

Suitability/Entropy/LES

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Outline

1 BASIC FACTS ABOUT THE NSE



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- 2 GALERKIN APPROX IN TORUS



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- 3 GALERKIN APPROX + DIRICHLET



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- 4 ARE SUITABLE SOLUTIONS USEFUL?



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- 5 NEW STABILIZATION/NUMERICAL TESTS



OUTLINE



Claude Louis Marie Henri
Navier



George Gabriel Stokes

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THE NAVIER-STOKES EQUATIONS

- u : velocity, p : pressure



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- Ω is a bounded fluid domain in \mathbb{R}^3

$$\left\{ \begin{array}{l} \partial_t u + u \cdot \nabla u + \nabla p - \nu \nabla^2 u = f \quad \text{in } \Omega \\ \nabla \cdot u = 0 \quad \text{in } \Omega, \\ u|_{\Gamma} = 0 \quad \text{or } u \text{ is periodic,} \\ u|_{t=0} = u_0, \end{array} \right.$$



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- u_0 is the initial data.
- f a source term.
- ρ is chosen equal to unity.
- ν is viscosity (inverse of Reynolds number).



EXISTENCE

- J. Leray (1934): introduces the notion of **turbulent solution**.
A turbulent solution is a **weak solution** in
 $u \in L^2(0, T; \mathbf{H}^1(\Omega)) \cap L^\infty(0, T; \mathbf{L}^2(\Omega))$
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 $\psi \in D(\mathbb{R}^3), \psi \geq 0, \int_{\mathbb{R}^3} \psi = 1, \psi_\epsilon(x) = \frac{1}{\epsilon} \psi\left(\frac{x}{\epsilon}\right).$

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- E. Hopf (1951) *et al.* uses the **Galerkin technique** to prove existence.



UNIQUENESS

- Are weak solutions **unique** in the large?



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- \Leftrightarrow Are weak solutions **classical** for T large?



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\Rightarrow Clay Institute 1M\$ prize.



SUITABLE WEAK SOLUTION

Definition (V. Scheffer (1976))

A NS weak solution is said to be **suitable weak solutions** iff (u, p) is a weak solution and

$$\partial_t\left(\frac{1}{2}u^2\right) + \nabla \cdot \left(u\left(\frac{1}{2}u^2 + p\right)\right) - \nu \nabla^2\left(\frac{1}{2}u^2\right) + \nu(\nabla u)^2 - f \cdot u \leq 0.$$

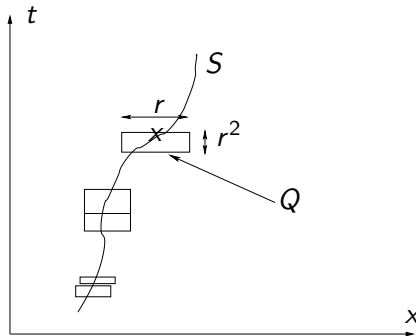
in $\mathcal{D}'((0, T) \times \Omega)$



SUITABLE WEAK SOLUTION

- Singular set: S

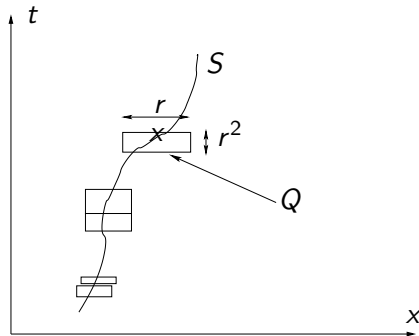
$$S = \{(x, t) \in \Omega \times]0, T[, u \notin L^\infty(V), \forall V \text{ s.t. } (x, t) \in V\}$$



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- $\mathcal{P}^1(S) = \lim_{\delta \rightarrow 0^+} \inf \{ \sum r_i^1, S \subset \cup Q(M_i, r_i), r_i < \delta \}$



SUITABLE WEAK SOLUTION

Theorem (Caffarelli-Kohn-Nirenberg (1982))

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Best partial regularity theorem to date.



EX 1: CONSTRUCTION BY MOLLIFICATION

- Leray's mollification

$$\partial_t u_\epsilon + (\psi_\epsilon * u_\epsilon) \cdot \nabla u_\epsilon + \nabla p_\epsilon - \nu \nabla^2 u_\epsilon = f$$



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Theorem (Leray (1934), Duchon-Robert (2000))

Unique weak solution for all $t > 0$ if $\alpha > \frac{d+2}{4}$, and $u_\epsilon \xrightarrow{\epsilon \rightarrow 0} u$ (up to subsequences) and u is *suitable*.



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- Set $N_\varepsilon = \frac{1}{\varepsilon}$
- Set $\mathbf{X}_\varepsilon = \dot{\mathbb{P}}_{N_\varepsilon}$ (velocity space).
- Let $P_\varepsilon : \dot{\mathbf{L}}^2(\Omega) \longrightarrow \mathbf{X}_\varepsilon$ be L^2 -projection.

$$\dot{\mathbf{L}}^2(\Omega) = \mathbf{X}_\varepsilon \oplus (\mathbf{X}_\varepsilon)^\perp$$



EX 2: CONSTRUCTION BY MOLLIFICATION/NLGM

- Solve for u_ϵ, p_ϵ s.t.

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Theorem

For all $\epsilon > 0$, problem is well-posed (existence + uniqueness).
 $u_\epsilon \rightarrow u, p_\epsilon \rightarrow p$ as $\epsilon \rightarrow 0$ (in appropriate spaces, up to subsequences), u and p are *suitable* weak solution to N.S.



EX 3: CONSTRUCTION BY HYPERVISCOSITY

- Add a vanishing hyperviscosity (Lions (1959)). Ω is the d -torus, d is the space dimension.

$$\left\{ \begin{array}{l} \partial_t u_\epsilon + u_\epsilon \cdot \nabla u_\epsilon + \nabla p_\epsilon - \nu \nabla^2 u_\epsilon + \epsilon^{2\alpha} (-\nabla^2)^\alpha u_\epsilon = f, \\ \nabla \cdot u_\epsilon = 0 \\ u_\epsilon \text{ is periodic,} \\ u_\epsilon|_{t=0} = u_0. \end{array} \right.$$



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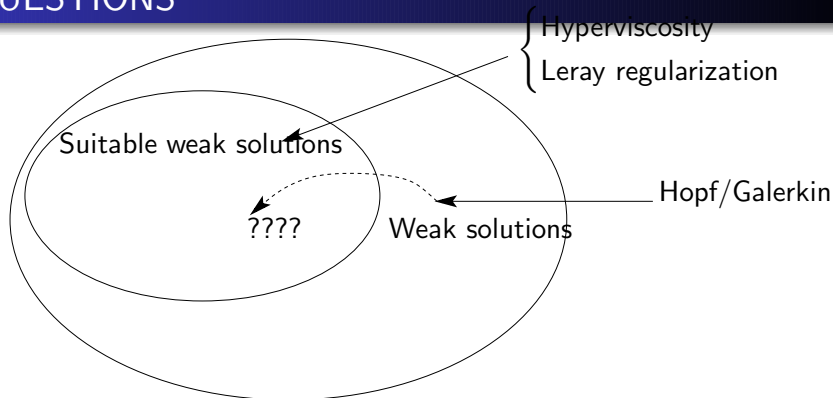
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Theorem (Lions (1959), Beirão da Veiga (1985))

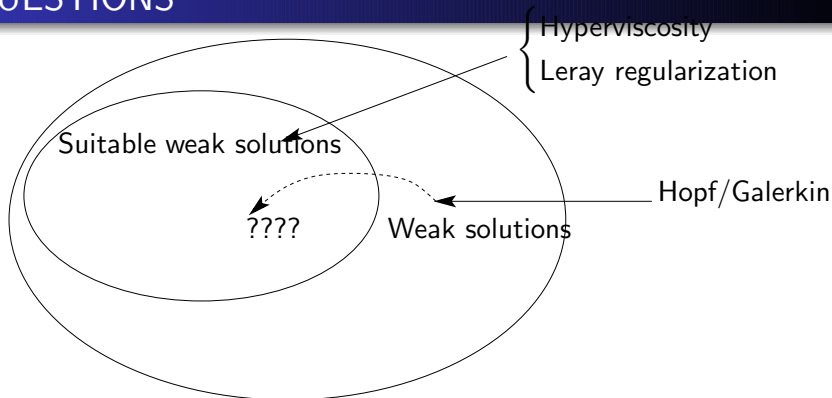
Unique weak solution for all $t > 0$ if $\alpha > \frac{d+2}{4}$, and $u_\epsilon \xrightarrow{\epsilon \rightarrow 0} u$ (up to subsequences) and u is *suitable*.



QUESTIONS



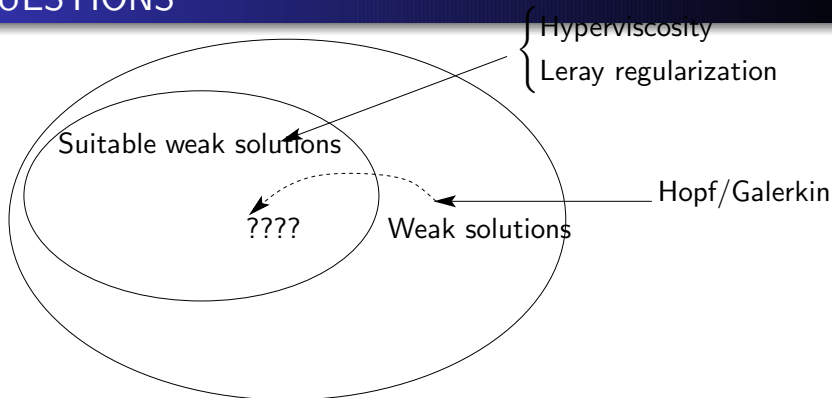
QUESTIONS



Q1: Is the set of suitable solutions a **proper** subset of weak solutions?



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Q2: Do the Galerkin solutions end up to be **suitable** after all?



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Boris Grigorievich Galerkin



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- Modify the nonlinear term as follows:

$$b_h(u, v, v) = \begin{cases} (u \cdot \nabla u + \frac{1}{2} u \nabla \cdot u, v) & \text{(Temam, 1967)} \\ ((\nabla \times u) \times u + \frac{1}{2} \nabla(\mathcal{K}_h(u^2)), v) \end{cases}$$

where $\mathcal{K}_h : L^2(\Omega) \longrightarrow M_h$, linear L^2 -stable approximation operator.



HYPOTHESES/DEFINITIONS (ctd.)

Definition (Discrete commutator property)

There is an operator $P_h \in \mathcal{L}(H_{\#}^1(\Omega); X_h)$ (resp. $Q_h \in \mathcal{L}(L^2(\Omega); M_h)$) such that for all ϕ in $W_{\#}^{2,\infty}(\Omega)$ (resp. all ϕ in $W_{\#}^{1,\infty}(\Omega)$) and all $v_h \in X_h$ (resp. all $q_h \in M_h$)

$$\|\phi v_h - P_h(\phi v_h)\|_{H^l} \leq c h^{1+m-l} \|v_h\|_{H^m} \|\phi\|_{W^{m+1,\infty}}, \quad 0 \leq l \leq m \leq 1$$

$$\|\phi q_h - Q_h(\phi q_h)\|_{L^2} \leq c h \|q_h\|_{L^2} \|\phi\|_{W^{1,\infty}}.$$



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- FE and wavelet-based approximation spaces **have the discrete commutator property** (local interpolation properties).
- Fourier-based approximation spaces **do not have the discrete commutator property** (No local interpolation properties).



GALERKIN FORMULATION

- Seek $u_h \in \mathcal{C}^1([0, T]; X_h)$ and $p_h \in \mathcal{C}^0([0, T]; M_h)$ such that for all $v_h \in X_h$, all $q_h \in M_h$, and all $t \in [0, T]$

$$\begin{cases} (\partial_t u_h, v) + b_h(u_h, u_h, v) - (p_h, \nabla \cdot v) + \nu(\nabla u_h, \nabla v) = \langle f, v \rangle, \\ (\nabla \cdot u_h, q) = 0, \\ u_h|_{t=0} = \mathcal{I}_h u_0, \end{cases}$$

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Theorem (Guermond (2006))

*Under the above hypotheses, if X_h and M_h have **the discrete commutator property**, the couple (u_h, p_h) convergences to a **suitable** solution to NS.*



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- Question was **open** since Scheffer (1977).



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- The main trick: **no boundary condition** \Rightarrow easy estimate on the pressure

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Question: Does the result hold for **Dirichlet BCs**?



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Johann Peter Gustav Lejeune Dirichlet

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- Discrete norm $\|v_h\|_{\mathbf{V}_h^s} := (A_h^s v_h, v_h)^{\frac{1}{2}}, \quad \forall s \in \mathbb{R}.$



PRELIMINARY RESULTS

Lemma

$\exists c_l > 0$ (*non-increasing function*) $\exists c_u > 0$ (*non-decreasing function*), independent of h :

$$c_l(|s|) \|v_h\|_{\tilde{\mathbf{H}}_0^s} \leq \|v_h\|_{\mathbf{V}_h^s} \leq c_u(|s|) \|v_h\|_{\tilde{\mathbf{H}}_0^s}, \quad \begin{cases} -\frac{1}{2} < s < \frac{3}{2}, & \text{lower,} \\ -\frac{3}{2} < s < \frac{3}{2}, & \text{upper} \end{cases}$$



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and for all $s \in (-\frac{3}{2}, 0]$

$$c_l(|s|) \|\nabla_h^2 v_h\|_{\tilde{\mathbf{H}}_0^s} \leq \|A_h v_h\|_{\mathbf{V}_h^s} \leq c_u(|s|) \|\nabla_h^2 v_h\|_{\tilde{\mathbf{H}}_0^s}, \quad \forall v_h \in \mathbf{V}_h.$$



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⇒ Hopf and Leray solutions are suitable



THE MAIN RESULT (ctd.)

Lemma

There is c independent of h so that,

$$\|\partial_t u_h\|_{H^{\tau-1}((0,T); \mathbf{H}^{-\alpha}(\Omega))} + \|u_h\|_{H^{\tau}((0,T); \mathbf{H}^{-\alpha}(\Omega))} \leq c,$$

for all $\alpha \in [\frac{1}{4}, \frac{1}{2})$ and for all $\tau < \bar{\tau} := \frac{2}{5}(1 + \alpha)$.



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Lemma

There is c independent of h such that for $s \in [\frac{3}{10}, \frac{1}{2}]$

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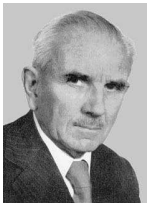
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FE/WAVELETS GALERKIN APPROX IN TORUS
GALERKIN APPROX WITH DIRICHLET BCs
ARE SUITABLE SOLUTIONS USEFUL?
NEW STABILIZATION/NUMERICAL TESTS
CONCLUSIONS/OPEN QUESTIONS

UNDER-RESOLVED SIMULATIONS
A NEW SUBGRID VISCOSITY MODEL?

OUTLINE



Jean Leray



Heinz Hopf

- 1 BASIC FACTS ABOUT THE NSE
- 2 GALERKIN APPROX IN TORUS
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- **Q:** Should we bother about **suitability**?
- **A:** Yes. (What does suitability means after all)?



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⇒ The singular **sub-scales** (if any) are dissipative (at very small scales energy is dissipated)



A NEW SUBGRID VISCOSITY MODEL?

- In under-resolved computations

$$R_h(x, t) := \partial_t u_h - \nu \nabla^2 u_h + u_h \cdot \nabla u_h + \nabla p_h - f \neq 0!$$



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- Under-resolved computations \Leftrightarrow There are singular subscales
- To guaranty that at the grid scale h , energy is well dissipated (**suitability**) we should have

$$R_h(x, t) \cdot u_h \leq 0, \quad \forall x, t$$



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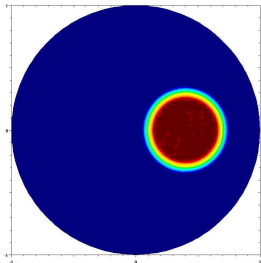
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- Note that $|D_h(x, t)| \rightarrow 0$ if there is no subgrid scale! (no consistency problem).



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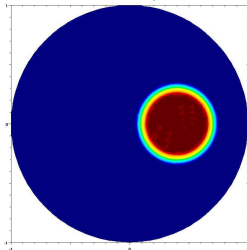


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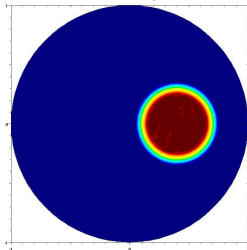
Stabilization? Entropy? Linear transport?

- Solve the transport equation $\partial_t u + \beta \cdot \nabla u = 0$



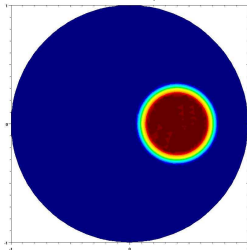
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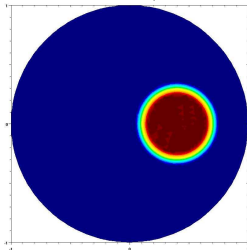


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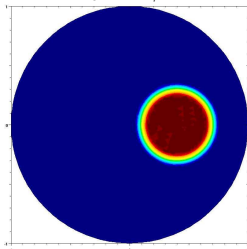


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- Solution method: Galerkin + entropy viscosity.



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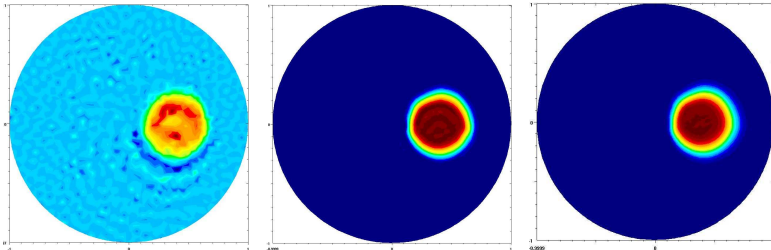
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- Convergence \mathbb{P}_k : L^1 -norm $\mathcal{O}(h^{\frac{k+1}{k+2}})$, L^2 -norm $\mathcal{O}(h^{\frac{k+1}{2(k+2)}})$?



Stabilization? Entropy? Linear transport?

- Numerical test \mathbb{P}_1 , $h = 0.05$, $T = 1$



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Buckley Leverett, \mathbb{P}_2 FE

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Non-convex fluxes (composite waves)

$$u(x, y, 0) = \begin{cases} 1, & \sqrt{x^2 + y^2} \leq 0.5 \\ 0, & \text{else} \end{cases}$$



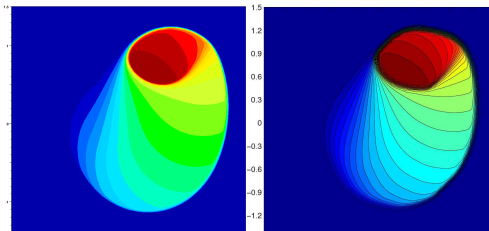
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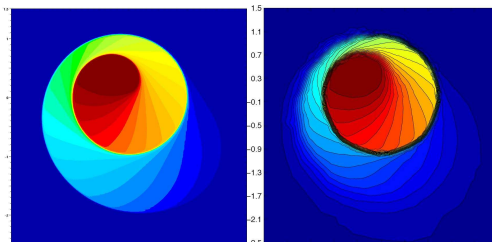
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Euler flows

- Solve 1D Euler equations for perfect gas,
 $(\gamma - 1)e = T = p/\rho$, $\gamma = 1.4$
- Entropy $S = \frac{\rho}{\gamma-1} \log(p/\rho^\gamma)$
- Entropy residual, $D_h(u) := \partial_t S + \partial_x(uS)$
- Define wave speed $\beta := |u| + (\gamma T)^{\frac{1}{2}}$



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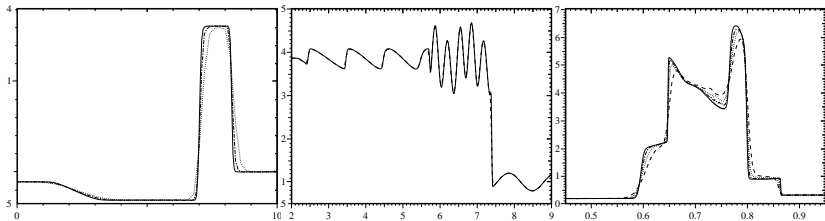


Figure: Lax shock tube, $t = 1.3$, 50, 100, 200 points. Shu-Osher shock tube, $t = 1.8$, 400, 800 points. Right: Woodward-Collela blast wave, $t = 0.038$, 200, 400, 800, 1600 points.



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- A new (very simple) stabilization technique has been proposed and tested.

