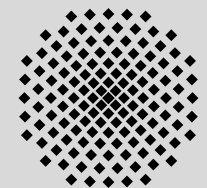


Explicit Discontinuous Galerkin Schemes for Direct and LES Simulations

Gregor Gassner, Frieder Lörcher, Claus-Dieter Munz

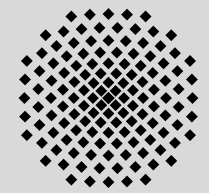
`gassner@iag.uni-stuttgart.de`

Institut for **A**erodynamics and **G**asdynamics
University of Stuttgart



Overview

- Diffusion
- Nodal Integration
- Space-time expansion
- Results



Diffusion I

- Scalar linear heat equation

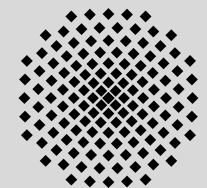
$$u_t - \Delta u = 0 \quad (1)$$

- Standard approach

- Resorting to a first order system (Bassi&Rebay 1997)

$$\begin{aligned} u_t - \nabla \cdot \mathbf{q} &= 0 \\ \mathbf{q} - \nabla u &= 0 \end{aligned} \quad (2)$$

- 'Classic' weak formulation for extended system
- We need numerical approximation of the solution u and \mathbf{q} at grid cell boundaries
- For 3D comp. NSE: 15 auxiliary variables are needed



Diffusion II

- Our approach
 - Use the original second order equation
 - Classic weak formulation

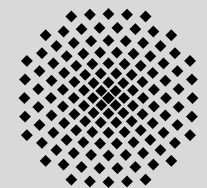
$$\begin{aligned}\langle u_t, \phi \rangle_Q &= \langle \Delta u, \phi \rangle_Q, \\ \langle u_t, \phi \rangle_Q &= (\nabla u \cdot \mathbf{n}, \phi)_{\partial Q} - \langle \nabla u, \nabla \phi \rangle_Q\end{aligned}\tag{3}$$

- second integration by parts ('ultra' weak formulation)

$$\langle u_t, \phi \rangle_Q = (\nabla u \cdot \mathbf{n}, \phi)_{\partial Q} - (u, \mathbf{n} \cdot \nabla \phi)_{\partial Q} + \langle u, \Delta \phi \rangle_Q\tag{4}$$

- New term strong related to symmetric term in SIP scheme (Nitsche 1971)
- no auxiliary variables
- Can be generalized to non-linear systems if viscous terms have the structure

$$U_t = \nabla \cdot \mathbf{F}(U, \nabla U) = \nabla \cdot (\underline{\mathbf{D}}(U) \nabla U)\tag{5}$$



Diffusion III

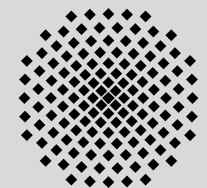
- We need approximation of u and $\nabla u \cdot \mathbf{n}$ at grid cell boundary
 - Diffusive generalized Riemann problem (dGRP)

$$u_t - u_{xx} = 0,$$
$$u(x, t = 0) = \begin{cases} u^- + x u_x^- & \text{for } x < 0, \\ u^+ + x u_x^- & \text{for } x > 0 \end{cases} \quad (6)$$

- dGRP solution yields numerical approximation

$$u_x|_{\partial Q} \approx \eta \frac{[[u]]}{\Delta x} + \{u_x\} \quad \text{and} \quad u|_{\partial Q} \approx \{u\} \quad (7)$$

- Strong relation to the SIP scheme
- η is a known quantity (depends on the order of the DG polynomial)



Diffusion IV

- Extension to non-linear diffusion systems

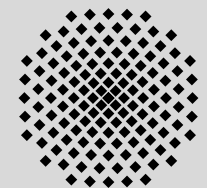
$$U_t - (\underline{D}(U)U_x)_x = 0 \quad (8)$$

- Linearization of the diffusion matrix

$$\underline{D}(U) \approx \underline{\tilde{D}} := \underline{D}(\{U\}) \quad (9)$$

- Diagonalization of $\underline{\tilde{D}}$ ('characteristic' variables)
- Eigenvalues of $\underline{\tilde{D}}$ are the diffusion coefficients of the decoupled system
- Use dGRP solution for each of the scalar heat equations
- Back transformation to conservative variables yields dGRP approximation

$$\underline{D}(U)U_x|_{\partial Q} \approx \underline{\tilde{D}} \left(\eta \frac{\llbracket U \rrbracket}{\Delta x} + \{U_x\} \right) \quad \text{and} \quad U|_{\partial Q} \approx \{U\} \quad (10)$$



Diffusion V

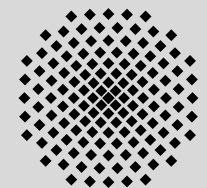
- SIP approach (Hartmann et al. 2006)

$$\underline{D}(U)U_x|_{\partial Q} \approx \eta_{SIP}\lambda_{max} \frac{[[U]]}{\Delta x} + \{D(U)U_x\} \quad \text{and} \quad U|_{\partial Q} \approx \{U\} \quad (11)$$

- For advection diffusion problems with **semi definite** diffusion matrix

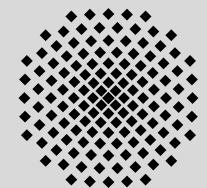
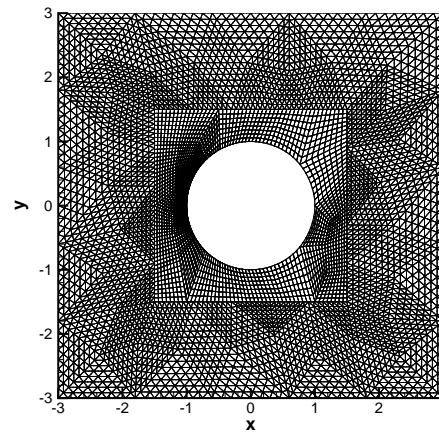
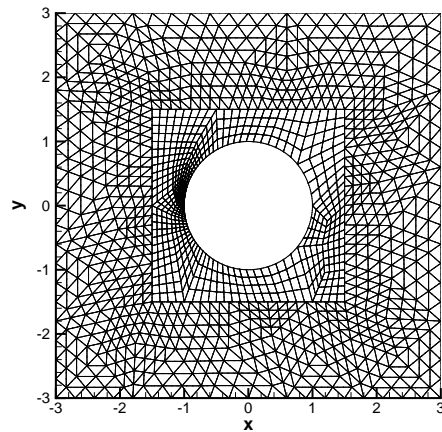
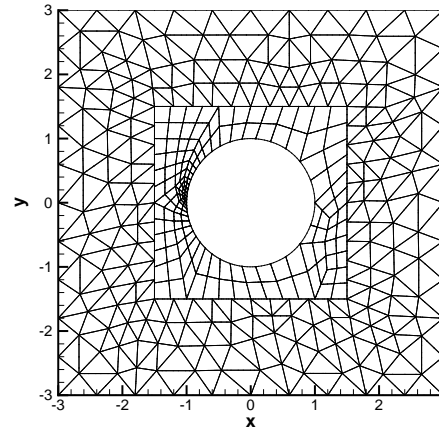
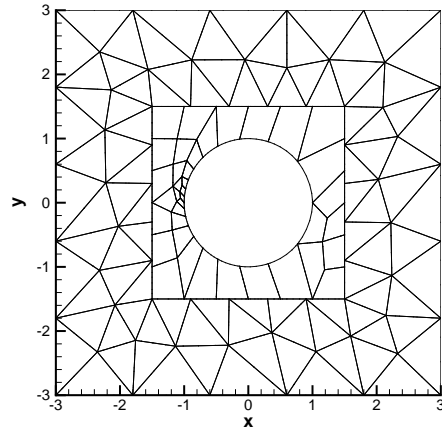
$$EOC_{SIP} = \begin{cases} p+1 & \text{for } p \text{ odd} \\ p & \text{for } p \text{ even} \end{cases} \quad (12)$$

- Compressible NSE (no diffusion in density equation)
 - SIP problem: density jump penalty in density equation
 - dGRP solution: if we multiply jump term with diffusion matrix, no direct penalty term in density equation
 - ⇒ **dGRP** approach yields **optimal EOC** for comp. NSE equations
- we tested several semi definite advection diffusion systems with same result



Diffusion VI

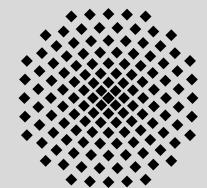
- EOC for compressible Navier-Stokes equations
 - grid sequence with triangles, quadrilaterals, hanging nodes and curved elements for h -convergence



Diffusion VII

- Exact analytical solution
 - Choose smooth analytical function
 - Insert function into comp. NSE
 - Use right hand side as source term
 - Solve inhomogeneous comp. NSE
- Results

refinement	$L_2(\rho e)$	EOC	$L_2(\rho e)$	EOC
	p=4		p=5	
0	$1,02E - 02$	-	$2,91E - 03$	-
1	$3,48E - 04$	4,9	$7,31E - 05$	5,3
2	$1,29E - 05$	4,8	$1,07E - 06$	6,1
3	$4,13E - 07$	5,0	$1,65E - 08$	6,0

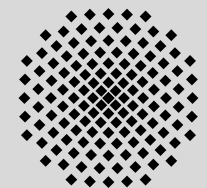


Nodal Integration I

- We choose a **modal** polynomial representation to simplify p -adaption, modal filtering and explicit time stepping. Furthermore, orthogonal hierarchical polynomial trial function are well suited for VMS scale separation.
- Problem: evaluation of the integrals
 - L_2 stability, if integrals computed exactly: not possible for comp. NSE
 - No loss of EOC, if integral evaluation is exact for linear problems
 - Standard approach: numerical integration (Gauss)

$$\left\langle f_1(u), \frac{\partial \phi}{\partial x_1} \right\rangle_Q \approx \sum_{j=1}^{(p+1)^d} f_1(u(\chi_j)) \frac{\partial \phi}{\partial x_1}(\chi_j) \omega_j \quad (13)$$

- Gauss points do not lie on the grid cell surface, therefore can not be reused for surface integral
- Example: $p = 5$ hexahedron needs 216 cubature points for volume integral and 216 cubature points for surface integral (= 432 flux evaluations)



Nodal Integration II

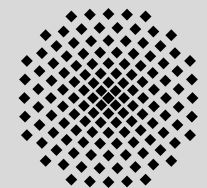
- Our approach: Nodal integration
 - Interpolate p th order polynomial of the non-linear flux f
 - Insert into integral and (pre-)integrate exactly

$$\left\langle f_1(u), \frac{\partial \phi}{\partial x_1} \right\rangle_Q \approx \sum_{i=1}^N f_1(u(\xi_i)) \left\langle \psi_i, \frac{\partial \phi}{\partial x_1} \right\rangle_Q \quad (14)$$

- number of interpolation points

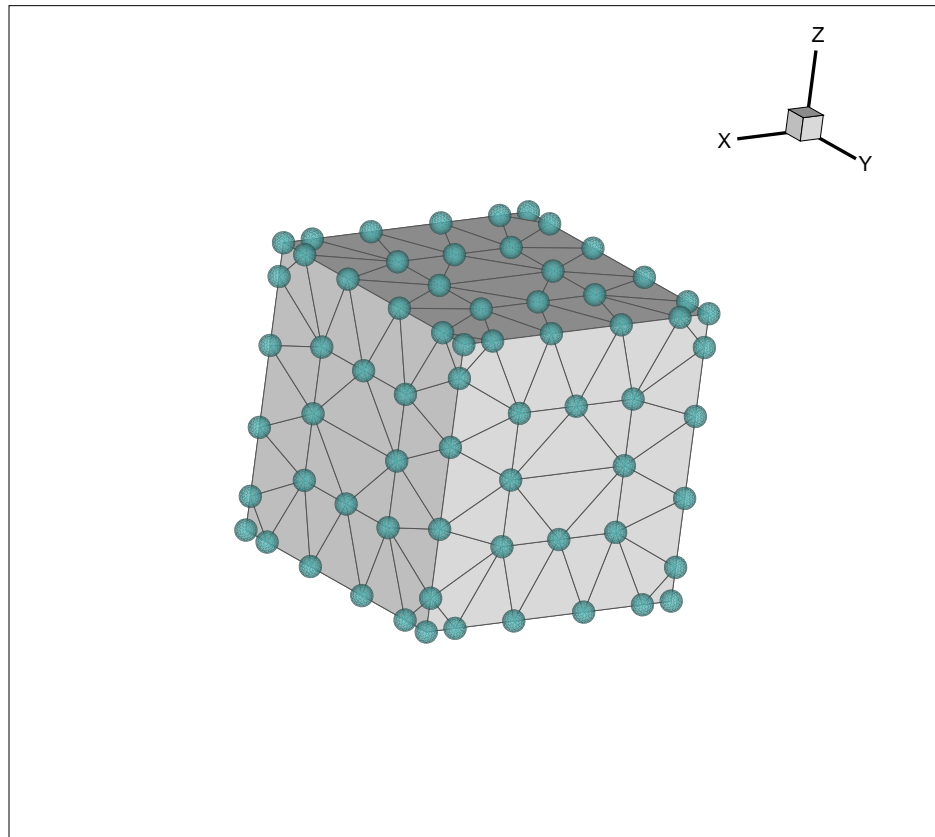
$$N_{min} = \frac{(p+d)!}{d!p!} \leq N \leq N_{max} = (p+1)^d \quad (15)$$

- we can choose int. points that **simultaneous** support a p th order accurate interpolation on the surface
- Example: Interpolation for $p = 5$ hexahedron with $N = 105$ points (432 flux evaluations for Gauss approach)

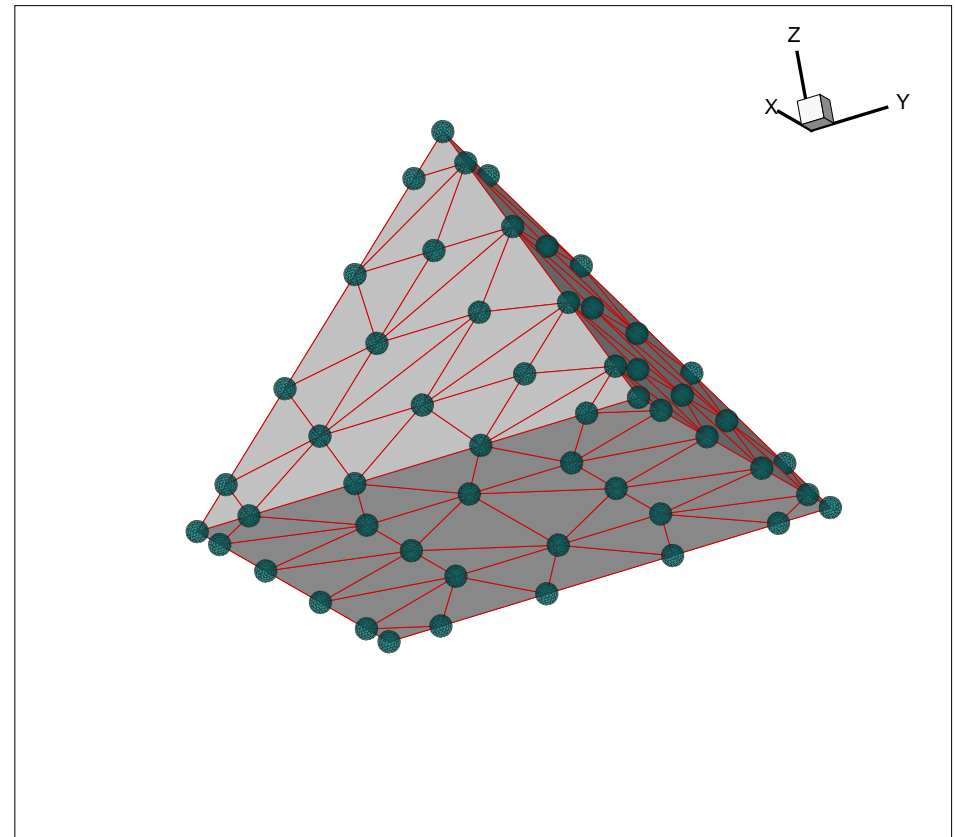


Nodal integration III

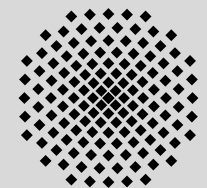
- Interpolation points for $p = 5$.



hexahedron with $N = 105$

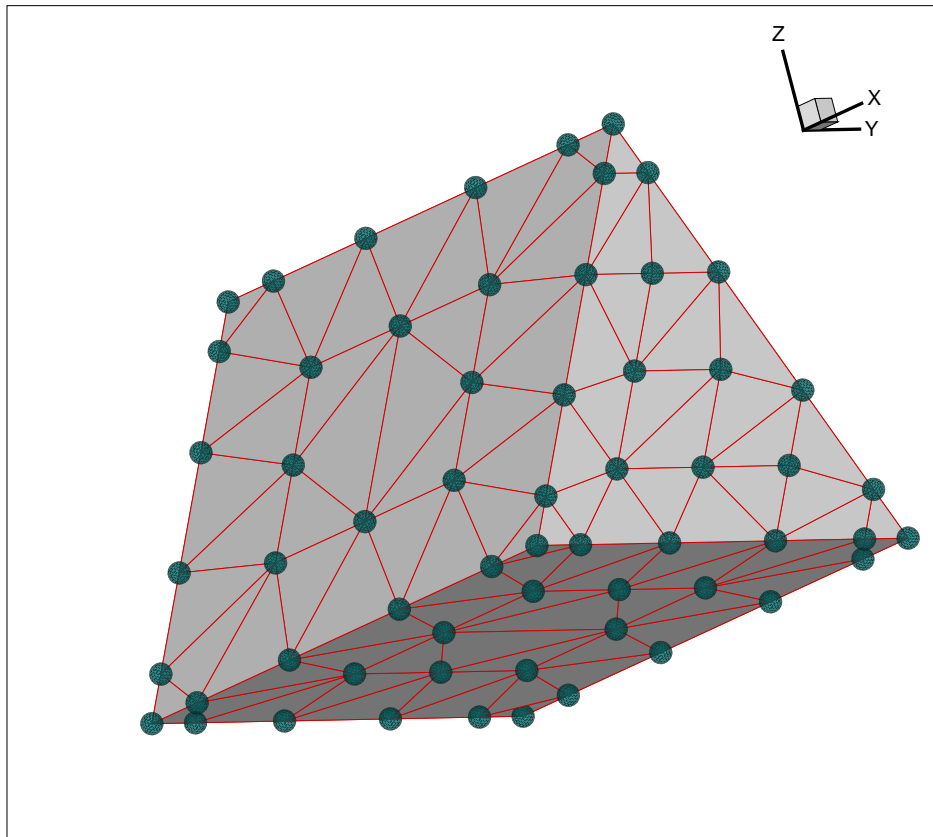


pyramid with $N = 74$

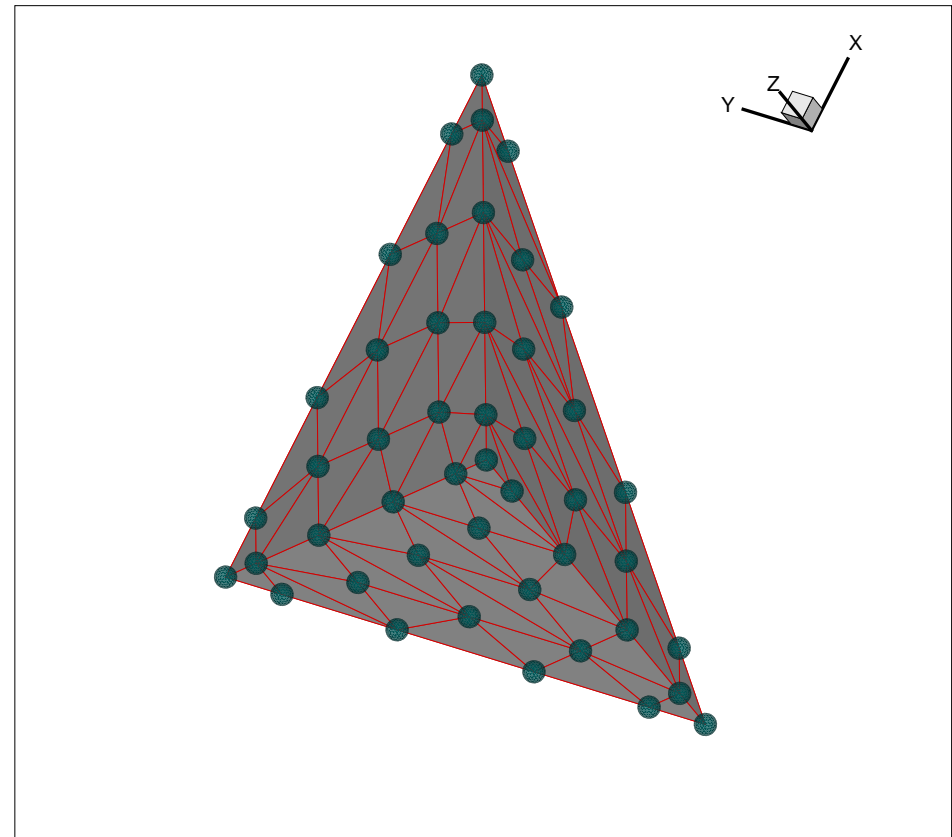


Nodal integration IV

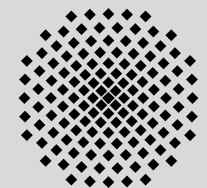
- Interpolation points for $p = 5$.



pentahedron with $N = 84$



tetrahedron with $N = 56$



Space Time Expansion I

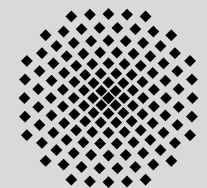
- As we aim to perform large scale computations on massive parallel clusters, we decided to use an explicit time discretization to fully utilize the locality of the DGM
 - Semi discrete DG formulation

$$\langle u_t, \phi \rangle_Q = \underbrace{S(u^\pm, \nabla u^\pm)}_{\text{surface terms}} + \underbrace{V(u, \nabla u)}_{\text{volume terms}} \quad (16)$$

- We simple integrate (16) from time level t_n to t_{n+1}

$$\langle u^{n+1}, \phi \rangle_Q = \langle u^n, \phi \rangle_Q + \int_{t_n}^{t_{n+1}} S(u^\pm, \nabla u^\pm) + V(u, \nabla u) dt \quad (17)$$

- We approximate the time integrals with Gauss quadrature
- We need a high order accurate approximation of the solution $\tilde{u} = \tilde{u}(x, t)$ in the space time grid cell $Q \times [t_n; t_{n+1}]$



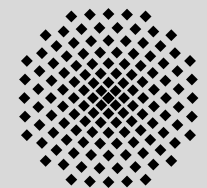
Space Time Expansion II

- We start with a Taylor space time expansion of the approximative solution at $t = t_n$ and $x = x_B$.

$$\tilde{u}(x, t) = \sum_{j=1}^p [(t - t_n) \partial_t + (x - x_B) \cdot \nabla]^j u|_{(x_B, t_n)} \quad (18)$$

- We want that $\tilde{u}(x, t = t_n) = u^n(x)$, thus the pure space derivatives are already available from the DG polynomial
- The mixed space-time and pure time derivatives are approximated using the so-called Cauchy-Kowalewskaya procedure:
 - Example: Burgers equation $u_t = -u u_x$

$$\begin{aligned} u_{tx} &= -u_x^2 - uu_{xx} \\ u_{tt} &= -u_t u_x - uu_{xt} \\ u_{tt} &= -(-uu_x)u_x - u(-u_x^2 - uu_{xx}) \\ &\dots \end{aligned} \quad (19)$$

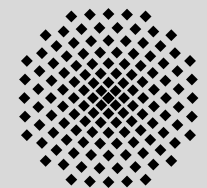


Space Time Expansion III

- The local time stepping feature
 - We drop global time levels t_n and introduce the local time level t_{n_i} for every grid cell Q_i
 - We can use the STE approach, just change integration limits to the grid cells local times

$$\left\langle u^{n+1}, \phi \right\rangle_{Q_i} = \left\langle u^n, \phi \right\rangle_{Q_i} + \int_{t_{n_i}}^{t_{n_i+1}} S(u^\pm, \nabla u^\pm) + V(u, \nabla u) dt \quad (20)$$

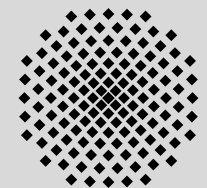
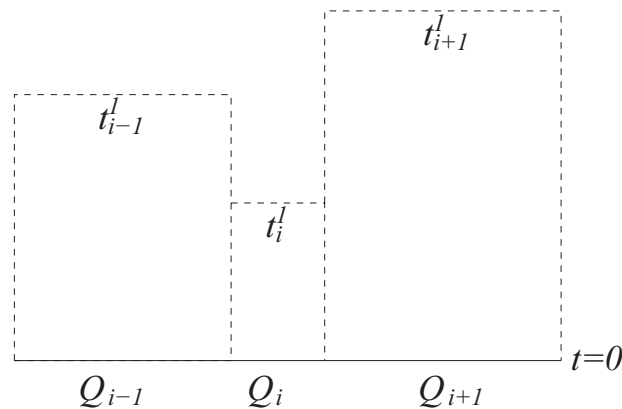
- The RHS of the STE-DG formulation consists of 2 parts:
 - The space-time volume integral needs only local data from the own grid cell. Using the Taylor expansion, this part could be easily evaluated using the approximation \tilde{u} at the space-time Gauss points.
 - The space-time surface integral needs **more care**, as not only local data, but also data from the adjacent grid cells is needed for the numerical fluxes.



Local Time Stepping I

- To get a stable, high order accurate time consistent and conservative scheme, every grid cell has to satisfy the following **evolve condition**

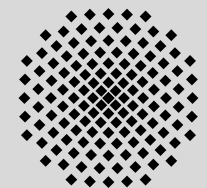
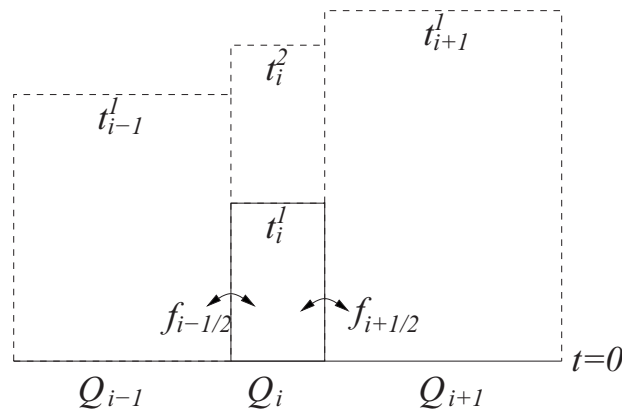
$$t_{n_i+1}^i \leq \min\{t_{n_{i-1}+1}^{i-1}, t_{n_{i+1}+1}^{i+1}\}$$



Local Time Stepping II

- To get a stable, high order accurate time consistent and conservative scheme, every grid cell has to satisfy the following **evolve condition**

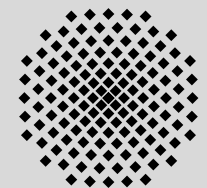
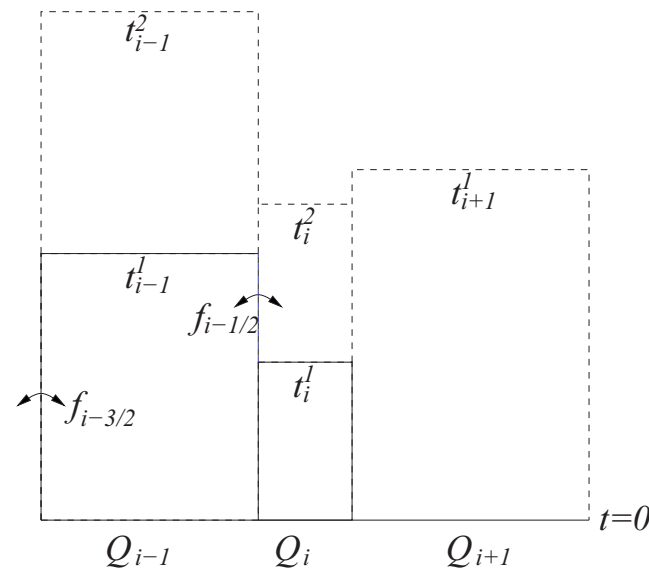
$$t_{n_i+1}^i \leq \min\{t_{n_{i-1}+1}^{i-1}, t_{n_{i+1}+1}^{i+1}\}$$



Local Time Stepping III

- To get a stable, high order accurate time consistent and conservative scheme, every grid cell has to satisfy the following **evolve condition**

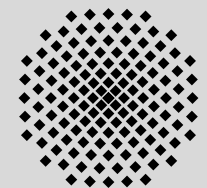
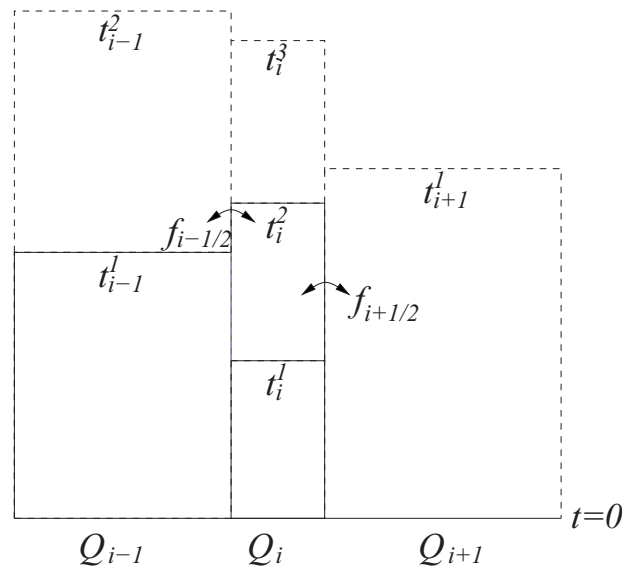
$$t_{n_i+1}^i \leq \min\{t_{n_{i-1}+1}^{i-1}, t_{n_{i+1}+1}^{i+1}\}$$



Local Time Stepping IV

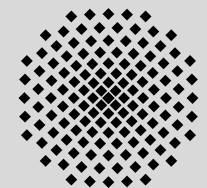
- To get a stable, high order accurate time consistent and conservative scheme, every grid cell has to satisfy the following **evolve condition**

$$t_{n_i+1}^i \leq \min\{t_{n_{i-1}+1}^{i-1}, t_{n_{i+1}+1}^{i+1}\}$$



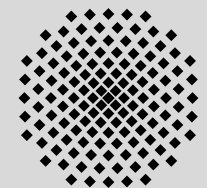
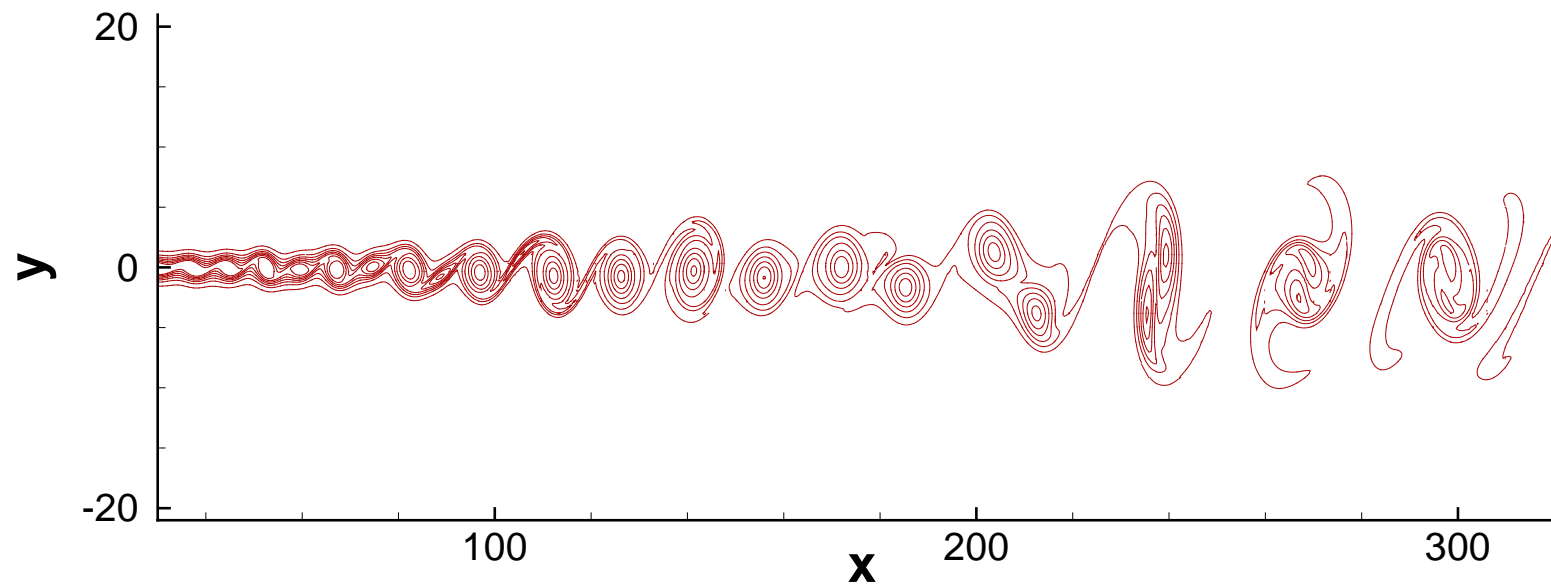
Local Time Stepping V

- Visualization of the local time stepping algorithm
 - 1D Euler equations
 - Irregular grid
 - As the problem is nonlinear, the local time steps depend on the solution!
 - Local time stepping space-time grid
- p -adaptation fully integrated in local time stepping algorithm
- p -adaptation STE-DG scheme fully parallelized and combined with dynamic load balancing
- first tests indicate excellent speed up and scale up efficiency for $\mathcal{O}(1000)$ processors
- more evaluations and applications are needed...



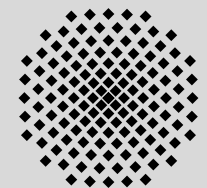
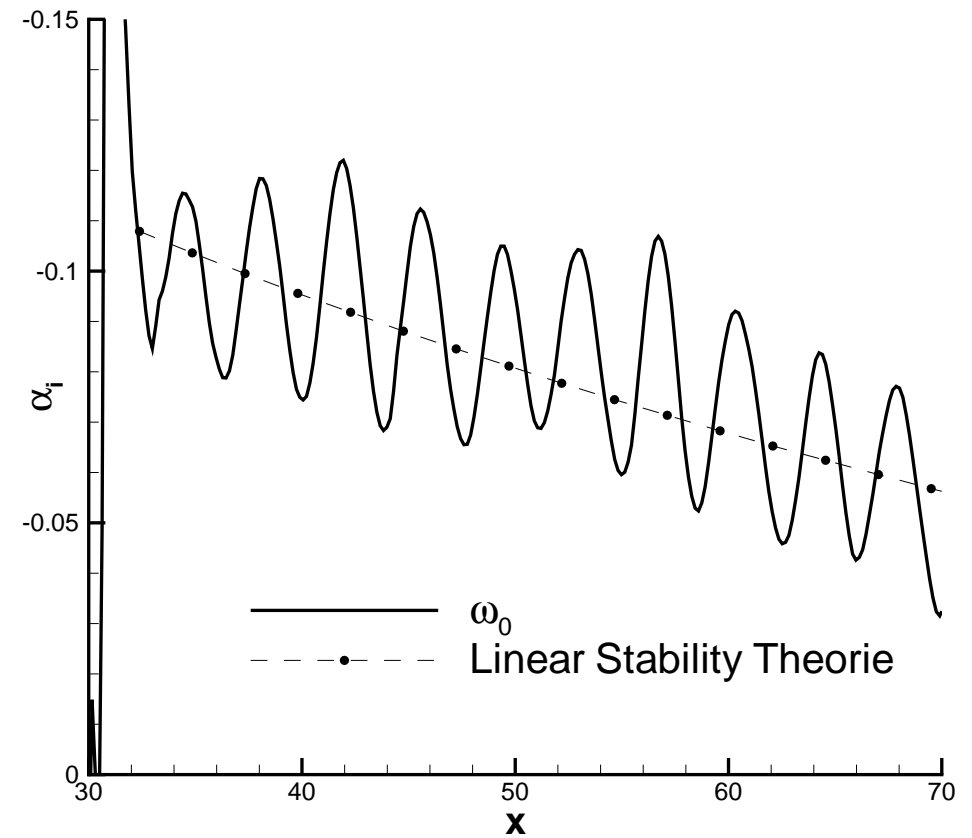
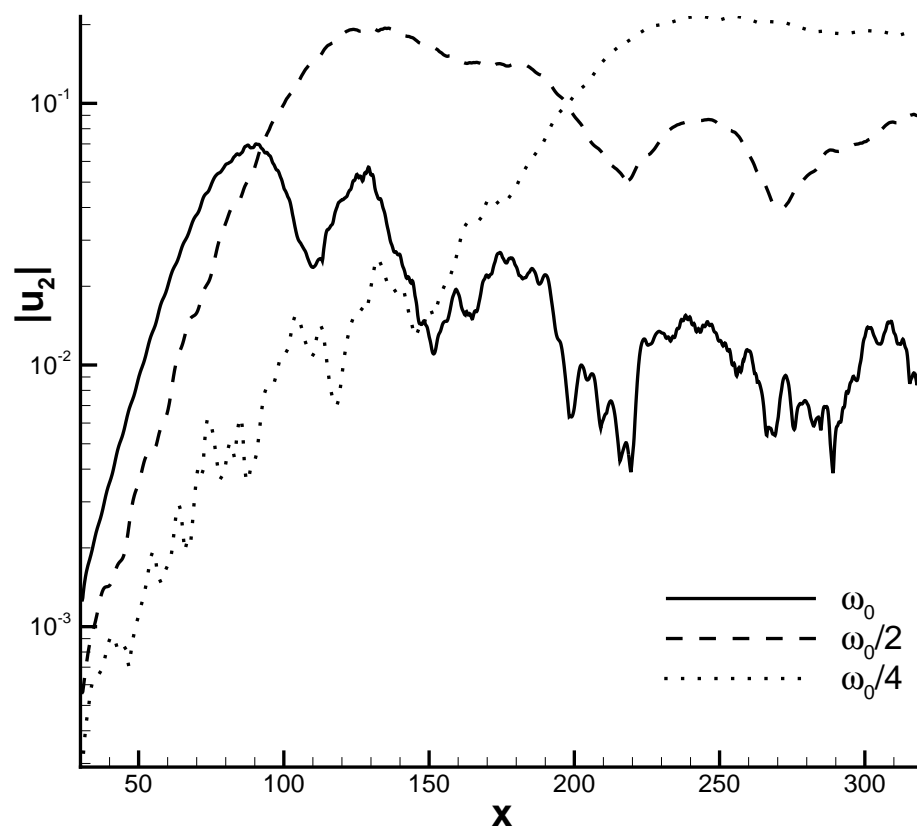
Result I

- Linear Stability Theory
 - Shear layer ($M_1 = 0.5$, $M_2 = 0.25$ and $Re_1 = 500$)
 - 2D compressible Navier-Stokes equations
 - ≈ 180.000 DOF $p = 6$ STE-DG



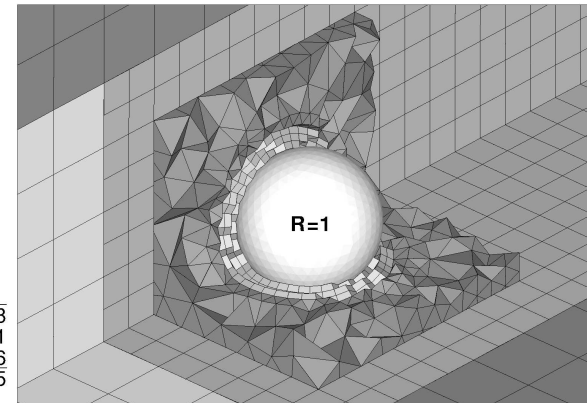
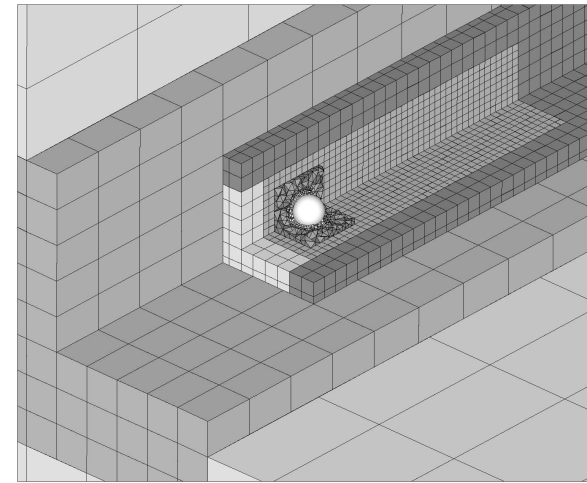
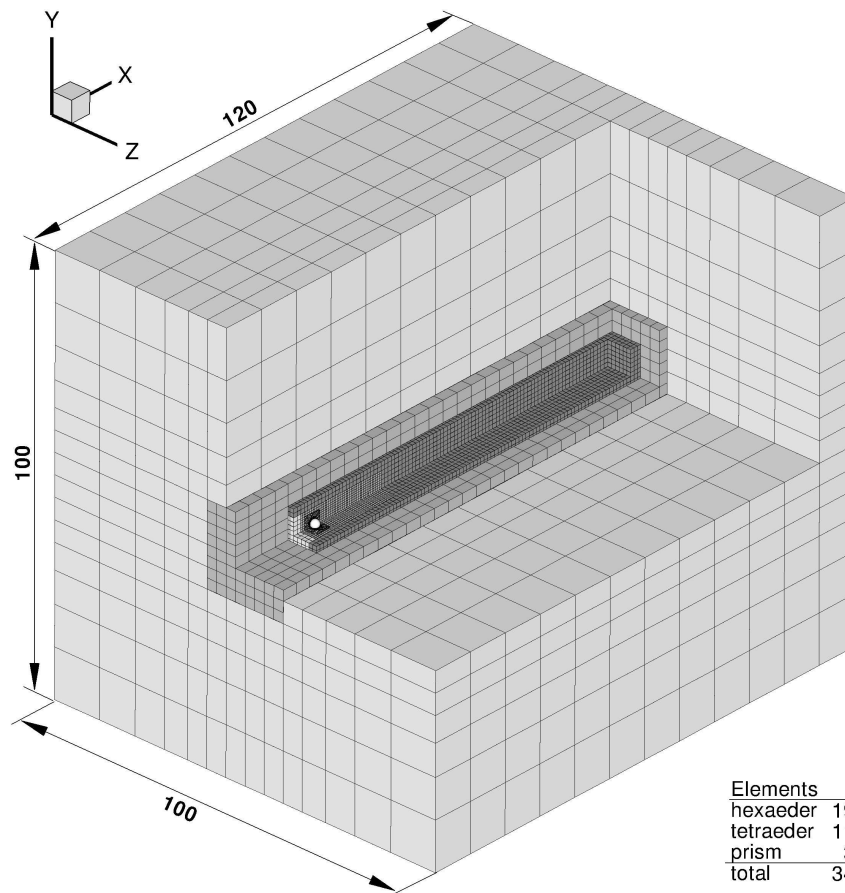
Result I

- Linear Stability Theory



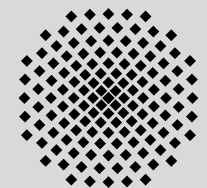
Result II

- Flow past a sphere ($M = 0.1$, $Re_D = 300$)



Elements	
hexaeder	19448
tetraeder	11211
prism	3876
total	34535

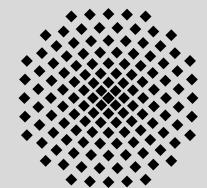
boundary layer $p = 5$, vortex street $p = 4$, far field $p = 1$ (≈ 1.2 Mio. DOF)



Result II

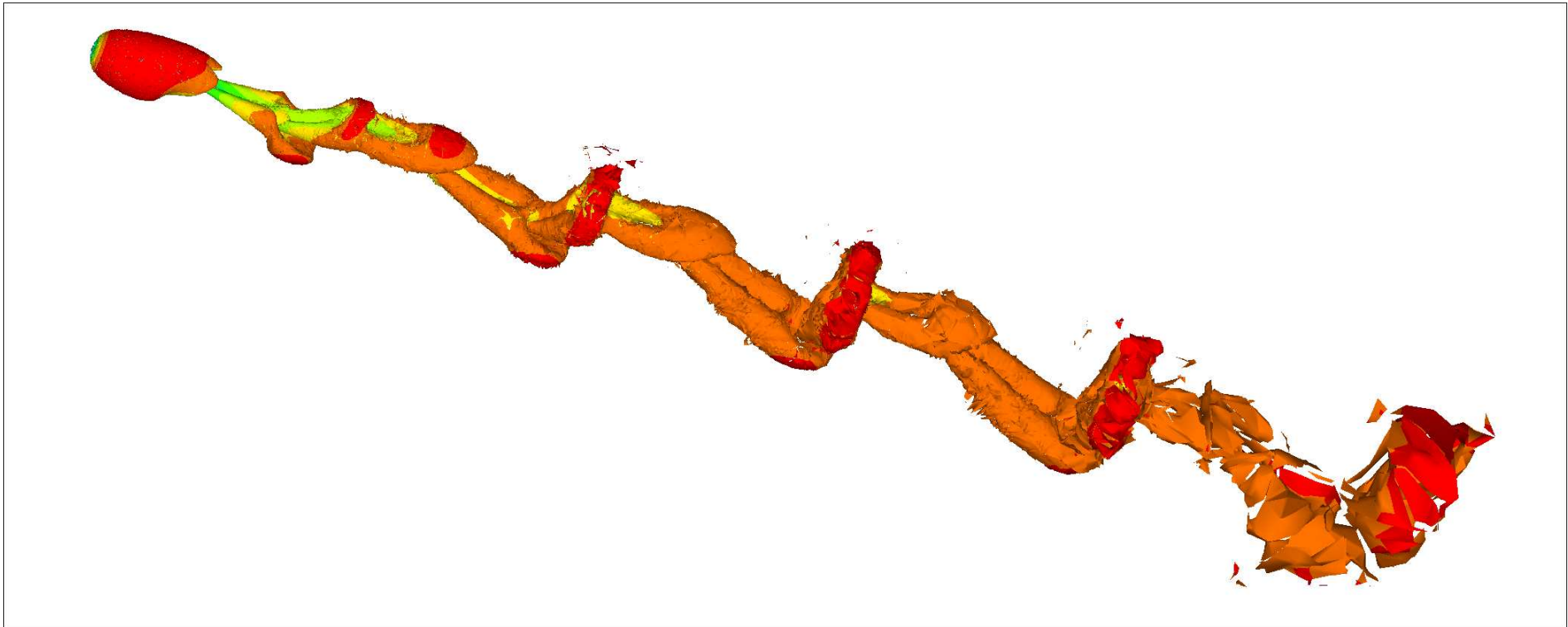
- Flow past a sphere ($M = 0.1$, $Re_D = 300$)
 - Drag, lifting coefficients and Strouhal number

	C_d	ΔC_d	C_l	ΔC_l	Str
	0.673	0.0031	-0.065	0.015	0.135
Tomboulides 1993	0.671	0.0028	—	—	0.136
Johnson&Patel 1999	0.656	0.0035	-0.069	0.016	0.137

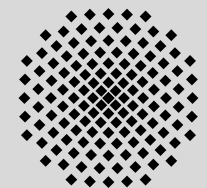


Result II

- Flow past a sphere ($M = 0.1$, $Re_D = 300$)
 - Vortex detection using λ_2 criterion

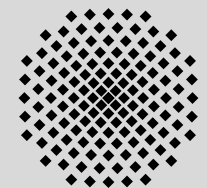


(a) λ_2 iso surface



Outlook

- Hybrid unstructured 3D Code
- Explicit, arbitrary high order accurate in space and time, well suited for high performance computing
- Time accurate local time stepping
- p -adaptation in space and time
- Current Task: combine VMS-LES with STE-DG scheme



Parallelization I

- Scale up test, constant load per processor (150.000 DOF)

np	1	8	64	512	2197
wallclock time(s)	803	816	818	834	851
efficiency (%)	100.0	98.5	98.2	96.4	94.5

- Speed up test: Shear layer (180.000 DOF)

np	50	100	400
Nb DOF/proc	~ 3600	~ 1800	~ 900
efficiency (%)	100.0	101.0	98.1

- Blast Wave Problem with p -adaptation

