

### 3. Motivating examples

#### ■ Inhomogenous diffusion equations

$$\gamma_\varepsilon(x)\dot{u}(t, x) = \operatorname{div} (A_\varepsilon(x)\nabla u) - f_\varepsilon(x, u(t, x)), \quad t > 0, x \in \Omega$$

(& suitable BC)

$L^2$ -type gradient system  $(\mathbf{X}, \mathcal{E}_\varepsilon, \mathbb{G}_\varepsilon)$  with  $\mathbf{X} = L^2(\Omega)$

$$(\mathbb{G}_\varepsilon v)(x) = \gamma_\varepsilon(x)v(x) \quad \Rightarrow \quad \mathcal{R}_\varepsilon(v) = \int_\Omega \frac{1}{2} \gamma_\varepsilon(x)v(x)^2 dx$$

$$\mathcal{E}_\varepsilon(u) = \int_\Omega \frac{1}{2} \nabla u \cdot A_\varepsilon(x)\nabla u + F_\varepsilon(x, u(x)) dx \quad F_\varepsilon(x, u) = \int_0^u f_\varepsilon(x, w) dw$$

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• Homogenization  $\gamma_\varepsilon(x) = g(x, \frac{x}{\varepsilon})$  and  $A_\varepsilon(x, \frac{x}{\varepsilon})$

**Aim:**  $\mathcal{R}_{\text{eff}}(v) = \int_\Omega \frac{\gamma_{\text{eff}}}{2}v^2 dx$  and  $\mathcal{E}_{\text{eff}}(u) = \int_\Omega \frac{1}{2}\nabla u \cdot A_{\text{eff}}\nabla u + F_{\text{eff}}(u) dx$

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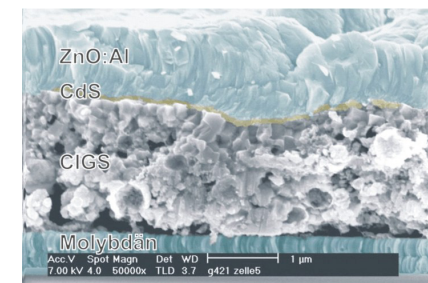
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- Dimension reduction (modeling of active interfaces)

$$x \in ]-1, 1[ \subset \mathbb{R}^1 \quad A_\varepsilon(x) = \begin{cases} \alpha & \text{for } |x| > \varepsilon/2, \\ \beta\varepsilon & \text{for } |x| < \varepsilon/2 \end{cases}$$

$$\mathcal{E}_{\text{eff}}(u) = \int_{-1}^0 \frac{\alpha}{2} u_x^2 dx + \underbrace{\frac{\beta}{2} (u(0^-) - u(0^+))^2}_{\text{gives interface conditions}} + \int_0^1 \frac{\alpha}{2} u_x^2 dx$$



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Fundamental work on evolutionary  $\Gamma$ -convergence:

**Sandier-Serfaty 2004 (Comm. Pure Appl. Math.):**  
*Gamma-convergence of gradient flows with applications to Ginzburg-Landau*

more recent, nicely readable survey:

Serfaty 2011: Gamma-convergence of gradient flows on Hilbert spaces and metric spaces and applications.

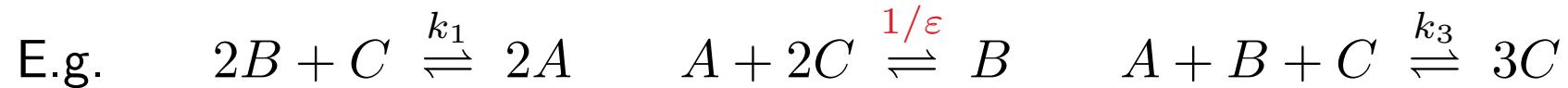
Ginzburg-Landau vortices:  $\psi(t, \cdot) : \Omega \rightarrow \mathbb{C} = \mathbb{R}^2$

$$(GLE) \quad c_\varepsilon \dot{\psi} = \Delta \psi + \frac{1}{\varepsilon^2} (1 - |\psi|^2) \psi \quad \& \quad \text{Neum.BC}$$

(GLE) is induced by the gradient system  $(\mathbf{X}, \mathcal{E}_\varepsilon, \mathcal{R}_\varepsilon)$  with  $\mathbf{X} = H^1(\Omega; \mathbb{C})$ ,  
 $\mathcal{E}_\varepsilon(\psi) = \int_\Omega \frac{1}{2} |\nabla \psi|^2 + \frac{1}{4\varepsilon^2} (1 - |\psi|^2)^2 dx$ , and  $\mathcal{R}_\varepsilon(\dot{\psi}) = \frac{1}{2 \log(1/\varepsilon)} \int_\Omega |\dot{\psi}|^2 dx$

Evol. $\Gamma$ -limit for  $\varepsilon \rightarrow 0 \rightsquigarrow$  ODE for vortex positions

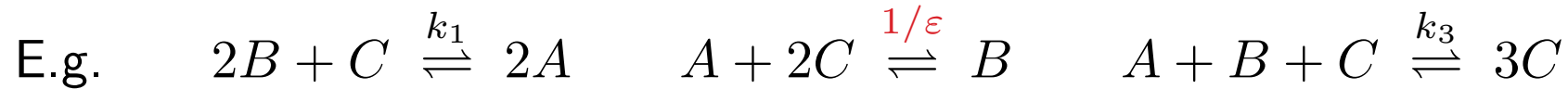
## ■ Chemical reaction systems with detailed balance



Fast reaction versus slow reactions  $k_1, k_3 = O(1)$

$$\begin{pmatrix} \dot{c}_A \\ \dot{c}_B \\ \dot{c}_C \end{pmatrix} = k_1 (c_B^2 c_C - c_A^2) \begin{pmatrix} 2 \\ -2 \\ -1 \end{pmatrix} + \frac{1}{\varepsilon} (c_A c_C^2 - c_B) \begin{pmatrix} -1 \\ 1 \\ -2 \end{pmatrix} + k_3 (c_A c_B c_C - c_C^3) \begin{pmatrix} -1 \\ -1 \\ 2 \end{pmatrix}$$

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Energy = relative entropy  $\mathcal{E}(\mathbf{c}) = \sum_{i=A,B,C} \lambda_{Bz}(c_i) \quad \lambda_{Bz}(z) = z \log z - z + 1$

$$\dot{\mathbf{c}} = -\mathbb{K}(\mathbf{c})D\mathcal{E}(\mathbf{c}) \text{ with } \mathbb{K}_\varepsilon(\mathbf{c}) = \mathbb{K}_{1,3}(\mathbf{c}) + \frac{1}{\varepsilon} \frac{c_A c_C^2 - c_B}{\log(c_A c_C^2 / c_B)} \begin{pmatrix} 1 & -1 & 2 \\ -1 & 1 & -2 \\ 2 & -2 & 4 \end{pmatrix}$$

Gradient system  $([0, \infty[^3, \mathcal{E}, \mathbb{K})$ :  $\mathcal{E}$  indep. of  $\varepsilon$  but  $\mathbb{K}_\varepsilon(\mathbf{c})$

Evol. $\Gamma$ -convergence proved in “Disser, Liero, Zinsl 2016 WIAS preprint 2227”.

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3.2.  $\Gamma$ -convergence for (static) functionals

3.3. An ODE problem

3.4. Homogenization

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## 5. Evolutionary variational inequality (EVI)

## 6. Rate-independent systems (RIS)

$\mathbf{X}$  sep./refl. Banach space and functionals  $\mathcal{J}_\varepsilon : \mathbf{X} \rightarrow \mathbb{R}_\infty = \mathbb{R} \cup \{\infty\}$

Definition (Weak/strong  $\Gamma$ -convergence and Mosco convergence)

**Weak  $\Gamma$ -convergence:**  $\mathcal{J}_\varepsilon \xrightarrow{\Gamma} \mathcal{J}$  if (G1w) and (G2w) hold:

$$(G1w) \quad u_\varepsilon \rightharpoonup u \implies \mathcal{J}(u) \leq \liminf_{\varepsilon \rightarrow 0} \mathcal{J}_\varepsilon(u_\varepsilon) \quad (\text{liminf estimate})$$

$$(G2w) \quad \forall \hat{u} \exists (\hat{u}_\varepsilon)_\varepsilon: \hat{u}_\varepsilon \rightharpoonup \hat{u} \quad \text{and} \quad \mathcal{J}(\hat{u}) = \lim_{\varepsilon \rightarrow 0} \mathcal{J}_\varepsilon(\hat{u}_\varepsilon) \quad (\text{ex. recovery seq.})$$



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**Mosco convergence**  $\mathcal{J}_\varepsilon \xrightarrow{M} \mathcal{J}$  if  $\mathcal{J}_\varepsilon \xrightarrow{\Gamma} \mathcal{J}$  and  $\mathcal{J}_\varepsilon \xrightarrow{\Gamma} \mathcal{J}$  hold  
(or simply (G1w) and (G2s))

### 3. Motivating examples

The (primal) dissipation potentials  $\mathcal{R}(u, \dot{u})$  is always **convex in  $\dot{u}$** .

The dual dissipation potential  $\mathcal{R}^*$  is always **convex in  $\xi$** .

$$\mathcal{R}^*(u, \xi) := \sup\{ \langle \xi, v \rangle - \mathcal{R}(u, v) \mid v \in \mathbf{X} \}$$

Theorem (Attouch 1984)

*Let  $\mathbf{X}$  be a reflexive Banach space and assume that all  $\mathcal{F}_n : \mathbf{X} \rightarrow \mathbb{R}_\infty$  are proper, convex, equicoercive and that  $(\mathcal{F}_n^*)^*$ . Then,*

$$\mathcal{F}_n \xrightarrow{\Gamma} \mathcal{F} \quad \Longleftrightarrow \quad \mathcal{F}_n^* \xrightarrow{\Gamma} \mathcal{F}^* .$$

(HU specialist are F. Bethke and N. Farchmin)

In particular, we have  $\mathcal{F}_n \xrightarrow{M} \mathcal{F} \Leftrightarrow \mathcal{F}_n^* \xrightarrow{M} \mathcal{F}^*$ .

Easy to remember via the well-known convergence result of linear functional analysis:

$$v_n \rightharpoonup v \text{ and } \xi_n \rightarrow \xi \text{ implies } \langle \xi_n, v_n \rangle \rightarrow \langle \xi, v \rangle$$

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$$\mathbf{X} = \mathbb{R}^2$$

$$\mathcal{E}_\varepsilon(u) = \frac{1}{2}u_1^2 + \frac{1}{2\varepsilon^2}(u_2 - \varepsilon u_1)^2 = \frac{1}{2}u \cdot A_\varepsilon u \quad \text{with } A_\varepsilon = \begin{pmatrix} 2 & -1/\varepsilon \\ -1/\varepsilon & 1/\varepsilon^2 \end{pmatrix}$$

$$\mathcal{R}_\varepsilon(\dot{u}) = \frac{1}{2}\dot{u}_1^2 + \frac{1}{2\varepsilon^\beta}\dot{u}_2^2 = \frac{1}{2}\dot{u} \cdot \mathbb{G}_\varepsilon \dot{u} \quad \text{with } \mathbb{G}_\varepsilon = \begin{pmatrix} 1 & 0 \\ 0 & 1/\varepsilon^\beta \end{pmatrix}$$

ODE reads  $\mathbb{G}_\varepsilon \dot{u}_\varepsilon = -A_\varepsilon u_\varepsilon$  with  $u_\varepsilon(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

Explicit solutions can be calculated for all  $\varepsilon > 0$ . We find, for all  $t \geq 0$ ,

$$\beta \in [0, 2[ : \quad u_\varepsilon(t) \rightarrow \begin{pmatrix} e^{-t} \\ 0 \end{pmatrix} \text{ as } \varepsilon \rightarrow 0$$

$$\beta = 2 : \quad u_\varepsilon(t) \rightarrow \begin{pmatrix} w(t) \\ 0 \end{pmatrix} \text{ as } \varepsilon \rightarrow 0$$

$$\beta > 2 : \quad u_\varepsilon(t) \rightarrow \begin{pmatrix} e^{-2t} \\ 0 \end{pmatrix} \text{ as } \varepsilon \rightarrow 0$$

where  $w(t) = \frac{1}{2\sqrt{5}}((\sqrt{5}+1)e^{-\mu_1 t} + (\sqrt{5}-1)e^{-\mu_2 t})$  with  $\mu_{1,2} = (3 \pm \sqrt{5})/2$

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$$\mathcal{E}_\varepsilon(u) = \frac{1}{2}u_1^2 + \frac{1}{2\varepsilon^2}(u_2 - \varepsilon u_1)^2 = \frac{1}{2}u \cdot A_\varepsilon u \quad \text{with } A_\varepsilon = \begin{pmatrix} 2 & -1/\varepsilon \\ -1/\varepsilon & 1/\varepsilon^2 \end{pmatrix}$$

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**What are the limits of the functionals?**

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$$\mathcal{E}_\varepsilon \xrightarrow{\text{pointwise}} \mathcal{E}_{\text{pw}} : u \mapsto \begin{cases} (\frac{1}{2} + \frac{1}{2})u_1^2 & \text{for } u_2 = 0, \\ \infty & \text{otherwise} \end{cases}$$

$$\mathcal{E}_\varepsilon \xrightarrow{\text{M}} \mathcal{E}_0 : u \mapsto \begin{cases} \frac{1}{2}u_1^2 & \text{for } u_2 = 0, \\ \infty & \text{otherwise} \end{cases} \quad \mathcal{R}_\varepsilon \xrightarrow{\text{M}} \mathcal{R}_0 : v \mapsto \begin{cases} \frac{1}{2}v_1^2 & \text{for } v_2 = 0, \\ \infty & \text{otherwise} \end{cases}$$

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$(\mathbb{R}^2, \mathcal{E}, \mathcal{R}_0)$  gives  $u(t) = \begin{pmatrix} e^{-t} \\ 0 \end{pmatrix}$  and

$(\mathbb{R}^2, \mathcal{E}_{\text{pw}}, \mathcal{R}_0)$  gives  $u(t) = \begin{pmatrix} e^{-2t} \\ 0 \end{pmatrix}$ .

$$\boxed{\beta < 2} \quad (\mathbb{R}^2, \mathcal{E}_\varepsilon, \mathcal{R}_\varepsilon) \xrightarrow{\text{evol}} (\mathbb{R}^2, \mathcal{E}, \mathcal{R}_0)$$

$\boxed{\beta = 2}$  **no** evolutionary  $\Gamma$  convergence

$$\boxed{\beta > 2} \quad (\mathbb{R}^2, \mathcal{E}_\varepsilon, \mathcal{R}_\varepsilon) \xrightarrow{\text{evol}} (\mathbb{R}^2, \mathcal{E}_{\text{pw}}, \mathcal{R}_0)$$



### 3. Motivating examples

$$\mathcal{E}_\varepsilon(u) = \frac{1}{2}u \cdot A_\varepsilon u, \quad \mathcal{R}_\varepsilon(\dot{u}) = \frac{1}{2}\dot{u} \cdot \mathbb{G}_\varepsilon \dot{u} \quad A_\varepsilon = \begin{pmatrix} 2 & -1/\varepsilon \\ -1/\varepsilon & 1/\varepsilon^2 \end{pmatrix}, \quad \mathbb{G}_\varepsilon = \begin{pmatrix} 1 & 0 \\ 0 & 1/\varepsilon^\beta \end{pmatrix}$$

Reason for non-convergence is seen via the energy-dissipation relation

$$\frac{d}{dt} \mathcal{E}_\varepsilon(u_\varepsilon(t)) = \langle D\mathcal{E}_\varepsilon(u_\varepsilon), \dot{u}_\varepsilon \rangle = -\langle \mathbb{G}_\varepsilon \dot{u}_\varepsilon, \dot{u}_\varepsilon \rangle = -((\dot{u}_{1,\varepsilon})^2 + (\dot{u}_{2,\varepsilon})^2 / \varepsilon^\beta)$$

$$\rightsquigarrow \mathcal{E}_\varepsilon(u_\varepsilon(T)) + \int_0^T 2\mathcal{R}_\varepsilon(\dot{u}_\varepsilon(t)) dt = \mathcal{E}_\varepsilon(u_\varepsilon(0)) = 1 \text{ (finite, indep. of } \varepsilon)$$

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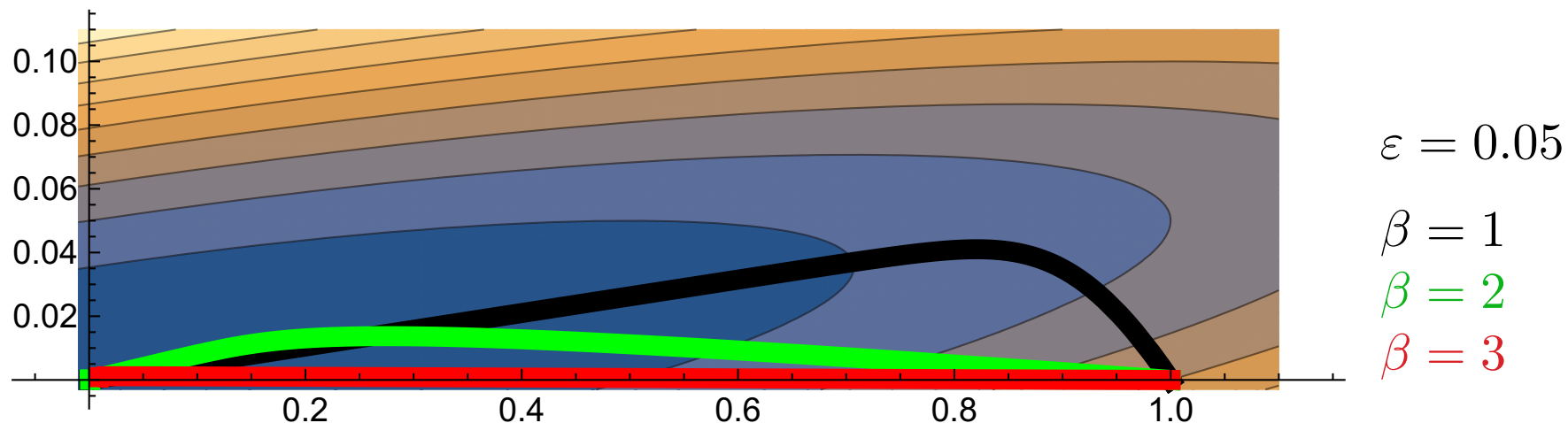
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- Energy landscape  $\mathcal{E}_\varepsilon(u)$  wants to have  $u_2 \approx \varepsilon u_1 \approx \varepsilon e^{-\lambda t}$
- Dissipation  $1 \geq \int_0^T 2\mathcal{R}_\varepsilon(\dot{u}_\varepsilon(t)) dt \geq \int_0^T (\dot{u}_{2,\varepsilon}(t))^2 / \varepsilon^\beta dt$ .

$\beta > 2$ : “dissipation” doesn’t allow solutions to move away from  $u_2 \equiv 0$ .



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We consider **one-dimensional homogenization** of a parabolic equation on  $x \in \Omega = ]0, \ell[$  for  $t > 0$ :

$$c\left(\frac{x}{\varepsilon}\right)\dot{u}(t, x) = \left(a\left(\frac{x}{\varepsilon}\right)u_x(t, x)\right)_x - b\left(\frac{x}{\varepsilon}\right)u(t, x) \quad u_x(t, 0) = 0 = u_x(t, \ell)$$

where  $a, b, c \in L^\infty(\mathbb{R})$  are 1-periodic and are  $\geq c_0 > 0$ .

Family of gradient system  $(L^2(\Omega), \mathcal{E}_\varepsilon, \mathcal{R}_\varepsilon)$  with

$$\mathcal{E}_\varepsilon(u) = \frac{1}{2} \int_{\Omega} a\left(\frac{x}{\varepsilon}\right)u_x(x)^2 + b\left(\frac{x}{\varepsilon}\right)u(x)^2 dx, \quad \mathcal{R}_\varepsilon(v) = \frac{1}{2} \int_{\Omega} c\left(\frac{x}{\varepsilon}\right)v(x)^2 dx$$

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**Aim:** Find  $\mathcal{E}_{\text{eff}}$  and  $\mathcal{R}_{\text{eff}}$  in the form

$$\mathcal{E}_{\text{eff}}(u) = \frac{1}{2} \int_{\Omega} a_{\text{eff}} u_x(x)^2 + b_{\text{eff}} u^2 dx, \quad \mathcal{R}_{\text{eff}}(v) = \frac{1}{2} \int_{\Omega} c_{\text{eff}} v^2 dx$$

$$a_{\text{eff}} = ?$$

$$b_{\text{eff}} = ?$$

$$c_{\text{eff}} = ?$$

#### Quadratic functionals:

$$\Psi_\varepsilon(v) = \frac{1}{2} \int_{\Omega} v(x) \cdot G_\varepsilon(x) v(x) dx \iff \Psi_\varepsilon^*(\xi) = \frac{1}{2} \int_{\Omega} \xi(x) \cdot G_\varepsilon(x)^{-1} \xi(x) dx$$

#### Lemma (One-dimensional homogenization)

Let  $G_\varepsilon(x) = \mathbb{G}(x/\varepsilon)$  with  $0 < c_0 \leq \mathbb{G}(y) \leq C_1$  and  $\mathbb{G}$  1-periodic.

In  $L^2(]x_1, x_2[)$  we have

**weak- $\Gamma$ :**  $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{harm}} : v \mapsto \frac{1}{2} \int_{x_1}^{x_2} v \cdot G_{\text{harm}} v dx$

**strong- $\Gamma$ :**  $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{arith}} : v \mapsto \frac{1}{2} \int_{x_1}^{x_2} v \cdot G_{\text{arithm}} v dx$

with  $G_{\text{harm}} = \left( \int_0^1 \mathbb{G}(y)^{-1} dy \right)^{-1} \leq G_{\text{arithm}} = \int_0^1 \mathbb{G}(y) dy$ .

### 3. Motivating examples

#### Quadratic functionals:

$$\Psi_\varepsilon(v) = \frac{1}{2} \int_{\Omega} v(x) \cdot G_\varepsilon(x) v(x) dx \iff \Psi_\varepsilon^*(x) = \frac{1}{2} \int_{\Omega} \xi(x) \cdot G_\varepsilon(x)^{-1} \xi(x) dx$$

#### Lemma (One-dimensional homogenization)

Let  $G_\varepsilon(x) = \mathbb{G}(x/\varepsilon)$  with  $0 < c_0 \leq \mathbb{G}(y) \leq C_1$  and  $\mathbb{G}$  1-periodic. In  $L^2(]x_1, x_2[)$  we have

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Proof of **weak- $\Gamma$** : Assume  $v_\varepsilon \rightharpoonup v$  in  $L^2(]a, b[)$ .

$$\Psi_\varepsilon(v_\varepsilon) = \frac{1}{2} \int_{x_1}^{x_2} v_\varepsilon \cdot G_\varepsilon v_\varepsilon dx =$$

$$\frac{1}{2} \int_{x_1}^{x_2} \underbrace{(G_\varepsilon v_\varepsilon - G_{\text{ha}} v) \cdot G_\varepsilon^{-1} (G_\varepsilon v_\varepsilon - G_{\text{ha}} v)}_{\geq 0} + \underbrace{2 G_\varepsilon v_\varepsilon \cdot G_\varepsilon^{-1} G_{\text{ha}} v}_{= v_\varepsilon \rightarrow v} - G_{\text{ha}} v \cdot \underbrace{G_\varepsilon^{-1} G_{\text{ha}} v}_{\xrightarrow{*} G_{\text{ha}}^{-1}} dx$$

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$$\begin{aligned} \Psi_\varepsilon(v_\varepsilon) &= \frac{1}{2} \int_{x_1}^{x_2} v_\varepsilon \cdot G_\varepsilon v_\varepsilon dx = \\ &= \frac{1}{2} \int_{x_1}^{x_2} \underbrace{(G_\varepsilon v_\varepsilon - G_{\text{ha}} v)}_{\geq 0} \cdot G_\varepsilon^{-1} (G_\varepsilon v_\varepsilon - G_{\text{ha}} v) + \underbrace{2G_\varepsilon v_\varepsilon \cdot G_\varepsilon^{-1} G_{\text{ha}} v}_{=v_\varepsilon \rightarrow v} - G_{\text{ha}} v \cdot \underbrace{G_\varepsilon^{-1} G_{\text{ha}} v}_{\xrightarrow{*} G_{\text{ha}}^{-1}} dx \end{aligned}$$

Hence,  $\liminf_{\varepsilon \rightarrow 0} \Psi_\varepsilon(v_\varepsilon) \geq \frac{1}{2} \int_{x_1}^{x_2} 0 + 2v \cdot G_{\text{ha}} v - v \cdot G_{\text{ha}} v dx = \Psi_{\text{harm}}(v)$



## Quadratic functionals:

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Given  $\hat{v}$  choose the **recovery sequence**  $\hat{v}_\varepsilon = G_\varepsilon^{-1} G_{\text{ha}} \hat{v} \rightharpoonup \hat{v}$  and first term = 0.

Hence,  $\Psi_\varepsilon(\hat{v}_\varepsilon) = \int_{x_1}^{x_2} 0 + G_{\text{ha}} \hat{v} \cdot G_\varepsilon^{-1} G_{\text{ha}} \hat{v} dx \rightarrow \Psi_{\text{harm}}(\hat{v})$

### 3. Motivating examples

#### Quadratic functionals:

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Proof of **strong- $\Gamma$**  is much simpler:

If  $v_\varepsilon \rightarrow v$  in  $L^2(]a, b[)$ , then

$$\Psi_\varepsilon(v_\varepsilon) = \frac{1}{2} \int_{x_1}^{x_2} v \cdot \underbrace{G_\varepsilon v}_{\rightarrow G_{\text{ar}} v} - 2v \cdot \underbrace{G_\varepsilon(v-v_\varepsilon)}_{\rightarrow 0} + \underbrace{(v-v_\varepsilon) \cdot G_\varepsilon(v-v_\varepsilon)}_{\rightarrow 0} dx \rightarrow \Psi_{\text{ar}}(v)$$

### 3. Motivating examples

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Result is compatible with Attouch's theorem:

$$\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{harm}} \iff \Psi_\varepsilon^* \xrightarrow{\Gamma} \Psi_{\text{harm}}^*$$

For this, simply use  $\text{arith}(\mathbb{G}^{-1}) = \text{harm}(\mathbb{G})^{-1}$ .

### 3. Motivating examples

**One-dimensional homogenization** for parabolic equation on  $x \in \Omega = ]0, \ell[$ :

$$c\left(\frac{x}{\varepsilon}\right)\dot{u}(t, x) = \left(a\left(\frac{x}{\varepsilon}\right)u_x(t, x)\right)_x - b\left(\frac{x}{\varepsilon}\right)u(t, x) \quad u_x(t, 0) = 0 = u_x(t, \ell)$$

where  $a, b, c \in L^\infty(\mathbb{R})$  are 1-periodic and are  $\geq c_0 > 0$ .

Gradient system  $(L^2(\Omega), \mathcal{E}_\varepsilon, \Psi_\varepsilon)$  with

$$\mathcal{E}_\varepsilon(u) = \frac{1}{2} \int_\Omega a\left(\frac{x}{\varepsilon}\right)u_x(x)^2 + b\left(\frac{x}{\varepsilon}\right)u(x)^2 dx, \quad \Psi_\varepsilon(v) = \frac{1}{2} \int_\Omega c\left(\frac{x}{\varepsilon}\right)v(x)^2 dx$$

- $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{harm}}$  or  $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{arith}}$  in the **dynamic space**  $L^2(\Omega)$
- Analogously the energy satisfies in the **energy space**  $H^1(\Omega) \Subset L^2(\Omega)$ 
  - $\mathcal{E}_\varepsilon \xrightarrow{\Gamma} \mathcal{E}_{\text{ha}} : u \mapsto \frac{1}{2} \int_\Omega a_{\text{harm}}u_x^2 + b_{\text{arith}}u^2 dx$  (weakly in  $H^1(\Omega)$ )
  - $\mathcal{E}_\varepsilon \xrightarrow{\Gamma} \mathcal{E}_{\text{ar}} : u \mapsto \frac{1}{2} \int_\Omega a_{\text{arith}}u_x^2 + b_{\text{arith}}u^2 dx$  (strongly in  $H^1(\Omega)$ )

### 3. Motivating examples

**One-dimensional homogenization** for parabolic equation on  $x \in \Omega = ]0, \ell[$ :

$$c\left(\frac{x}{\varepsilon}\right)\dot{u}(t, x) = \left(a\left(\frac{x}{\varepsilon}\right)u_x(t, x)\right)_x - b\left(\frac{x}{\varepsilon}\right)u(t, x) \quad u_x(t, 0) = 0 = u_x(t, \ell)$$

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■  $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{harm}}$  or  $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{arith}}$  in the **dynamic space**  $L^2(\Omega)$

■ Analogously the energy satisfies in the **energy space**  $H^1(\Omega) \subset L^2(\Omega)$

$$\mathcal{E}_\varepsilon \xrightarrow{\Gamma} \mathcal{E}_{\text{ha}} : u \mapsto \frac{1}{2} \int_\Omega a_{\text{harm}}u_x^2 + b_{\text{arith}}u^2 dx \quad (\text{weakly in } H^1(\Omega))$$

$$\mathcal{E}_\varepsilon \xrightarrow{\Gamma} \mathcal{E}_{\text{ar}} : u \mapsto \frac{1}{2} \int_\Omega a_{\text{arith}}u_x^2 + b_{\text{arith}}u^2 dx \quad (\text{strongly in } H^1(\Omega))$$

We will use later:  $\mathcal{E}_\varepsilon \xrightarrow{M} \mathcal{E}_{\text{ha}}$  (Mosco in  $L^2(\Omega)$ )

↗ expected limit eqn  $c_{\text{eff}}u_t = a_{\text{harm}}u_{xx} - b_{\text{arith}}u$  with  $c_{\text{eff}} \in \{c_{\text{harm}}, c_{\text{arith}}\}$

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# Overview

1. Introduction
2. Gradient systems
3. Motivating examples
- 4. Energy-dissipation formulations**
5. Evolutionary variational inequality (EVI)
6. Rate-independent systems (RIS)

# Overview

## 1. Introduction

## 2. Gradient systems

## 3. Motivating examples

## 4. Energy-dissipation formulations

4.1. Equivalent formulations via Legendre transform

4.2. The Sandier-Serfaty approach using EDP

4.3. Choice of GS determines effective equation

4.4. General evolutionary  $\Gamma$ -convergence using EDP

4.5. From viscous to rate-independent friction

## 5. Evolutionary variational inequality (EVI)

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## 6. Rate-independent systems (RIS)

**Legendre-Fenchel theory** for a reflexive Banach space

$\Psi : \mathbf{X} \rightarrow \mathbb{R}_\infty$  proper, convex, lower semicontinuous

Legendre transform  $\Psi^* = \mathcal{L}\Psi : \mathbf{X}^* \rightarrow \mathbb{R}_\infty$  with

$$\Psi^*(\xi) := \sup\{ \langle \xi, v \rangle - \Psi(v) \mid v \in \mathbf{X} \}$$

**Basic properties:**

- $\mathcal{L}(\mathcal{L}\Psi) = \Psi$  or  $\Psi^{**} = \Psi$
- Young-Fenchel estimate:  $\forall v \in \mathbf{X} \quad \forall \xi \in \mathbf{X}^* : \Psi(v) + \Psi^*(\xi) \geq \langle \xi, v \rangle$
- $\Psi(v) = \frac{1}{2} \langle Gv, v \rangle \implies \Psi^*(\xi) = \frac{1}{2} \langle \xi, G^{-1}\xi \rangle$
- $\Psi(v) = \frac{1}{p} \|v\|_{\mathbf{X}}^p \implies \Psi^*(\xi) = \frac{1}{p^*} \|\xi\|_{\mathbf{X}^*}^{p^*} \quad \text{for } 1 < p < \infty, p^* = p/(p-1)$



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$$\Psi^{**} = \Psi \text{ and } \Psi(v) + \Psi^*(\xi) \geq \langle \xi, v \rangle$$

**Subdifferential of convex  $\Psi$**

$$\partial\Psi(v) = \{ \eta \in \mathbf{X}^* \mid \forall w \in \mathbf{X} : \Psi(w) \geq \Psi(v) + \langle \eta, w-v \rangle \} \subset \mathbf{X}^*$$

If  $\Psi \in C^1(\mathbf{X}; \mathbb{R})$  and convex, then  $\partial\Psi(v) = \{D\Psi(v)\}$ .

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If  $\Psi \in C^1(\mathbf{X}; \mathbb{R})$  and convex, then  $\partial\Psi(v) = \{D\Psi(v)\}$ .

Theorem (Fenchel equivalence)

$$(i) \quad \xi \in \partial\Psi(v) \iff (ii) \quad v \in \partial\Psi^*(\xi) \iff (iii) \quad \Psi(v) + \Psi^*(\xi) \leq \langle \xi, v \rangle$$

## 4. Energy-dissipation formulations

$$(i) \quad \xi \in \partial\Psi(v) \iff (ii) \quad v \in \partial\Psi^*(\xi) \iff (iii) \quad \Psi(v) + \Psi^*(\xi) \leq \langle \xi, v \rangle$$

Generalized gradient system  $(\mathbf{X}, \mathcal{E}, \mathcal{R})$

Energy funct.  $\mathcal{E} : [0, T] \times \mathbf{X} \rightarrow \mathbb{R}_\infty$ , dissipation pot.  $\mathcal{R}(u, \cdot) : \mathbf{X} \rightarrow [0, \infty]$

$$(i) \quad 0 \in \partial_{\dot{u}} \mathcal{R}(u(t), \dot{u}(t)) + D\mathcal{E}(t, u(t)) \in \mathbf{X}^* \text{ for a.a. } t \in [0, T]$$

force balance in  $\mathbf{X}^*$

Biot's equation 1954

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force balance in  $\mathbf{X}^*$

Biot's equation 1954

Dual dissipation potential  $\mathcal{R}^*(u, \xi) = \mathcal{L}(\mathcal{R}(u, \cdot))(\xi)$

$$(ii) \quad \dot{u}(t) \in \partial_\xi \mathcal{R}^*(u(t), -D\mathcal{E}(t, u(t))) \in \mathbf{X} \text{ for a.a. } t \in [0, T]$$

rate equation in  $\mathbf{X}$

Onsager's equation 1931

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rate equation in  $\mathbf{X}$

Onsager's equation 1931

$$(iii) \quad \mathcal{R}(u(t), \dot{u}(t)) + \mathcal{R}^*(u(t), -D\mathcal{E}(t, u(t))) \leq \langle -D\mathcal{E}(t, u(t)), \dot{u}(t) \rangle$$

power balance in  $\mathbb{R}$  (equivalent to equality by Young-Fenchel)

De Giorgi's  $(\Psi, \Psi^*)$  formulation 1980

## 4. Energy-dissipation formulations

$$(i) \quad 0 \in \partial_{\dot{u}} \mathcal{R}(u, \dot{u}) + D\mathcal{E}(t, u) \quad (ii) \quad \dot{u} \in \partial_{\xi} \mathcal{R}^*(u, -D\mathcal{E}(t, u))$$

Theorem (Energy-Dissipation Principle (EDP), De Giorgi'80)

Assume that  $\mathcal{E}$  satisfies the **chain rule on  $X$** , then  $u \in W^{1,1}([0, T]; X)$  solves (i) or (ii) if and only if **(EDE)** holds:

$$(EDE) \quad \mathcal{E}(T, u(T)) + \int_0^T \mathcal{R}(u(t), \dot{u}(t)) + \mathcal{R}^*(u(t), -D\mathcal{E}(t, u(t))) dt \leq \mathcal{E}(0, u(0)) + \int_0^T \partial_s \mathcal{E}(s, u(s)) ds$$

Final energy + dissipated energy = initial energy + external work

## 4. Energy-dissipation formulations

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Proof:  $\int_0^T -\langle D\mathcal{E}(t, u), \dot{u} \rangle dt \stackrel{YF}{\leq} \int_0^T \mathcal{R}(u, \dot{u}) + \mathcal{R}^*(u, -D\mathcal{E}) dt$   
 $\stackrel{(EDE)}{\leq} \mathcal{E}(0, u(0)) + \int_0^T \partial_t \mathcal{E} dt - \mathcal{E}(T, u(T)) \stackrel{Ch.Rule}{=} \int_0^T -\langle D\mathcal{E}(t, u), \dot{u} \rangle dt$   
 $\Rightarrow$  all estimates are equalities  $\Rightarrow$  Young-Fenchel estimate is equality a.e. QED

# 4. Energy-dissipation formulations

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Fundamental and more general tool **Chain-Rule Estimate (CR)**

$\mathcal{E} : X \rightarrow \mathbb{R}_{\infty}$  satisfies **CRE**, if

$$\left. \begin{array}{l} u \in W^{1,p}([0, T]; X), \quad \xi \in L^{p'}([0, T]; X^*) \\ \xi(t) \in \partial \mathcal{E}(u(t)) \end{array} \right\} \implies \frac{d}{dt} \mathcal{E}(u(t)) \geq \langle \xi(t), \dot{u}(t) \rangle$$

(e.g. always true for lsc and convex  $\mathcal{E}(\cdot)$ )



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$0 \in \partial_{\dot{u}} \mathcal{R}_\varepsilon(u, \dot{u}) + D\mathcal{E}_\varepsilon(u) \xleftrightarrow{\text{Fenchel}} (\text{EDE}) = \text{Energy-Dissipation Estimate}$

$$(\text{EDE}) \quad \mathcal{E}_\varepsilon(u^\varepsilon(t)) + \int_0^T \mathcal{R}_\varepsilon(u^\varepsilon, \dot{u}^\varepsilon) + \mathcal{R}_\varepsilon^*(u^\varepsilon, -D\mathcal{E}_\varepsilon(u^\varepsilon)) dt \leq \mathcal{E}_\varepsilon(u^\varepsilon(0))$$

Evolutionary  $\Gamma$  convergence based on (EDP)

- Sandier-Serfaty'04 (general approach)
- here: improved version of M-Rossi-Savare'12 (CVPDE)  $\mathcal{R}(u, v) = \Psi(v)$

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Theorem (Mosco convergence implies evolutionary  $\Gamma$ -convergence)

$\mathbf{X}$  reflexive,  $\exists c, C, \lambda_c > 0, p > 1$  such that  $\mathcal{E}_\varepsilon(\cdot) + \lambda_c \|\cdot\|_{\mathbf{X}}^2$  is convex,  
 $\Psi_\varepsilon(v) \geq c\|v\|_{\mathbf{X}}^p - C, \Psi_\varepsilon^*(\xi) \geq c\|\xi\|_{\mathbf{X}^*}^p - C, \mathcal{E}_\varepsilon(u) \geq c\|u\|_{\mathbf{Z}} - C$  with  $\mathbf{Z} \in \mathbf{X}$

$(\mathcal{E}_\varepsilon \xrightarrow{M} \mathcal{E}_0 \ \& \ \Psi_\varepsilon \xrightarrow{\Gamma} \Psi_0 \text{ in } \mathbf{X}) \implies (\mathbf{X}, \mathcal{E}_\varepsilon, \Psi_\varepsilon) \xrightarrow{evol} (\mathbf{X}, \mathcal{E}_0, \Psi_0)$

Compatibility:  $\mathcal{E}_\varepsilon \xrightarrow{M} \mathcal{E}_0$  and  $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_0$  in SAME topology  $\mathbf{X}$

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- M-Rossi-Savare'12 (CVPDE) uses the much stronger statement  $\Psi_\varepsilon \xrightarrow{M} \Psi_0$  (but allows state dependence via  $\mathcal{R}_\varepsilon(u, \dot{u})$ )
- Sina Reichelt 2014 (not contained in my survey article, see Reichelt-Liero SIMA'16 and exercise in tutorial)

$\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_0$  is sufficient

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Application to our two simple problems (coercivity of  $\Psi$  defines  $\mathbf{X}$ )

## ODE model on $\mathbf{X} = \mathbb{R}^2$

We always have  $\mathcal{E}_\varepsilon \xrightarrow{M} \mathcal{E}$  and  $\mathcal{R}_\varepsilon \xrightarrow{M} \mathcal{R}$ .

$$\mathcal{R}(v) = \frac{1}{2}(v_1^2 + v_2^2/\varepsilon^\beta) \text{ and } \mathcal{R}^*(\xi) = \frac{1}{2}(\xi_1^2 + \varepsilon^\beta \xi_2^2)$$

Theorem is applicable for  $\beta = 0$  only,  
 because of needed equicoercivity of for  $\Psi_\varepsilon^*$ .

# 4. Energy-dissipation formulations

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Application to our two simple problems (coercivity of  $\Psi$  defines  $\mathbf{X}$ )

**Homogenization:**  $c\|v\|_{L^2}^2 \leq \Psi_\varepsilon(v) \leq C\|v\|_{L^2}^2 \implies \mathbf{X} = L^2(0, \ell).$

$$\mathcal{E}_\varepsilon(u) = \frac{1}{2} \int_0^\ell a_\varepsilon u_x^2 + b_\varepsilon u^2 dx : \quad \mathcal{E}_\varepsilon \xrightarrow{M} \mathcal{E}_0 \text{ in } \mathbf{X} = L^2(0, \ell) \quad \oplus$$

$$\Psi_\varepsilon(v) = \frac{1}{2} \int_0^\ell c(x/\varepsilon)v(x)^2 dx : \quad \text{not } \xrightarrow{M}, \text{ but } \Psi_\varepsilon \xrightarrow{\Gamma} \Psi_{\text{strong}} \quad \oplus$$

Theorem is applicable and gives  $c_{\text{eff}} = c_{\text{arith}} \geq c_{\text{harm}}$ .

$$c_\varepsilon u_t = (a_\varepsilon u_x)_x - b_\varepsilon u \xrightarrow{\text{evol}} c_{\text{arith}} u_t = (a_{\text{harm}} u_x)_x - b_{\text{arith}} u$$

**Sketch of proof of theorem:**  $u_\varepsilon$  are solutions of (i) = (EDB) $_\varepsilon$ :

$$\mathcal{E}_\varepsilon(u_\varepsilon(T)) + \int_0^T (\Psi_\varepsilon(\dot{u}_\varepsilon) + \Psi_\varepsilon^*(\xi_\varepsilon)) dt = \mathcal{E}_\varepsilon(u_\varepsilon(0)) \text{ where } -\xi_\varepsilon(t) \in D\mathcal{E}_\varepsilon(u_\varepsilon(t))$$

■ Uniform coercivity of  $\mathcal{E}_\varepsilon$ ,  $\Psi_\varepsilon$ . and  $\Psi_\varepsilon^*$  yield uniform a priori bounds

$$\|u_\varepsilon\|_{L^\infty([0,T];\mathbf{Z})} + \|u_\varepsilon\|_{W^{1,p}([0,T];\mathbf{X})} + \|\xi_\varepsilon\|_{L^p([0,T];\mathbf{X}^*)} \leq C$$

■ We find convergent subsequences (still  $\varepsilon = \varepsilon_k$ )

$$u_\varepsilon(t) \rightarrow u(t) \text{ in } \mathbf{X}, \quad u_\varepsilon \rightharpoonup u \text{ in } W^{1,p}([0,T];\mathbf{X}), \quad \xi_\varepsilon \rightharpoonup \xi \text{ in } L^p([0,T];\mathbf{X}^*)$$



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(Sebastian Hensel is specialist for versions of Ioffe's theorem!)

The integrand  $(\varepsilon, \xi) \mapsto \Psi_\varepsilon^*(\xi)$  for the functional  $\mathcal{J}(\varepsilon, \xi) := \int_0^T \Psi_\varepsilon^*(\xi) dt$

- is seq. weakly lower semicontinuous
- and convex in  $\xi \in \mathbf{X}^*$

The convergence  $\varepsilon_\varepsilon := \varepsilon \rightarrow 0$  is strong, while the convergence  $\xi_\varepsilon \rightharpoonup \xi$  is weak.

## 4. Energy-dissipation formulations

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Ioffe's theorem doesn't apply as  $\dot{u}_\varepsilon \rightharpoonup \dot{u}$  in  $W^{1,p}(0, T; \mathbf{X})$  and  $\Psi_\varepsilon \xrightarrow{\Gamma} \Psi_0.$

Reichelt's lemma (see tutorial): It still holds because of  $u_\varepsilon(t) \xrightarrow{\mathbf{X}} u(t).$

- Taking  $\varepsilon \rightarrow 0$  in  $(\text{EDE})_\varepsilon$  and using **well-prepared initial cond. (i.c.)** gives

$$\mathcal{E}_0(u(T)) + \int_0^T (\Psi_0(\dot{u}) + \Psi_0^*(\xi)) dt \leq \lim_{\varepsilon \rightarrow 0} \mathcal{E}_\varepsilon(u_\varepsilon(0)) \stackrel{\text{(i.c.)}}{=} \mathcal{E}_0(u(0))$$

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**Strong-weak closedness of  $D\mathcal{E}_\varepsilon$**  if  $\mathcal{E}_\varepsilon \xrightarrow{M} \mathcal{E}_0$  &  $\mathcal{E}_\varepsilon$   $\lambda_c$ -convex (cf. Attouch'84)

$$u_\varepsilon \rightarrow u \text{ in } \mathbf{X} \text{ and } -D\mathcal{E}_\varepsilon(u_\varepsilon) \ni \xi_\varepsilon \rightharpoonup \xi \text{ in } \mathbf{X}^* \Rightarrow -\xi \in D\mathcal{E}_0(u)$$

- Now the Energy-Dissipation Principle shows

- that  $u$  is a solution and
- that  $\mathcal{E}_\varepsilon(t) \rightarrow \mathcal{E}_0(u(t))$  for all  $t \in [0, T].$  QED

## 4. Energy-dissipation formulations

Main tool is Strong-Weak Closedness of the graph of  $(D\mathcal{E}_\varepsilon)_{\varepsilon \in ]0,1[}$

$$(SWC) \quad u_\varepsilon \rightarrow u \text{ in } \mathbf{X} \text{ and } D\mathcal{E}_\varepsilon(u_\varepsilon) \ni \xi_\varepsilon \rightarrow \xi \text{ in } \mathbf{X}^* \Rightarrow \xi \in D\mathcal{E}(u)$$

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This is a consequence of  $\Gamma$ -convergence and convexity!

Theorem (Convexity and  $\xrightarrow{\Gamma}$  imply (SWC), cf. Attouch 1984)

*If all  $\mathcal{E}_\varepsilon$  are lsc and convex, then  $\mathcal{E}_\varepsilon \xrightarrow{M} \mathcal{E}_0$  implies (SWC).*

Proof: Assume  $u_\varepsilon \rightarrow u$ ,  $\xi_\varepsilon \rightarrow \xi$ , and  $\mathcal{E}_\varepsilon(u_\varepsilon) \rightarrow e_*$

Then convexity gives  $(Cvx)_\varepsilon \quad \mathcal{E}_\varepsilon(w) \geq \mathcal{E}_\varepsilon(u_\varepsilon) + \langle \xi_\varepsilon, w - u_\varepsilon \rangle$

For given  $\hat{u}$  the  $\Gamma_s$ -convergence gives a rec. seq.  $\hat{u}_\varepsilon$  with  $\hat{u}_\varepsilon \rightarrow \hat{u}$ ,  $\mathcal{E}_\varepsilon(\hat{u}_\varepsilon) \rightarrow \mathcal{E}_0(\hat{u})$

Hence, setting  $w = \hat{u}_\varepsilon$  in  $(Cvx)_\varepsilon$  gives  $\underbrace{\mathcal{E}_\varepsilon(\hat{u}_\varepsilon)}_{\rightarrow \mathcal{E}_0(\hat{u})} \geq \underbrace{\mathcal{E}_\varepsilon(u_\varepsilon)}_{\rightarrow e_*} + \underbrace{\langle \xi_\varepsilon, \hat{u}_\varepsilon - u_\varepsilon \rangle}_{\rightarrow \langle \xi, \hat{u} - u \rangle}$

Taking the limit  $\varepsilon \rightarrow 0$  we obtain the relation  $\mathcal{E}_0(\hat{u}) \geq e_* + \langle \xi, \hat{u} - u \rangle$

Choose  $\hat{u} = u$  we see that  $\mathcal{E}_0(u) \geq e_*$  but  $\Gamma_s$ -liminf gives  $e_* \geq \mathcal{E}_0(u)$ . Thus,  $e_* = \mathcal{E}_0(u)$  and we conclude  $\xi \in \partial\mathcal{E}_0(u)$  as desired. □

$$(EDE) \quad \mathcal{E}_\varepsilon(u_\varepsilon(t)) + \int_0^T (\Psi_\varepsilon(\dot{u}_\varepsilon) + \Psi_\varepsilon^*(-D\mathcal{E}_\varepsilon(u_\varepsilon))) dt \leq \mathcal{E}_\varepsilon(u_\varepsilon(0))$$

---

The **Sandier-Serfaty [2004]** approach is more general.

They do assume

neither Strong-Weak Closedness of  $(\partial\mathcal{E}_\varepsilon)_{\varepsilon \in [0,1]}$

nor the Mosco convergence of  $\Psi_\varepsilon \xrightarrow{M} \Psi$

Instead they assume

$$(i) \quad v_\varepsilon \rightarrow v \implies \liminf_{\varepsilon \rightarrow 0} \Psi_\varepsilon(v_\varepsilon) \geq \Psi_0(v) \quad (\text{w-}\Gamma\text{-liminf})$$

$$(ii) \quad u_\varepsilon \rightarrow u \implies \liminf_{\varepsilon \rightarrow 0} \Psi_\varepsilon^*(D\mathcal{E}_\varepsilon(u_\varepsilon)) \geq \Psi_0^*(D\mathcal{E}_0(u)) \quad (\text{dual w-}\Gamma\text{-liminf})$$

Clearly, (SWC) &  $\Psi_\varepsilon \xrightarrow{M} \Psi$  imply (i) and (ii) but not vice-versa.