. The described integration procedure can be applied independently on each Ω^I and the fluxes and potentials fulfill the continuity conditions.

Concluding remarks

Results remain true for

(Glitzky '08)

- more general state equations $u_i = \bar{u}_i g_i(v_i - \bar{v}_i)$
mass fluxes

$$j_i = -\bar{u}_i g_i'(v_i - \bar{v}_i) \mathbf{S}_i \nabla v_i$$

(inverse of the Hessian of the free energy for species X_i)

- charged species X_i with charge numbers q_i
add Poisson equation for the electrostatic potential v_0

$$-\nabla \cdot (\mathbf{S}_0 \nabla v_0) = f + \sum_{i=1}^m q_i u_i \quad \text{on } \mathbb{R}_+ \times \Omega$$

$$\nu \cdot (\mathbf{S}_0 \nabla v_0) + \tau v_0 = f^\Gamma \quad \text{on } \mathbb{R}_+ \times \Gamma$$

mass fluxes

$$j_i = -\bar{u}_i g_i'(v_i - \bar{v}_i) \mathbf{S}_i \nabla (v_i + q_i v_0)$$

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