

Some achievements in multiscale subgrid modelling

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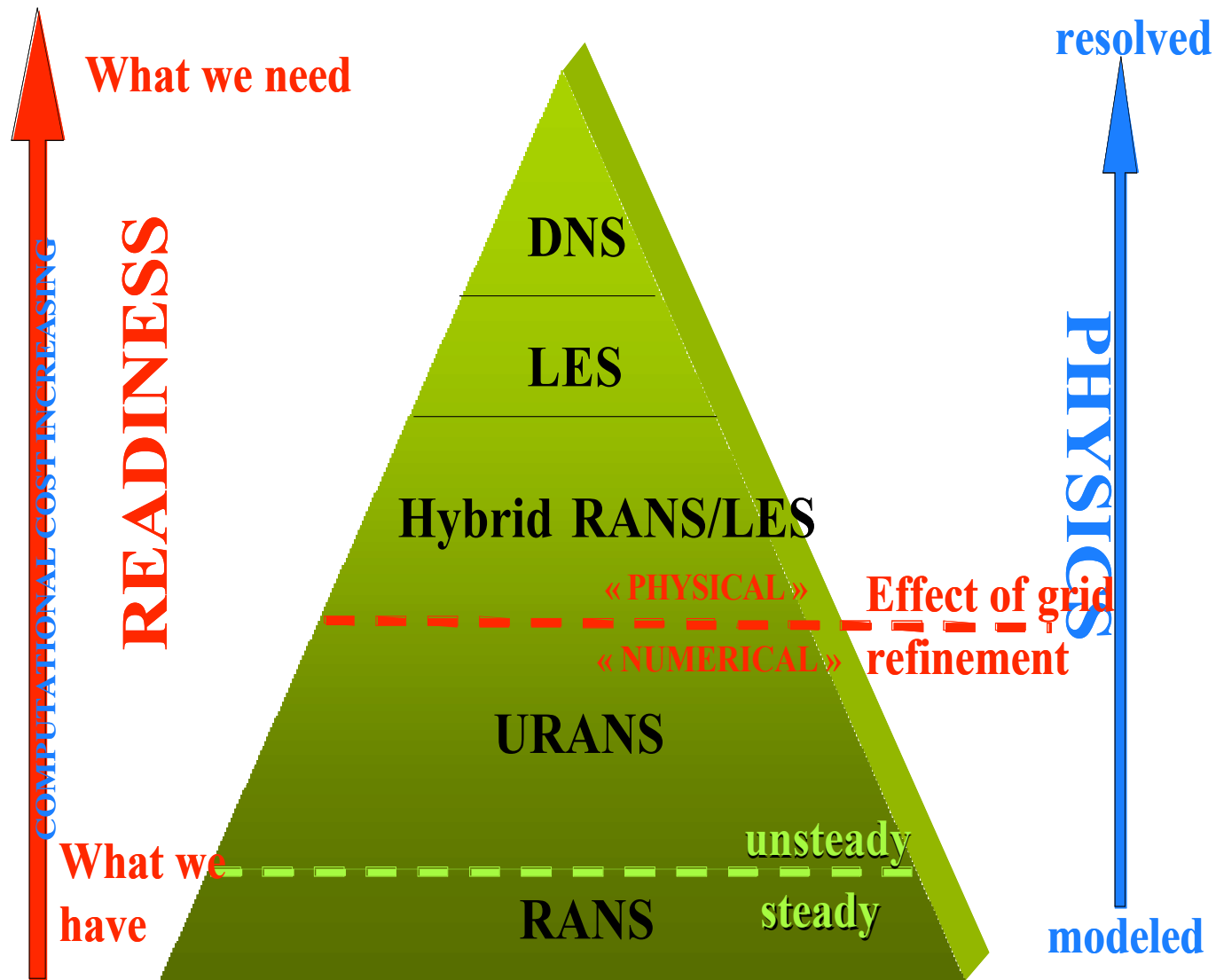
M. Ciardi (Cambridge Univ.), D. Lakehal (ETHZ), J. Meyers (KUL), V. Levasseur (Dassault)

Workshop VMS 2008

Saarbrück, Germany

23-24 june, 2008

Hierarchy of CFD approaches



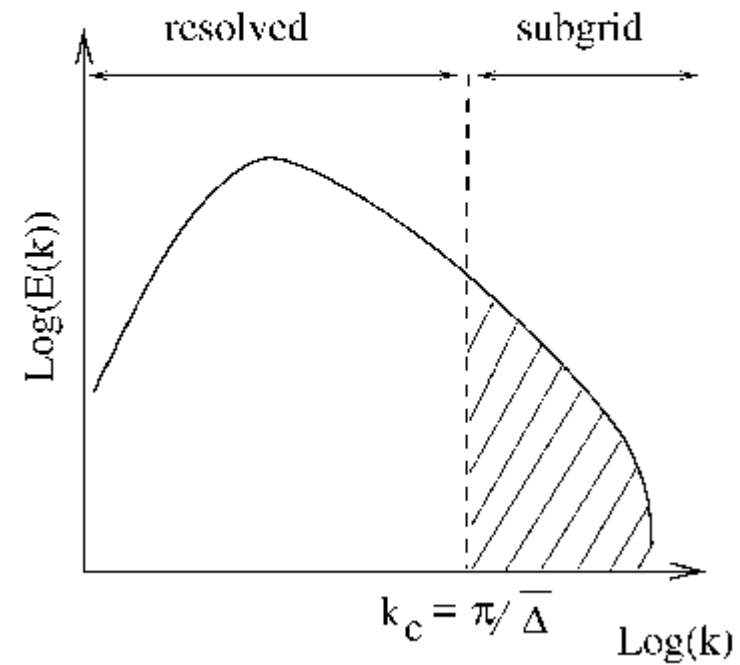
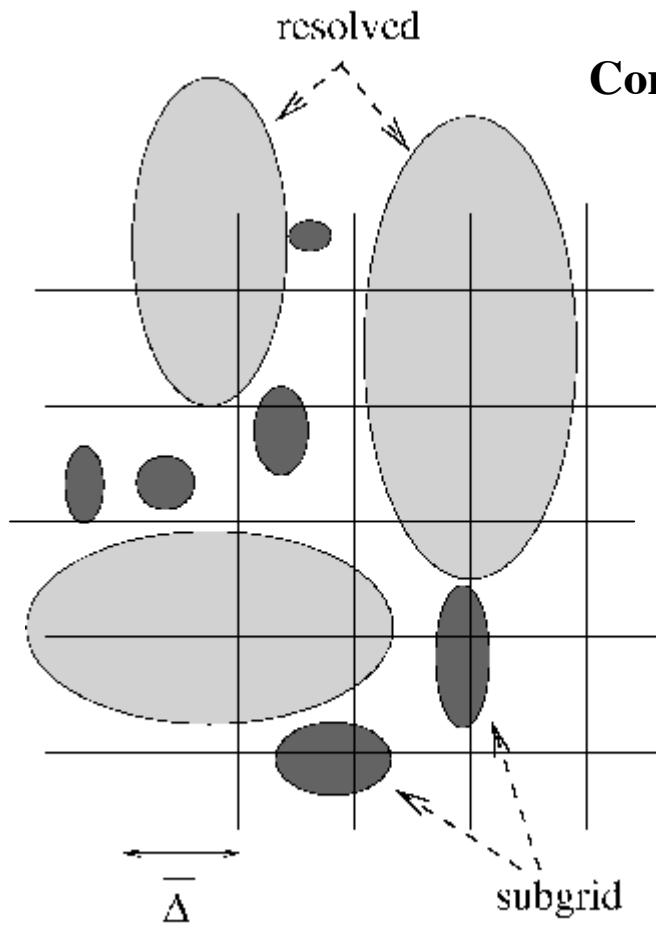
From "Multiscale and multiresolution approaches in turbulence"
P.Sagaut, S.Deck, M. Terracol, Imperial College Press, 2006

The LES concept

$$\bar{u}(x, t) = G \star u(x, t)$$

Computed 'filtered'
solution

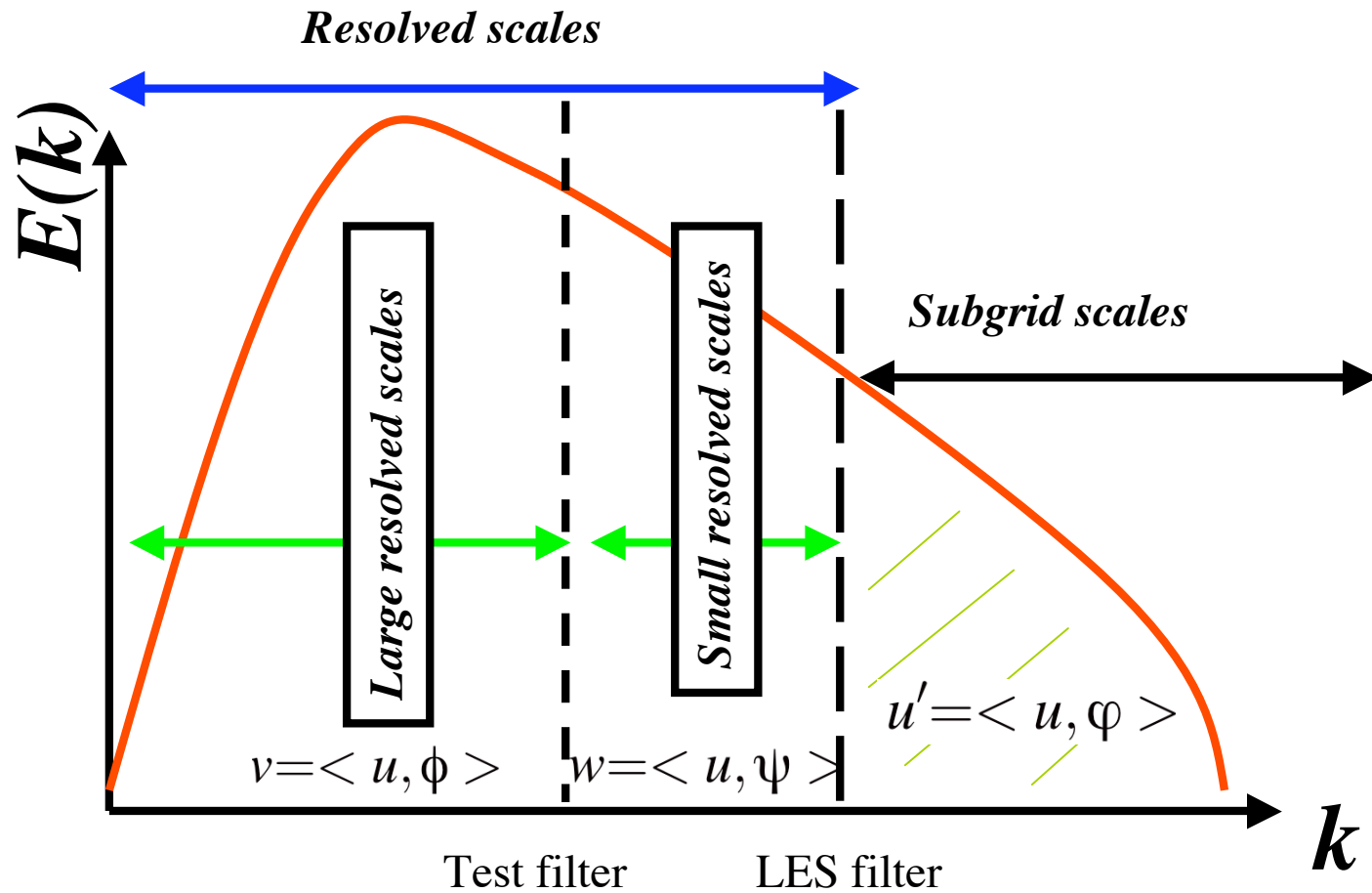
Exact solution



- Brief reminder about VMS-LES methods
- A few examples of use at UPMC
- A remark dealing with numerical errors/subgrid model coupling
- On the value of the constant in VMS Smagorinsky model

Schematic view of VMS-LES strategies

$$\frac{\partial u}{\partial t} + F(u, u) = 0 \quad u = \bar{u} \oplus u' = v \oplus w \oplus u'$$



Neglected (or modeled)

$$\frac{\partial v}{\partial t} + \langle F(\bar{u}, \bar{u}), \phi \rangle = -2 \langle F(\bar{u}, u'), \phi \rangle - \langle F(u', u'), \phi \rangle$$

$$\frac{\partial w}{\partial t} + \langle F(\bar{u}, \bar{u}), \psi \rangle = -2 \langle F(\bar{u}, u'), \psi \rangle - \langle F(u', u'), \psi \rangle$$

modeled

$$\frac{\partial u'}{\partial t} + \langle F(u, u), \varphi \rangle = 0$$

Ignored (or simplified PDEs ?)

General form of VMS subgrid viscosities

$$r.h.s. = \langle \mathbf{v}_t(h, X) (\nabla Y + \nabla^t Y), Z \rangle$$

Subgrid viscosity Strain tensor Selected subspace

Viscosity test field:

- $X=v$: Large - ??? Model
- $X=w$: Small - ??? Model
- $X=v+w$: All - ??? Model

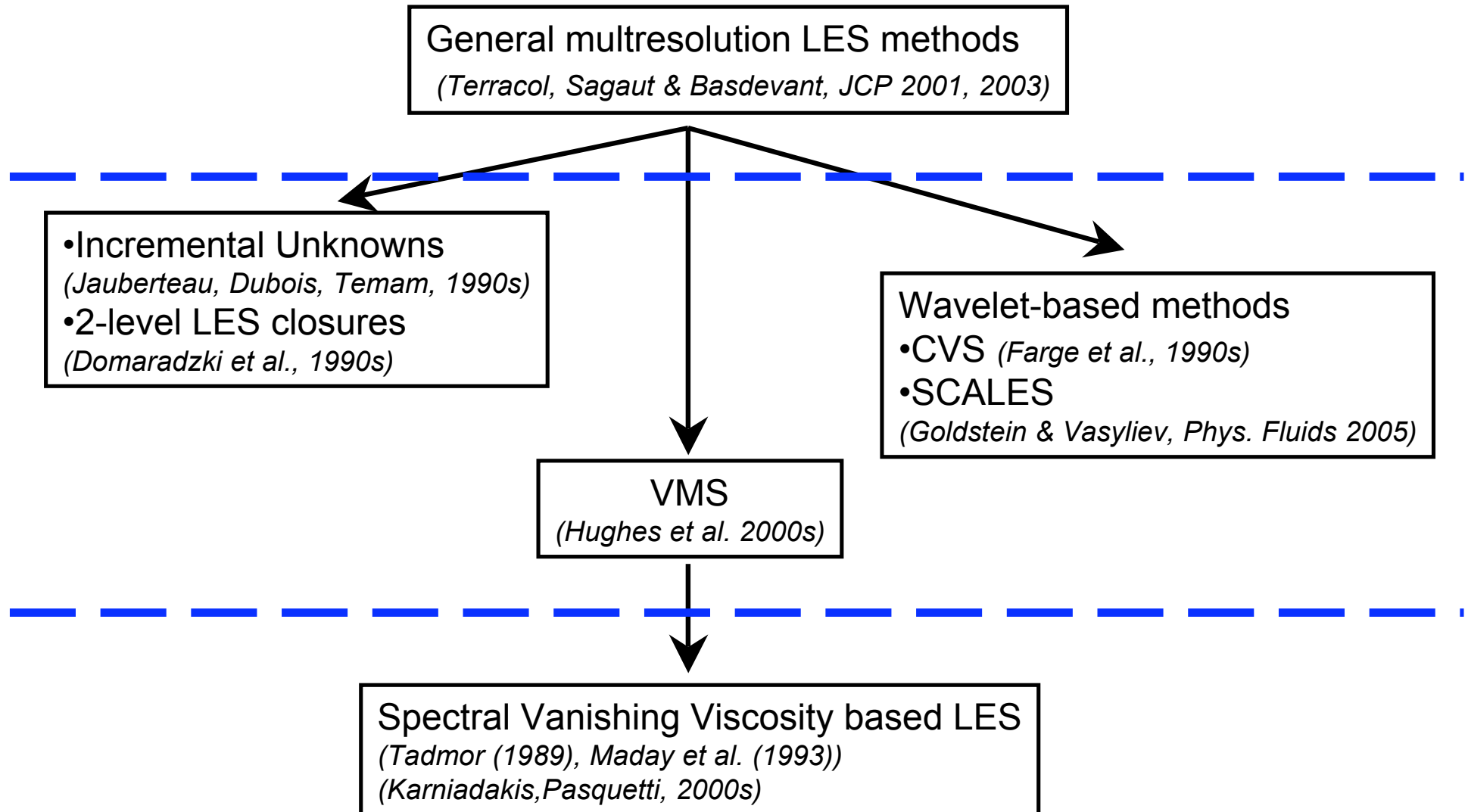
Strain test field:

- $Y=v$: ???- Large Model
- $Y=w$: ???-Small Model
- $Y= v+w$: ???- All Model

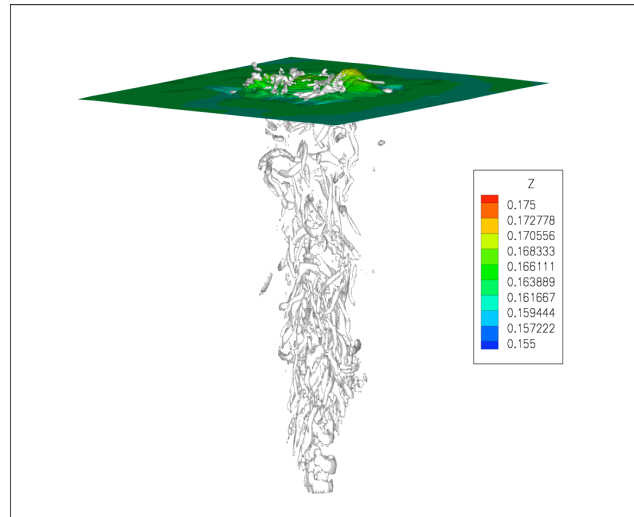
Target subspace:

- $Z=\phi$
- $Z= \psi$
- $Z= \phi + \psi$

Hierarchy of multiscale/level LES methods

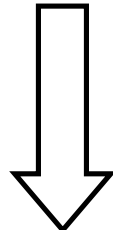


Jet/free surface interaction
(with TREFLE Lab.)



Numerical methods	Physical model	Flow	Reference
Spectral method	incompressible	Isotropic turbulence	<i>Sagaut & Levasseur, Phys. Fluids, 2005</i>
Spectral/FD	incompressible	Channel flow	<i>Meyers & Sagaut, Phys. Fluids, 2007</i>
Unstructured GLS	compressible	Isotropic turbulence	<i>Levasseur & al., CMAME</i>
Unstructured GLS	compressible	Weapon bay	<i>Levasseur & al., JFS, in press</i>
Unstructured FV	compressible	Isotropic turbulence, channel flow	<i>Sagaut & Ciardi, Phys. Fluids, 2006</i>
Spectral method	incompressible	Free surface channel	<i>Reboux & al. Phys. Fluids, 2006</i>
Structured FD	incompressible	Free surface flow	<i>Moreau & al., submitted</i>

$$\frac{\delta u_h}{\delta t} + F_h(u_h, u_h) = 0$$

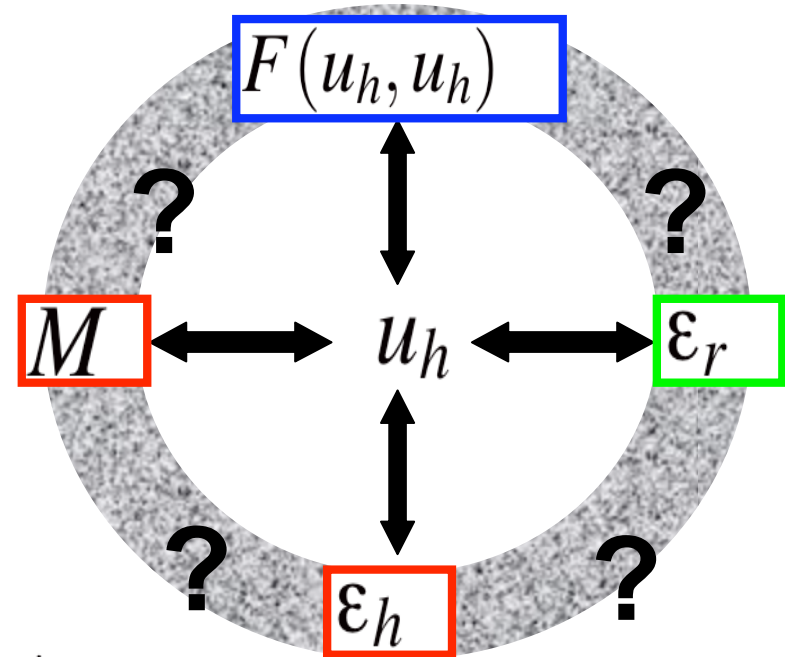


$$\frac{\partial u_h}{\partial t} + F(u_h, u_h) = M + \varepsilon_r + \varepsilon_h$$

Exact subgrid model

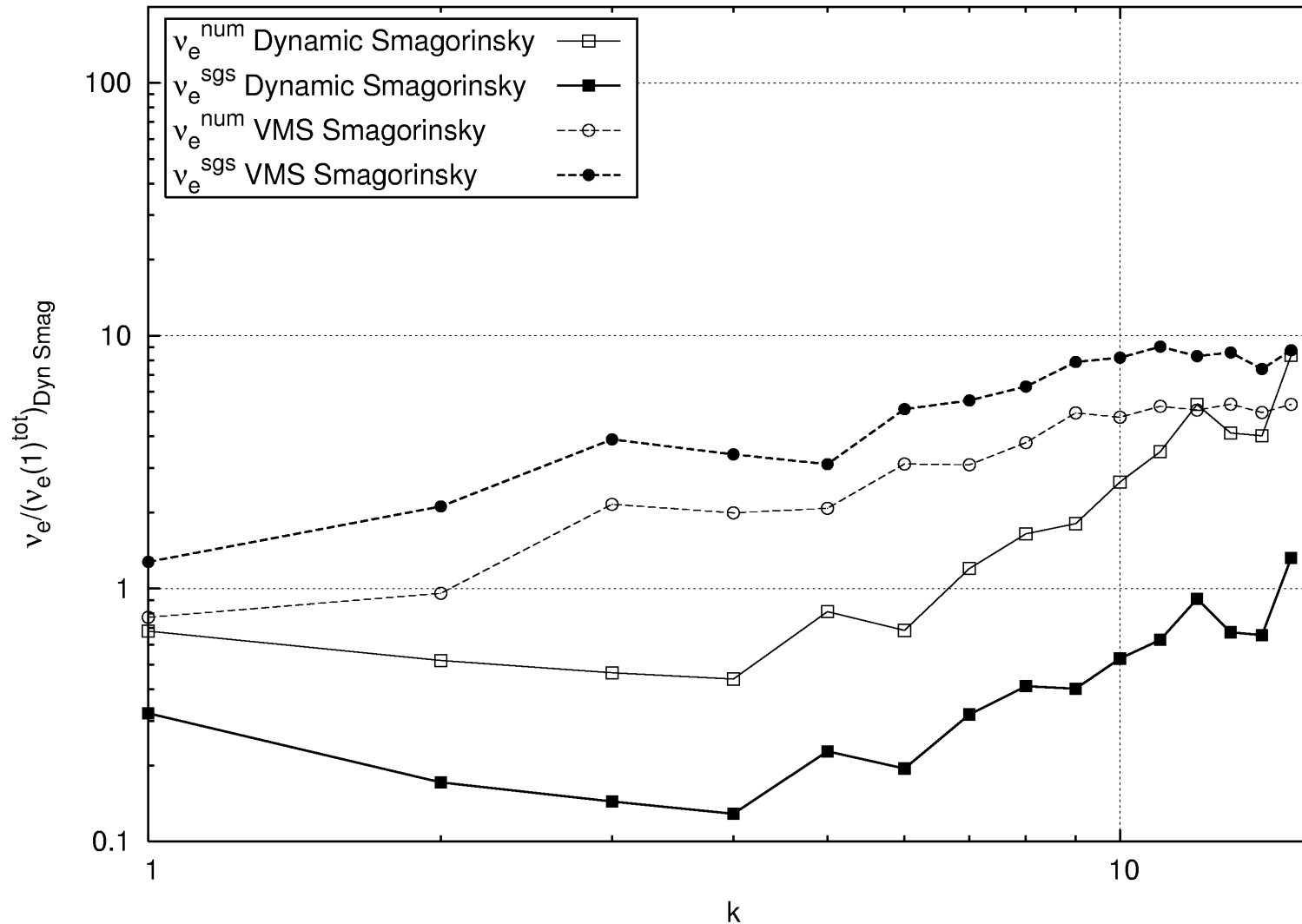
Subgrid modelling error

numerical error



Numerical error/subgrid model interactions

(Ciardi & al., *J. Comput. Phys.*, 2005)
 (Sagaut & Ciardi, *Phys. Fluids*, 2006)



- A consistent subgrid model must account for exact energy transfers between resolved and subgrid scales

⇒ sensitivity to two non-dimensional parameters

$$L/\Delta \quad \eta/\Delta$$

- Classical subgrid models designed for (asymptotic canonical case)

$$L/\Delta \gg 1 \quad \eta/\Delta \ll 1$$

⇒ problems in realistic cases, DNS not recovered satisfactorily

Pope's spectrum model (*Pope, 2000*)

$$E(k) = K_0 \varepsilon^{2/3} k^{-5/3} f_L(kL) f_\eta(kL Re_L^{-3/4})$$

$$f_L(x) = \left(\frac{x}{\sqrt{x^2 + c_L}} \right)^{11/3} \quad \text{Large scale part (flow-dependent)}$$

$$f_\eta(x) = \exp \left(-c_\beta \left((x^4 + c_\eta^4)^{1/4} - c_\eta \right) \right) \quad \text{Viscous range}$$

Small-Small VMS Smagorinsky

- Subgrid model

$$m_{ij} = -[2C_{s1}^2 \Delta^2 |\bar{S}'| |\bar{S}'_{ij}|]'$$

- Induced subgrid dissipation

$$\varepsilon_{t,s1} = (C_{s1} \Delta)^2 \langle 2\bar{S}'_{ij} \bar{S}'_{ij} \rangle^{3/2} = (C_{s1} \Delta)^2 \left(2 \int_0^\infty k^2 (H'(k))^2 \bar{E}(k) dk \right)^{3/2}$$

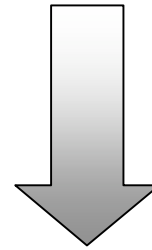
Test filter

Resolved field spectrum

Viscous effects and total dissipation

Normalized total dissipation

$$1 = \frac{\varepsilon_t}{\varepsilon} + \frac{\varepsilon_v}{\varepsilon}$$



$$1 = \gamma_1^2 C_{s1}^2 \left(\frac{3\alpha}{2}\right)^{3/2} \pi^2 (1 - \beta^{4/3})^{3/2} \Psi_1^{3/2} + \left(\frac{3\alpha}{2}\right) \frac{(\gamma \pi L / \Delta)^{4/3}}{Re_L} \Phi$$



Viscous dissipation
neglected in usual analysis

Inertial-range consistency requires to account for this contribution

Scaling functions

$$\beta = \Delta' / \Delta$$

$$\gamma = \frac{\Delta}{\pi} \left(\frac{4}{3} \int_0^{\infty} x^{1/3} G^2(x) dx \right)^{3/4} \quad \gamma_1 = \left(\frac{4 \int_0^{\infty} k^{1/3} (H'(k))^2 (G(k))^2 dk}{3 (\pi/\Delta)^{4/3} (1 - \beta^{4/3})} \right)^{3/4}$$

$$\Phi(L/\Delta, Re_L) = \frac{4}{3} \frac{1}{(\gamma\pi L/\Delta)^{4/3}} \int_0^{\infty} x^{1/3} G^2(x/L) f_L(x) f_{\eta}(x Re_L^{-3/4}) dx$$

$$\Psi_1 \left(\frac{L}{\Delta}, Re_L \right) = \frac{4 \int_0^{\infty} x^{1/3} (H'(x/L))^2 (G(x/L))^2 f_L(x) f_{\eta}(x Re_L^{-3/4}) dx}{3 \int_0^{\infty} x^{1/3} (H'(x/L))^2 (G(x/L))^2 dx}$$

⇒ filter & spectrum dependent

$$\begin{aligned}
 C_{s1} &= \frac{C_{s,\infty}}{\gamma_1} \frac{\Psi_1^{-3/4}}{(1 - \beta^{4/3})^{3/4}} \sqrt{1 - \left(\frac{\gamma L}{C_{s,\infty} \Delta}\right)^{4/3} \frac{1}{Re_L} \Phi} \\
 &= \frac{C_{s,\infty}}{\gamma_1} \frac{\Psi_1^{-3/4}}{(1 - \beta^{4/3})^{3/4}} \sqrt{1 - \left(\frac{\gamma \eta}{C_{s,\infty} \Delta}\right)^{4/3} \Phi}.
 \end{aligned}$$

⇒ a single universal value doesn't exist !

All-Small VMS Smagorinsky

- Subgrid model

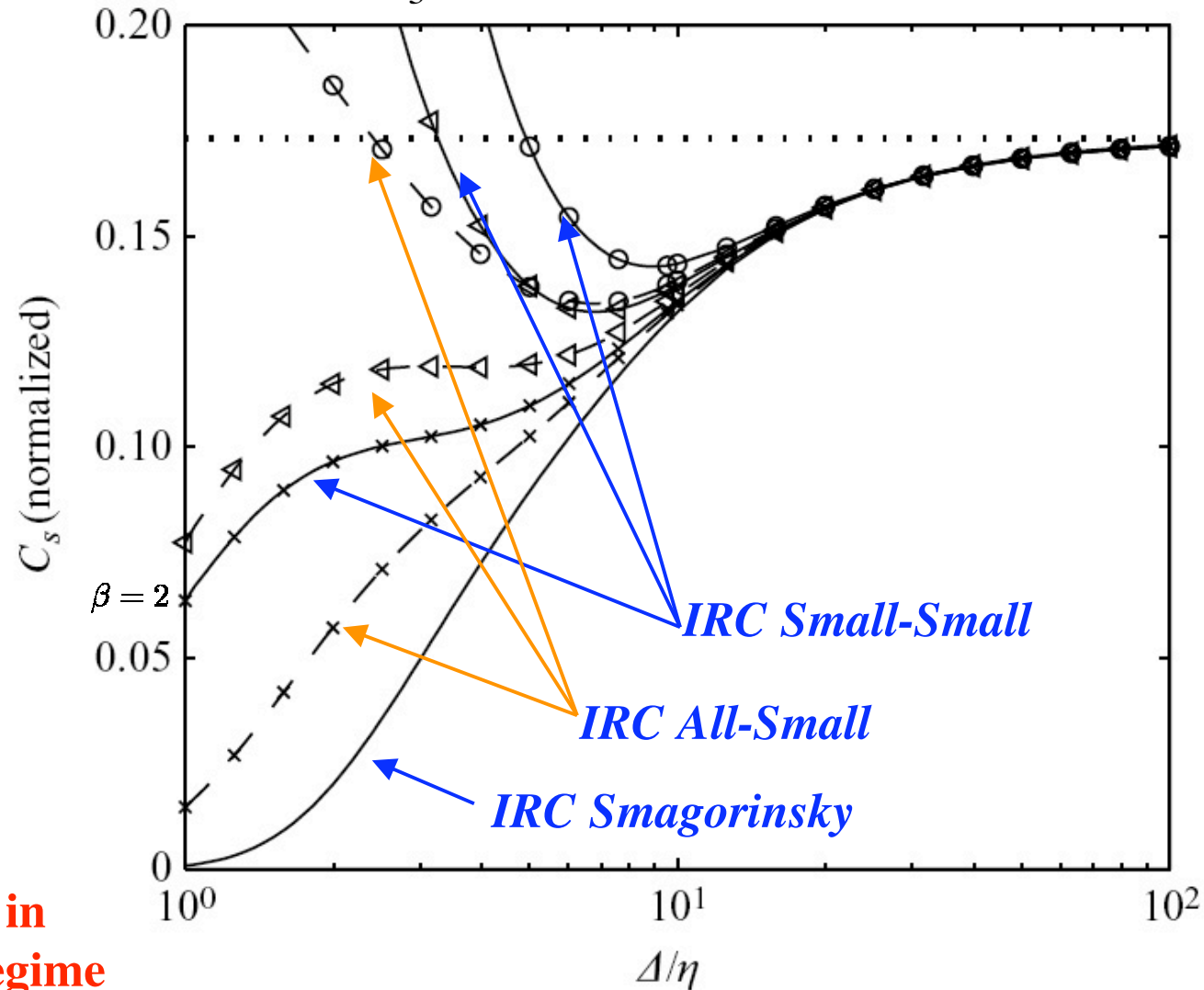
$$m_{ij} = -[2C_{a1}^2 \Delta^2 |\bar{S}| \bar{S}'_{ij}]'$$

- IRC variant

$$\begin{aligned}
 C_{a1} &= C_{s,\infty} \frac{\Phi^{-1/4} \Psi_1^{-1/2}}{\gamma^{1/3} \gamma_1^{2/3} \sqrt{1 - \beta^{4/3}}} \sqrt{1 - \left(\frac{\gamma L}{C_{s,\infty} \Delta} \right)^{4/3} \frac{1}{Re_L} \Phi} \\
 &= C_{s,\infty} \frac{\Phi^{-1/4} \Psi_1^{-1/2}}{\gamma^{1/3} \gamma_1^{2/3} \sqrt{1 - \beta^{4/3}}} \sqrt{1 - \left(\frac{\gamma \eta}{C_{s,\infty} \Delta} \right)^{4/3} \Phi}.
 \end{aligned}$$

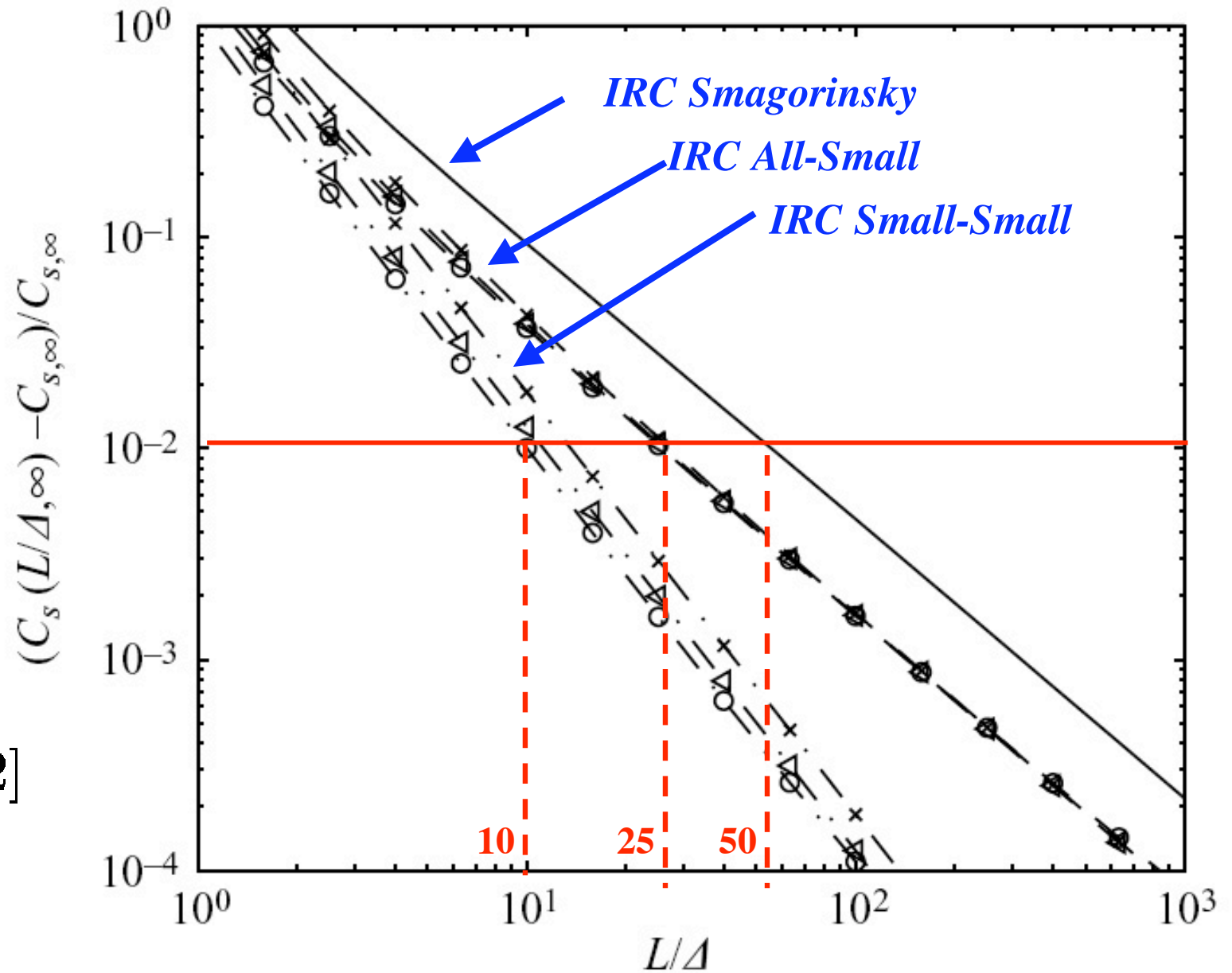
IRC constant behavior

$L/\Delta \gg 1$ Sharp cutoff filters
 $\beta = \frac{4}{3}$



**Transition in
 Quasi-DNS regime**

Deviation from asymptotic value



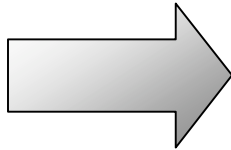
$\beta \in [1, 2]$

Asymptotic behavior

$$L/\Delta \gg 1 \quad Re_L = +\infty$$

$$\delta\Phi = \frac{4}{3} \frac{\int_0^{\pi L/\Delta} x^{1/3} (1 - f_L(x)) dx}{(\pi L/\Delta)^{4/3}}$$

$$\begin{aligned} \delta\Phi &\approx \frac{4}{3} \frac{\int_0^{\pi L/\ell} x^{1/3} (1 - f_L(x)) dx + \int_{\pi L/\ell}^{\pi L/\Delta} C x^{1/3} x^{-p} dx}{(\pi L/\Delta)^{4/3}} \\ &= C' (L/\Delta)^{-4/3} + C'' (L/\Delta)^{-p}, \end{aligned}$$



$$\delta C_s \sim (L/\Delta)^{-\min(4/3, p)}$$

$$\delta C_{a1} \sim (L/\Delta)^{-\min(4/3, p)}$$

$$\delta C_{s1} \sim (L/\Delta)^{-p}.$$

$$p \in [1, 2]$$

- IRC models not tractable \Leftrightarrow approximations are necessary
- Introducing 2 non-dimensional parameter (e.g. Smagorinsky model)

$$R = \frac{\mathbf{v}_t}{\mathbf{v}} \quad Q = \frac{(C_\infty \Delta / \gamma)^2 \sqrt{2 \bar{S}_{ij} \bar{S}_{ij}}}{\mathbf{v}}$$

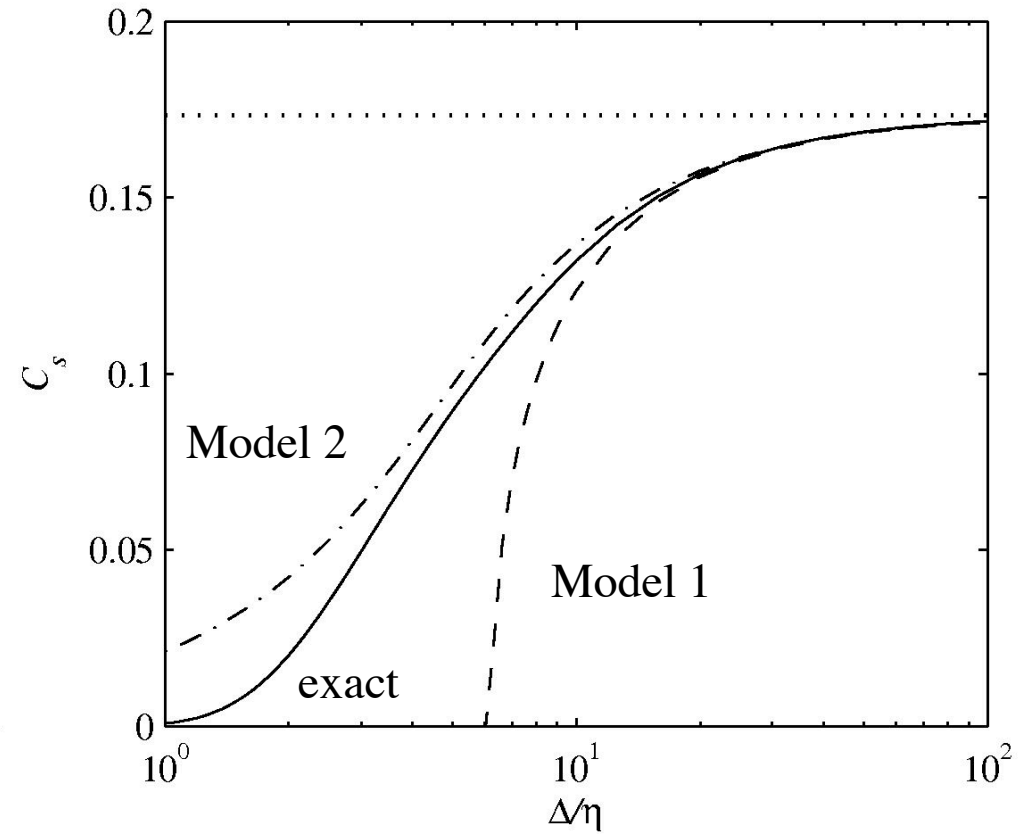
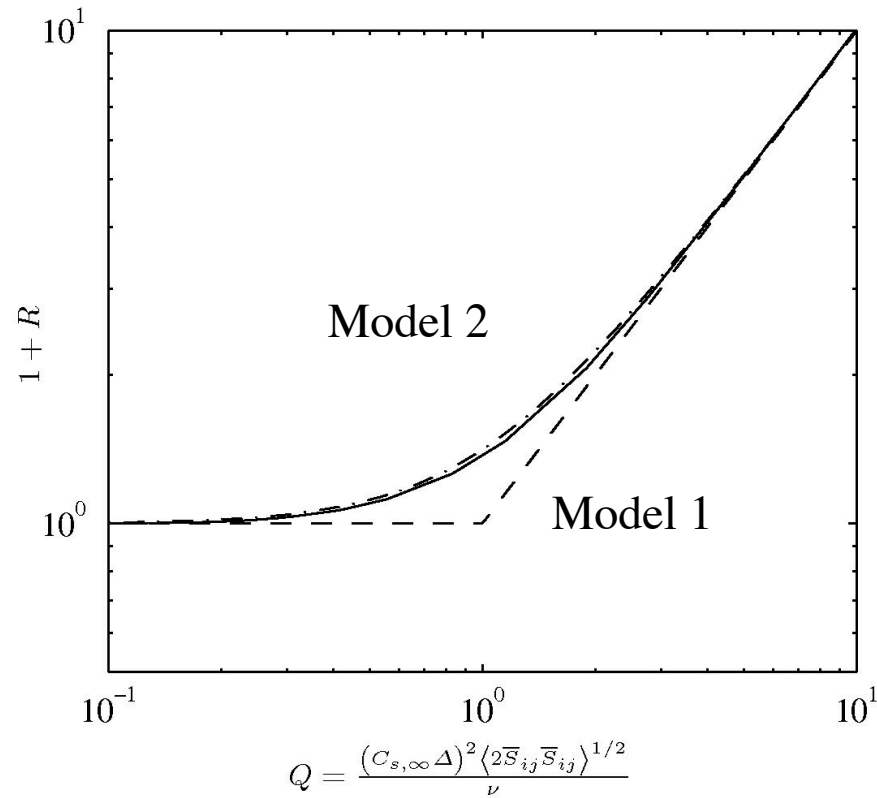
A closure is expressed as $F(R, Q) = 0$

Original model $R = Q \Leftrightarrow \mathbf{v}_t = \mathbf{v}_{Lilly}$

Model 1: $R = \max(Q - 1, 0) \Leftrightarrow \mathbf{v}_t = \max(\mathbf{v}_{Lilly} - \mathbf{v}, 0)$

Model 2: $(R + 1)^2 - Q^2 - 1 = 0 \Leftrightarrow \mathbf{v}_t = \sqrt{\mathbf{v}_{Lilly}^2 + \mathbf{v}^2} - \mathbf{v}$

Remapping for Smagorinsky model



- Small-Small model (quadratic mapping)

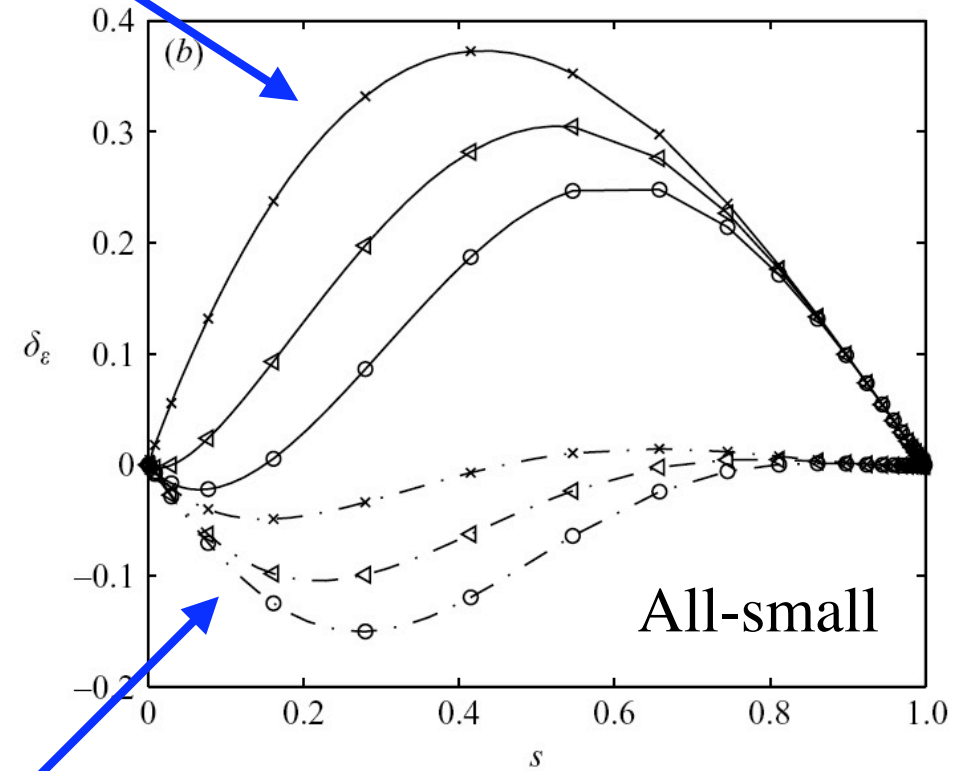
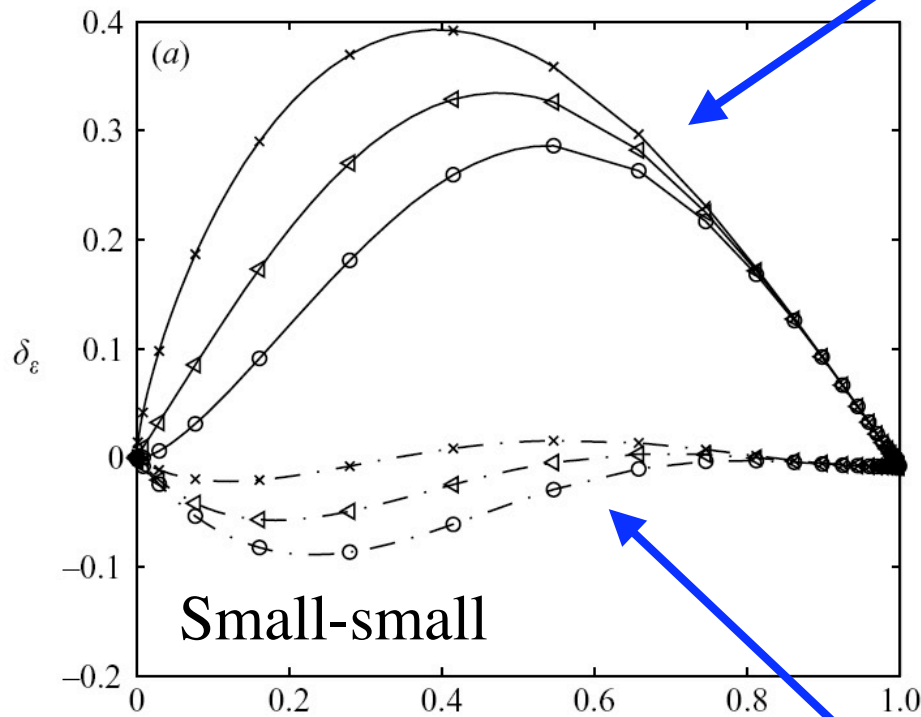
$$v'_{t,s1} = \frac{(\gamma/\gamma_1)^{4/3}}{1 - \beta^{4/3}} \left(\sqrt{\left(\frac{C_{s,\infty}\Delta}{\gamma}\right)^4 \frac{(\gamma/\gamma_1)^{4/3} \langle 2\bar{S}'_{ij}\bar{S}'_{ij} \rangle}{1 - \beta^{4/3}} + v^2} - v \right)$$

- All-Small model (quadratic mapping)

$$v'_{t,a1} = \frac{(\gamma/\gamma_1)^{4/3}}{1 - \beta^{4/3}} \left(\sqrt{(C_{s,\infty}\Delta/\gamma)^4 \langle 2\bar{S}_{ij}\bar{S}_{ij} \rangle + v^2} - v \right)$$

Normalized dissipation error

Non-IRC



IRC approximations

- Inertial-range consistent VMS subgrid models defined
- Asymptotic behavior may be complex !
- Recovery of DNS is model-dependent
- IRC approach can be applied to any VMS models
- IRC analysis may be extended considering numerical dissipation

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Thank you!