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Finite element construction

Scientific Computing Winter 2016/2017

Lecture 18

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Recap (Finite Elements)

Heat conduction revisited: Derivation of weak formulation

+– Sobolev space theory provides the necessary framework to formulate existence and uniqueness of solutions of PDEs.

- ▶ Heat conduction equation with homogeneous Dirichlet boundary conditions:

$$\begin{aligned} -\nabla \cdot \lambda \nabla u &= f \text{ in } \Omega \\ u &= 0 \text{ on } \partial\Omega \end{aligned}$$

Multiply and integrate with an arbitrary *test function* from $C_0^\infty(\Omega)$:

$$\begin{aligned} - \int_{\Omega} \nabla \cdot \lambda \nabla u v \, dx &= \int_{\Omega} f v \, dx \\ \int_{\Omega} \lambda \nabla u \nabla v \, dx &= \int_{\Omega} f v \, dx \end{aligned}$$

Weak formulation of homogeneous Dirichlet problem

- ▶ Search $u \in H_0^1(\Omega)$ such that

$$\int_{\Omega} \lambda \nabla u \nabla v \, dx = \int_{\Omega} f v \, dx \quad \forall v \in H_0^1(\Omega)$$

- ▶ Then,

$$a(u, v) := \int_{\Omega} \lambda \nabla u \nabla v \, dx$$

is a self-adjoint bilinear form defined on the Hilbert space $H_0^1(\Omega)$.

It is bounded due to Cauchy-Schwarz:

$$|a(u, v)| = |\lambda| \left| \int_{\Omega} \nabla u \nabla v \, dx \right| \leq \|u\|_{H_0^1(\Omega)} \cdot \|v\|_{H_0^1(\Omega)}$$

- ▶ $f(v) = \int_{\Omega} f v \, dx$ is a linear functional on $H_0^1(\Omega)$. For Hilbert spaces V the dual space V' (the space of linear functionals) can be identified with the space itself.

The Lax-Milgram lemma

Let V be a Hilbert space. Let $a : V \times V \rightarrow \mathbb{R}$ be a self-adjoint bilinear form, and f a linear functional on V . Assume a is coercive, i.e.

$$\exists \alpha > 0 : \forall u \in V, a(u, u) \geq \alpha \|u\|_V^2.$$

Then the problem: find $u \in V$ such that

$$a(u, v) = f(v) \quad \forall v \in V$$

admits one and only one solution with an a priori estimate

$$\|u\|_V \leq \frac{1}{\alpha} \|f\|_{V'}$$

Heat conduction revisited

Let $\lambda > 0$. Then the weak formulation of the heat conduction problem: search $u \in H_0^1(\Omega)$ such that

$$\int_{\Omega} \lambda \nabla u \nabla v \, dx = \int_{\Omega} f v \, dx \quad \forall v \in H_0^1(\Omega)$$

has an unique solution.

Proof: $a(u, v)$ is coercive:

$$a(u, v) = \int_{\Omega} \lambda \nabla u \nabla u \, dx = \lambda \|u\|_{H_0^1(\Omega)}^2$$

□

Weak formulation of inhomogeneous Dirichlet problem

$$\begin{aligned} -\nabla \cdot \lambda \nabla u &= f \text{ in } \Omega \\ u &= g \text{ on } \partial\Omega \end{aligned}$$

If g is smooth enough, there exists a *lifting* $u_g \in H^1(\Omega)$ such that $u_g|_{\partial\Omega} = g$. Then, we can re-formulate:

$$\begin{aligned} -\nabla \cdot \lambda \nabla (u - u_g) &= f + \nabla \cdot \lambda \nabla u_g \text{ in } \Omega \\ u - u_g &= 0 \text{ on } \partial\Omega \end{aligned}$$

- Search $u \in H^1(\Omega)$ such that

$$\begin{aligned} u &= u_g + \phi \\ \int_{\Omega} \lambda \nabla \phi \nabla v \, dx &= \int_{\Omega} f v \, dx + \int_{\Omega} \lambda \nabla u_g \nabla v \, dx \quad \forall v \in H_0^1(\Omega) \end{aligned}$$

Here, necessarily, $\phi \in H_0^1(\Omega)$ and we can apply the theory for the homogeneous Dirichlet problem.

Weak formulation of Robin problem

$$\begin{aligned} -\nabla \cdot \lambda \nabla u &= f \text{ in } \Omega \\ \lambda \nabla u \cdot \mathbf{n} + \alpha(u - g) &= 0 \text{ on } \partial\Omega \end{aligned}$$

Multiply and integrate with an arbitrary *test function* from $C_c^\infty(\Omega)$:

$$\begin{aligned} - \int_{\Omega} (\nabla \cdot \lambda \nabla u) v \, dx &= \int_{\Omega} f v \, dx \\ \int_{\Omega} \lambda \nabla u \nabla v \, dx + \int_{\partial\Omega} (\lambda \nabla u \cdot \mathbf{n}) v \, ds &= \int_{\Omega} f v \, dx \\ \int_{\Omega} \lambda \nabla u \nabla v \, dx + \int_{\partial\Omega} \alpha u v \, ds &= \int_{\Omega} f v \, dx + \int_{\partial\Omega} \alpha g v \, ds \end{aligned}$$

Weak formulation of Robin problem II

- ▶ Let

$$a^R(u, v) := \int_{\Omega} \lambda \nabla u \nabla v \, dx + \int_{\partial\Omega} \alpha uv \, ds$$
$$f^R(v) := \int_{\Omega} fv \, dx + \int_{\partial\Omega} \alpha gv \, ds$$

The integrals over $\partial\Omega$ must be understood in the sense of the trace space $H^{\frac{1}{2}}(\partial\Omega)$.

- ▶ Search $u \in H^1(\Omega)$ such that

$$a^R(u, v) = f^R(v) \quad \forall v \in H^1(\Omega)$$

- ▶ If $\lambda > 0$ and $\alpha > 0$ then $a^R(u, v)$ is coercive.

Neumann boundary conditions

Homogeneous Neumann:

$$\lambda \nabla u \cdot \mathbf{n} = 0 \text{ on } \partial\Omega$$

Inhomogeneous Neumann:

$$\lambda \nabla u \cdot \mathbf{n} = g \text{ on } \partial\Omega$$

Weak formulation:

- ▶ Search $u \in H^1(\Omega)$ such that

$$\int_{\omega} \nabla u \nabla v dx = \int_{\partial\Omega} g v ds \quad \forall v \in H^1(\Omega)$$

Not coercive due to the fact that we can add an arbitrary constant to u and $a(u, u)$ stays the same!

Further discussion on boundary conditions

- ▶ Mixed boundary conditions:
One can have different boundary conditions on different parts of the boundary. In particular, if Dirichlet or Robin boundary conditions are applied on at least a part of the boundary of measure larger than zero, the bilinear form becomes coercive.
- ▶ Natural boundary conditions: Robin, Neumann
These are imposed in a “natural” way in the weak formulation
- ▶ Essential boundary conditions: Dirichlet
Explicitly imposed on the function space
- ▶ Coefficients $\lambda, \alpha \dots$ can be functions.

The Dirichlet penalty method

- ▶ Robin problem: search $u_\alpha \in H^1(\Omega)$ such that

$$\int_{\Omega} \lambda \nabla u_\alpha \nabla v \, dx + \int_{\partial\Omega} \alpha u_\alpha v \, ds = \int_{\Omega} f v \, dx + \int_{\partial\Omega} \alpha g v \, ds \quad \forall v \in H^1(\Omega)$$

- ▶ Dirichlet problem: search $u \in H^1(\Omega)$ such that

$$u = u_g + \phi \quad \text{where } u_g|_{\partial\Omega} = g$$
$$\int_{\Omega} \lambda \nabla \phi \nabla v \, dx = \int_{\Omega} f v \, dx + \int_{\Omega} \lambda \nabla u_g \nabla v \quad \forall v \in H_0^1(\Omega)$$

- ▶ Penalty limit:

$$\lim_{\alpha \rightarrow \infty} u_\alpha = u$$

- ▶ Formally, the convergence rate is quite low
- ▶ Implementing Dirichlet boundary conditions directly leads to a number of technical problems
- ▶ Implementing the penalty method is technically much simpler
- ▶ Proper way of handling the parameter leads to exact fulfillment of Dirichlet boundary condition in the floating point precision

The Galerkin method I

- ▶ Weak formulations “live” in Hilbert spaces which essentially are infinite dimensional
- ▶ For computer representations we need finite dimensional approximations
- ▶ The finite volume method provides one possible framework which in many cases is close to physical intuition. However, its error analysis is hard.
- ▶ The Galerkin method and its modifications provide a general scheme for the derivation of finite dimensional approximations

The Galerkin method II

- ▶ Let V be a Hilbert space. Let $a : V \times V \rightarrow \mathbb{R}$ be a self-adjoint bilinear form, and f a linear functional on V . Assume a is coercive with coercivity constant α , and continuity constant γ .
- ▶ Continuous problem: search $u \in V$ such that

$$a(u, v) = f(v) \quad \forall v \in V$$

- ▶ Let $V_h \subset V$ be a finite dimensional subspace of V
- ▶ “Discrete” problem \equiv Galerkin approximation:
Search $u_h \in V_h$ such that

$$a(u_h, v_h) = f(v_h) \quad \forall v_h \in V_h$$

By Lax-Milgram, this problem has a unique solution as well.

Céa's lemma

- ▶ What is the connection between u and u_h ?
- ▶ Let $v_h \in V_h$ be arbitrary. Then

$$\begin{aligned}\alpha \|u - u_h\|^2 &\leq a(u - u_h, u - u_h) \quad (\text{Coercivity}) \\ &= a(u - u_h, u - v_h) + a(u - u_h, v_h - u_h) \\ &= a(u - u_h, u - v_h) \quad (\text{Galerkin Orthogonality}) \\ &\leq \gamma \|u - u_h\| \cdot \|u - v_h\| \quad (\text{Boundedness})\end{aligned}$$

- ▶ As a result

$$\|u - u_h\| \leq \frac{\gamma}{\alpha} \inf_{v_h \in V_h} \|u - v_h\|$$

- ▶ Up to a constant, the error of the Galerkin approximation is the error of the best approximation of the solution in the subspace V_h .

From the Galerkin method to the matrix equation

- ▶ Let $\phi_1 \dots \phi_n$ be a set of basis functions of V_h .
- ▶ Then, we have the representation $u_h = \sum_{j=1}^n u_j \phi_j$
- ▶ In order to search $u_h \in V_h$ such that

$$a(u_h, v_h) = f(v_h) \quad \forall v_h \in V_h$$

it is actually sufficient to require

$$\begin{aligned} a(u_h, \phi_i) &= f(\phi_i) \quad (i = 1 \dots n) \\ a\left(\sum_{j=1}^n u_j \phi_j, \phi_i\right) &= f(\phi_i) \quad (i = 1 \dots n) \\ \sum_{j=1}^n a(\phi_j, \phi_i) u_j &= f(\phi_i) \quad (i = 1 \dots n) \end{aligned}$$

$$AU = F$$

with $A = (a_{ij})$, $a_{ij} = a(\phi_i, \phi_j)$, $F = (f_i)$, $f_i = F(\phi_i)$, $U = (u_i)$.

- ▶ Matrix dimension is $n \times n$. Matrix sparsity ?

Obtaining a finite dimensional subspace

- ▶ Let $\Omega = (a, b) \subset \mathbb{R}^1$
- ▶ Let $a(u, v) = \int_{\Omega} \lambda(x) \nabla u \nabla v dx$.
- ▶ Analysis I provides a finite dimensional subspace: the space of sin/cos functions up to a certain frequency \Rightarrow *spectral method*
- ▶ Ansatz functions have global support \Rightarrow full $n \times n$ matrix
- ▶ OTOH: rather fast convergence for smooth data
- ▶ Generalization to higher dimensions possible
- ▶ Big problem in irregular domains: we need the eigenfunction basis of some operator. . .
- ▶ Spectral methods are successful in cases where one has regular geometry structures and smooth/constant coefficients – e.g. “Spectral Einstein Code”

The finite element idea

- ▶ Choose basis functions with local support. In this case, the matrix becomes sparse, as only integrals of basis function pairs with overlapping support contribute to the matrix.
- ▶ Linear finite elements in $\Omega = (a, b) \subset \mathbb{R}^1$:
- ▶ Partition $a = x_1 \leq x_2 \leq \dots \leq x_n = b$
- ▶ Basis functions (for $i = 1 \dots n$)

$$\phi_i(x) = \begin{cases} \frac{x-x_{i-1}}{x_i-x_{i-1}}, & i > 1, x \in (x_{i-1}, x_i) \\ \frac{x_{i+1}-x}{x_{i+1}-x_i}, & i < n, x \in (x_i, x_{i+1}) \\ 0, & \text{else} \end{cases}$$

- ▶ Any function $u_h \in V_h = \text{span}\{\phi_1 \dots \phi_n\}$ is piecewise linear, and the coefficients in the representation $u_h = \sum_{i=1}^n u_i \phi_i$ are the values $u_h(x_i)$.
- ▶ Fortunately, we are working with a weak formulation, and weak derivatives are well defined !

1D matrix elements

($\lambda = 1$, $x_{i+1} - x_i = h$) - The integrals are nonzero for $i = j, i + 1 = j, i - 1 = j$

Let $j = i + 1$

$$\begin{aligned} a_{ij} = a(\phi_i, \phi_{i+1}) &= \int_{\Omega} \nabla \phi_i \nabla \phi_j dx = \int_{x_i}^{x_{i+1}} \nabla \phi_i \nabla \phi_j dx = - \int_{x_i}^{x_{i+1}} \frac{1}{h^2} dx \\ &= \frac{1}{h} dx \end{aligned}$$

Similarly, $a(\phi_i, \phi_{i-1}) = -\frac{1}{h}$

For $1 < i < N$:

$$\begin{aligned} a_{ii} = a(\phi_i, \phi_i) &= \int_{\Omega} \nabla \phi_i \nabla \phi_i dx = \int_{x_{i-1}}^{x_{i+1}} \nabla \phi_i \nabla \phi_i dx = \int_{x_{i-1}}^{x_{i+1}} \frac{1}{h^2} dx \\ &= \frac{2}{h} dx \end{aligned}$$


For $i = 1$ or $i = N$, $a(\phi_i, \phi_i) = \frac{1}{h}$

1D matrix elements II

Adding the boundary integrals yields

$$A = \begin{pmatrix} \alpha + \frac{1}{h} & -\frac{1}{h} & & & \\ -\frac{1}{h} & \frac{2}{h} & -\frac{1}{h} & & \\ & -\frac{1}{h} & \frac{2}{h} & -\frac{1}{h} & \\ & & \ddots & \ddots & \ddots \\ & & & -\frac{1}{h} & \frac{2}{h} & -\frac{1}{h} \\ & & & & -\frac{1}{h} & \frac{2}{h} & -\frac{1}{h} \\ & & & & & -\frac{1}{h} & \frac{1}{h} + \alpha \end{pmatrix}$$

... the same matrix as for the finite volume method...



Finite Elements in higher dimensions

(after Ern/Guermond)

Definition of a Finite Element (Ciarlet)

Triplet $\{K, P, \Sigma\}$ where

- ▶ $K \subset \mathbb{R}^d$: compact, connected Lipschitz domain with non-empty interior
- ▶ P : finite dimensional vector space of functions $p : K \rightarrow \mathbb{R}^m$ (mostly, $m = 1, m = d$)
- ▶ $\Sigma = \{\sigma_1 \dots \sigma_s\} \subset \mathcal{L}(P, \mathbb{R})$: set of linear forms defined on P called *local degrees of freedom* such that the mapping

$$\Lambda_\Sigma : P \rightarrow \mathbb{R}^s$$
$$p \mapsto (\sigma_1(p) \dots \sigma_s(p))$$

is bijective, i.e. Σ is a basis of $\mathcal{L}(P, \mathbb{R})$.

Local shape functions

- ▶ Due to bijectivity of Λ_Σ , for any finite element $\{K, P, \Sigma\}$, there exists a basis $\{\theta_1 \dots \theta_s\} \subset P$ such that

$$\sigma_i(\theta_j) = \delta_{ij} \quad (1 \leq i, j \leq s)$$

- ▶ Elements of such a basis are called *local shape functions*

Unisolvence

- ▶ Bijectivity of Λ_{Σ} is equivalent to the condition

$$\forall (\alpha_1 \dots \alpha_s) \in \mathbb{R}^s \exists! p \in P \text{ such that } \sigma_i(p) = \alpha_i \quad (1 \leq i \leq s)$$

i.e. for any given tuple of values $a = (\alpha_1 \dots \alpha_s)$ there is a unique polynomial $p \in P$ such that $\Lambda_{\Sigma}(p) = a$.

- ▶ Equivalent to *unisolvence*:

$$\begin{cases} \dim P = |\Sigma| = s \\ \forall p \in P : \sigma_i(p) = 0 (i = 1 \dots s) \Rightarrow p = 0 \end{cases}$$

Lagrange finite elements

- ▶ A finite element $\{K, P, \Sigma\}$ is called *Lagrange* finite element (or *nodal* finite element) if there exist a set of points $\{a_1 \dots a_s\} \subset K$ such that

$$\sigma_i(p) = p(a_i) \quad 1 \leq i \leq s$$

- ▶ $\{a_1 \dots a_s\}$: *nodes* of the finite element
- ▶ *nodal basis: $\{\theta_1 \dots \theta_s\} \subset P$ such that

$$\theta_j(a_i) = \delta_{ij} \quad (1 \leq i, j \leq s)$$

Hermite finite elements

- ▶ All or a part of degrees of freedoms defined by derivatives of p in some points

Local interpolation operator

- ▶ Let $\{K, P, \Sigma\}$ be a finite element with shape function bases $\{\theta_1 \dots \theta_s\}$. Let $V(K)$ be a normed vector space of functions $v : K \rightarrow \mathbb{R}^m$ such that
 - ▶ $P \subset V(K)$
 - ▶ The linear forms in Σ can be extended to be defined on $V(K)$
- ▶ *local interpolation operator*

$$\mathcal{I}_K : V(K) \rightarrow P$$

$$v \mapsto \sum_{i=1}^s \sigma_i(v) \theta_i$$

- ▶ P is invariant under the action of \mathcal{I}_K , i.e. $\forall p \in P, \mathcal{I}_K(p) = p$:
 - ▶ Let $p = \sum_{j=1}^s \alpha_j \theta_j$ Then,

$$\begin{aligned} \mathcal{I}_K(p) &= \sum_{i=1}^s \sigma_i(p) \theta_i = \sum_{i=1}^s \sum_{j=1}^s \alpha_j \sigma_i(\theta_j) \theta_i \\ &= \sum_{i=1}^s \sum_{j=1}^s \alpha_j \delta_{ij} \theta_i = \sum_{j=1}^s \alpha_j \theta_j \end{aligned}$$

Local Lagrange interpolation operator

- ▶ Let $V(K) = (C^0(K))^m$

$$\mathcal{I}_K : V(K) \rightarrow P$$

$$v \mapsto I_K v = \sum_{i=1}^s v(\mathbf{a}_i) \theta_i$$

Simplices

- ▶ Let $\{a_0 \dots a_d\} \subset \mathbb{R}^d$ such that the d vectors $a_1 - a_0 \dots a_d - a_0$ are linearly independent. Then the convex hull K of $a_0 \dots a_d$ is called *simplex*, and $a_0 \dots a_d$ are called *vertices* of the simplex.
- ▶ *Unit simplex*: $a_0 = (0 \dots 0)$, $a_1 = (0, 1 \dots 0) \dots a_d = (0 \dots 0, 1)$.

$$K = \left\{ x \in \mathbb{R}^d : x_i \geq 0 \ (i = 1 \dots d) \text{ and } \sum_{i=1}^d x_i \leq 1 \right\}$$

- ▶ A general simplex can be defined as an image of the unit simplex under some affine transformation
- ▶ F_i : face of K opposite to a_i
- ▶ \mathbf{n}_i : outward normal to F_i

Barycentric coordinates

- ▶ Let K be a simplex.
- ▶ Functions λ_i ($i = 0 \dots d$):

$$\lambda_i : \mathbb{R}^d \rightarrow \mathbb{R}$$

$$x \mapsto \lambda_i(x) = 1 - \frac{(x - a_i) \cdot \mathbf{n}_i}{(a_j - a_i) \cdot \mathbf{n}_i}$$

where a_j is any vertex of K situated in F_i .

- ▶ For $x \in K$, one has

$$\begin{aligned} 1 - \frac{(x - a_i) \cdot \mathbf{n}_i}{(a_j - a_i) \cdot \mathbf{n}_i} &= \frac{(a_j - a_i) \cdot \mathbf{n}_i - (x - a_i) \cdot \mathbf{n}_i}{(a_j - a_i) \cdot \mathbf{n}_i} \\ &= \frac{(a_j - x) \cdot \mathbf{n}_i}{(a_j - a_i) \cdot \mathbf{n}_i} = \frac{\text{dist}(x, F_i)}{\text{dist}(a_i, F_i)} \\ &= \frac{\text{dist}(x, F_i) |F_i| / d}{\text{dist}(a_i, F_i) |F_i| / d} \\ &= \frac{\text{dist}(x, F_i) |F_i|}{|K|} \end{aligned}$$

i.e. $\lambda_i(x)$ is the ratio of the volume of the simplex $K_i(x)$ made up of x and the vertices of F_i to the volume of K .

Barycentric coordinates II

- ▶ $\lambda_i(a_j) = \delta_{ij}$
- ▶ $\lambda_i(x) = 0 \quad \forall x \in F_i$
- ▶ $\sum_{i=0}^d \lambda_i(x) = 1 \quad \forall x \in \mathbb{R}^d$
(just sum up the volumes)
- ▶ $\sum_{i=0}^d \lambda_i(x)(x - a_i) = 0 \quad \forall x \in \mathbb{R}^d$
(due to $\sum \lambda_i(x)x = x$ and $\sum \lambda_i a_i = x$ as the vector of linear coordinate functions)
- ▶ Unit simplex:
 - ▶ $\lambda_0(x) = 1 - \sum_{i=1}^d x_i$
 - ▶ $\lambda_i(x) = x_i$ for $1 \leq i \leq d$

Polynomial space \mathbb{P}_k

- ▶ Space of polynomials in $x_1 \dots x_d$ of total degree $\leq k$ with real coefficients $\alpha_{i_1 \dots i_d}$:

$$\mathbb{P}_k = \left\{ p(x) = \sum_{\substack{0 \leq i_1 \dots i_d \leq k \\ i_1 + \dots + i_d \leq k}} \alpha_{i_1 \dots i_d} x_1^{i_1} \dots x_d^{i_d} \right\}$$

- ▶ Dimension:

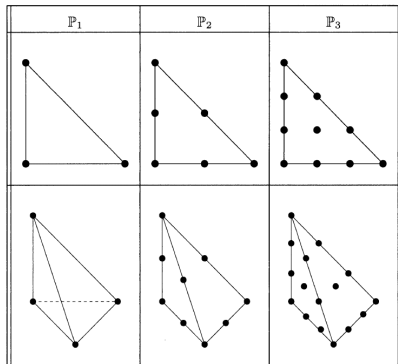
$$\dim \mathbb{P}_k = \binom{d+k}{k} = \begin{cases} k+1, & d=1 \\ \frac{1}{2}(k+1)(k+2), & d=2 \\ \frac{1}{6}(k+1)(k+2)(k+3), & d=3 \end{cases}$$

$$\dim \mathbb{P}_1 = d+1$$

$$\dim \mathbb{P}_2 = \begin{cases} 3, & d=1 \\ 6, & d=2 \\ 10, & d=3 \end{cases}$$

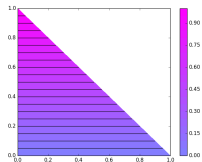
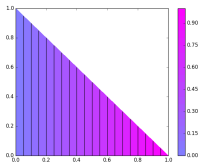
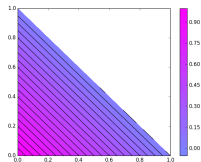
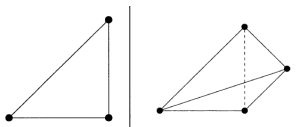
\mathbb{P}_k simplex finite elements

- ▶ K : simplex spanned by $a_0 \dots a_d$ in \mathbb{R}^d
- ▶ $P = \mathbb{P}_k$, such that $s = \dim P_k$
- ▶ For $0 \leq i_0 \dots i_d \leq k$, $i_0 + \dots + i_d = k$, let the set of nodes be defined by the points $a_{i_1 \dots i_d; k}$ with barycentric coordinates $(\frac{i_0}{k} \dots \frac{i_d}{k})$.
Define Σ by $\sigma_{i_1 \dots i_d; k}(p) = p(a_{i_1 \dots i_d; k})$.



\mathbb{P}_1 simplex finite elements

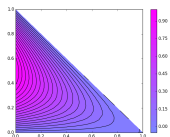
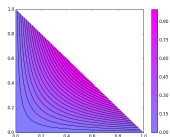
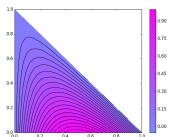
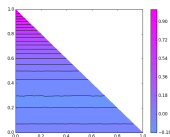
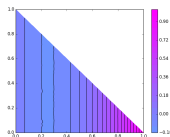
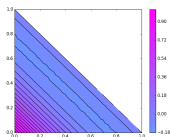
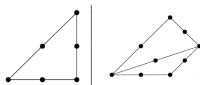
- ▶ K : simplex spanned by $a_0 \dots a_d$ in \mathbb{R}^d
- ▶ $P = \mathbb{P}_1$, such that $s = d + 1$
- ▶ Nodes \equiv vertices
- ▶ Basis functions \equiv barycentric coordinates



\mathbb{P}_2 simplex finite elements

- ▶ K : simplex spanned by $a_0 \dots a_d$ in \mathbb{R}^d
- ▶ $P = \mathbb{P}_2$, Nodes \equiv vertices + edge midpoints
- ▶ Basis functions:

$$\lambda_i(2\lambda_i - 1), (0 \leq i \leq d); \quad 4\lambda_i\lambda_j, \quad (0 \leq i < j \leq d) \quad (\text{"edge bubbles"})$$



Cuboids

- ▶ Given intervals $I_i = [c_i, d_i]$, $i = 1 \dots d$ such that $c_i < d_i$.
- ▶ *Cuboid*:

$$K = \prod_{i=1}^d [c_i, d_i]$$

- ▶ Local coordinate vector $(t_1 \dots t_d) \in [0, 1]^d$
- ▶ Unique representation of $x \in K$: $x_i = c_i + t_i(d_i - c_i)$ for $i = 1 \dots d$.
- ▶ Bijective mapping $[0, 1]^d \rightarrow K$.

Polynomial space \mathbb{Q}_k

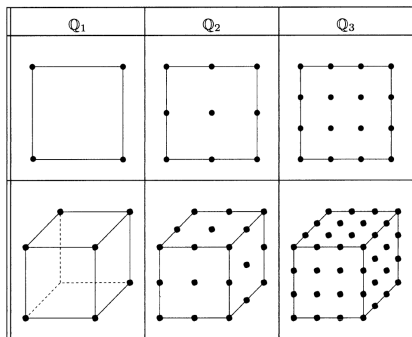
- ▶ Space of polynomials of degree at most k in each variable
- ▶ $d = 1 \Rightarrow \mathbb{Q}_k = \mathbb{P}_k$
- ▶ $d > 1$:

$$\mathbb{Q}_k = \left\{ p(x) = \sum_{0 \leq i_1 \dots i_d \leq k} \alpha_{i_1 \dots i_d} x_1^{i_1} \dots x_d^{i_d} \right\}$$

- ▶ $\dim \mathbb{Q}_k = (k + 1)^d$

\mathbb{Q}_k cuboid finite elements

- ▶ K : cuboid spanned by intervals $[c_i, d_i]$, $i = 1 \dots d$
- ▶ $P = \mathbb{Q}_k$
- ▶ For $0 \leq i_0 \dots i_d \leq k$, let the set of nodes be defined by the points $a_{i_1 \dots i_d; k}$ with local coordinates $(\frac{i_0}{k} \dots \frac{i_d}{k})$.
Define Σ by $\sigma_{i_1 \dots i_d; k}(p) = p(a_{i_1 \dots i_d; k})$.



General finite elements

- ▶ Simplicial finite elements can be defined on triangulations of polygonal domains. During the course we will stick to this case.
- ▶ A curved domain Ω may be approximated by a polygonal domain Ω_h which is then triangulated. During the course, we will ignore this difference.
- ▶ As we have seen, more general elements are possible: cuboids, but and $T_m|_{\hat{F}} = T_n|_{\hat{F}}$ also prismatic elements etc.
- ▶ Curved geometries are possible. Isoparametric finite elements use and $T_m|_{\hat{F}} = T_n|_{\hat{F}}$ the polynomial space to define a mapping of some polyhedral reference element to an element with curved boundary

Conformal triangulations

- ▶ Let \mathcal{T}_h be a subdivision of the polygonal domain $\Omega \subset \mathbb{R}^d$ into non-intersecting compact simplices K_m , $m = 1 \dots n_e$:

$$\bar{\Omega} = \bigcup_{m=1}^{n_e} K_m$$

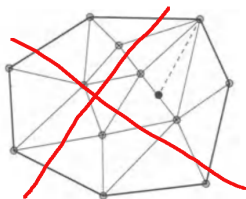
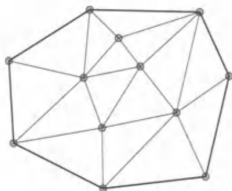
- ▶ Each simplex can be seen as the image of a affine transformation of a reference (e.g. unit) simplex \hat{K} :

$$K_m = T_m(\hat{K})$$

- ▶ We assume that it is conformal, i.e. if K_m, K_n have a $d - 1$ dimensional intersection $F = K_m \cap K_n$, then there is a face \hat{F} of \hat{K} and renumberings of the vertices of K_n, K_m such that $F = T_m(\hat{F}) = T_n(\hat{F})$ and $T_m|_{\hat{F}} = T_n|_{\hat{F}}$

Conformal triangulations II

- ▶ $d = 1$: Each intersection $F = K_m \cap K_n$ is either empty or a common vertex
- ▶ $d = 2$: Each intersection $F = K_m \cap K_n$ is either empty or a common vertex or a common edge



- ▶ $d = 3$: Each intersection $F = K_m \cap K_n$ is either empty or a common vertex or a common edge or a common face
- ▶ Triangulations corresponding to simplicial complexes are conformal
- ▶ Delaunay triangulations are conformal

Reference finite element

- ▶ Let $\{\widehat{P}, \widehat{K}, \widehat{\Sigma}\}$ be a fixed finite element
- ▶ Let T_K be some affine transformation and $K = T_K(\widehat{K})$
- ▶ There is a linear bijective mapping ψ_K between functions on K and functions on \widehat{K} :

$$\begin{aligned}\psi_K : V(K) &\rightarrow V(\widehat{K}) \\ f &\mapsto f \circ T_K\end{aligned}$$

- ▶ Let
 - ▶ $K = T_K(\widehat{K})$
 - ▶ $P_K = \{\psi_K^{-1}(\widehat{p}); \widehat{p} \in \widehat{P}\}$,
 - ▶ $\Sigma_K = \{\sigma_{K,i}, i = 1 \dots s : \sigma_{K,i}(p) = \widehat{\sigma}_i(\psi_K(p))\}$ Then $\{K, P_K, \Sigma_K\}$ is a finite element.

Commutativity of interpolation and reference mapping

► $\mathcal{I}_{\hat{K}} \circ \psi_K = \psi_K \circ \mathcal{I}_K,$

i.e. the following diagram is commutative:

$$\begin{array}{ccc} V(K) & \xrightarrow{\psi_K} & V(\hat{K}) \\ \downarrow \mathcal{I}_K & & \downarrow \mathcal{I}_{\hat{K}} \\ P_K & \xrightarrow{\psi_K} & P_{\hat{K}} \end{array}$$

Global interpolation operator \mathcal{I}_h

- ▶ Let $\{K, P_K, \Sigma_K\}_{K \in \mathcal{T}_h}$ be a triangulation of Ω .
- ▶ Domain:

$$D(\mathcal{I}_h) = \{v \in (L^1(\Omega))^m \text{ such that } \forall K \in \mathcal{T}_h, v|_K \in V(K)\}$$

- ▶ For all $v \in D(\mathcal{I}_h)$, define $\mathcal{I}_h v$ via

$$\mathcal{I}_h v|_K = \mathcal{I}_K(v|_K) = \sum_{i=1}^s \sigma_{K,i}(v|_K) \theta_{K,i} \quad \forall K \in \mathcal{T}_h,$$

Assuming $\theta_{K,i} = 0$ outside of K , one can write

$$\mathcal{I}_h v = \sum_{K \in \mathcal{T}_h} \sum_{i=1}^s \sigma_{K,i}(v|_K) \theta_{K,i},$$

mapping $D(\mathcal{I}_h)$ to the *approximation space*

$$W_h = \{v_h \in (L^1(\Omega))^m \text{ such that } \forall K \in \mathcal{T}_h, v_h|_K \in P_K\}$$

H^1 -Conformal approximation using Lagrangian finite elements

- ▶ Let V be a Banach space of functions on Ω . The approximation space W_h is said to be V -conformal if $W_h \subset V$.
- ▶ Non-conformal approximations are possible, we will stick to the conformal case.
- ▶ Conformal subspace of W_h with zero jumps at element faces:

$$V_h = \{v_h \in W_h : \forall n, m, K_m \cap K_n \neq \emptyset \Rightarrow (v_h|_{K_m})_{K_m \cap K_n} = (v_h|_{K_n})_{K_m \cap K_n}\}$$

- ▶ Then: $V_h \subset H^1(\Omega)$.

Happy Holidays!

Next lecture: Jan. 5, 2017