# Random walks conditioned to stay in Weyl chambers of type C and D

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#### Abstract

We construct the conditional versions of a multidimensional random walk given that it does not leave the Weyl chambers of type C and of type D, respectively, in terms of a Doob h-transform. Furthermore, we prove functional limit theorems for the rescaled random walks. This is an extension of recent work by Eichelsbacher and König who studied the analogous conditioning for the Weyl chamber of type A. Our proof follows recent work by Denisov and Wachtel who used martingale properties and a strong approximation of random walks by Brownian motion. Therefore, we are able to keep minimal moment assumptions. Finally, we present an alternate function that is amenable to an h-transform in the Weyl chamber of type C.

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#### 1 Introduction

In his classical work [9] Dyson established a connection between dynamical versions of random matrices and non-colliding random particle systems. Indeed, the eigenvalue process of a  $k \times k$  Hermitian Brownian motion has the same distribution as the evolution of k independent standard Brownian motions conditioned never to collide (which means that they are in the same order at all times). This process can also be characterised by saying that a k-dimensional Brownian motion is conditioned on never leaving the Weyl chamber of type A,  $W^{A} = \{x = (x_1, \ldots, x_k) \in \mathbb{R}^k \colon x_1 < \ldots < x_k\}$ .

This conditional process, called *Dyson's Brownian motion*, attracted the interest of various researchers. Several discrete versions were considered. Recently, Eichelsbacher and König [10] constructed, in great generality, the analogous random walk version, i.e., the conditional version of a random walk on  $\mathbb{R}^k$  given that it never leaves  $W^A$ . This result and its proof were recently improved by Denisov and Wachtel [7]. It is the aim

of this paper to extend their analysis to the two cases of the Weyl chambers of type C and D, see Section 2.

Let us first describe the random walk version for the type-A chamber. To fix notation, let  $S(n) = (S_1(n), \ldots, S_k(n))$  denote the position of a random walk in  $\mathbb{R}^k$  started at  $x \in \mathbb{R}^k$  after n steps with components  $S_j(n) = x_j + \xi_j^{(1)} + \ldots + \xi_j^{(n)}, 1 \le j \le k$ , where  $\{\xi_j^{(i)}: 1 \le j \le k, i \in \mathbb{N}\}$  is a family of independent identically distributed random variables. In particular, S(0) = x. We write  $\mathbb{P}_x$  and  $\mathbb{E}_x$  for the corresponding probability measure and expectation.

Actually one can understand conditioning to never leave  $W^A$  in two ways. If  $\tau^A = \inf\{n \in \mathbb{N}_0 \colon S(n) \notin W^A\}$  denotes the exit time from  $W^A$ , then on the one hand one can mean the conditional distribution of the path given the event  $\{\tau^A > m\}$  asymptotically as m grows to infinity, that is,

$$\widehat{\mathbb{P}}_x(S(n) \in dy) = \lim_{m \to \infty} \mathbb{P}_x(S(n) \in dy \mid \tau^{A} > m), \qquad x, y \in W^{A}.$$

On the other hand, one can make a change of measure by Doob's h-transform [8]. Necessary for this procedure is to find a function h which is strictly positive on  $W^{A}$  and regular for the restriction of the transition kernel to  $W^{A}$ , i. e.,

$$\mathbb{E}_x[h(S(1)); \tau^{A} > 1] = h(x), \quad x \in W^{A}.$$

Then a new probability transition function on  $W^{A}$  is defined by

$$\widehat{\mathbb{P}}_x^{(h)}(S(n) \in dy) = \mathbb{P}_x(S(n) \in dy; \tau^{\mathcal{A}} > n) \frac{h(y)}{h(x)}, \qquad x, y \in W^{\mathcal{A}}.$$

The corresponding Markov chain is called the h-transform on  $W^A$ . A priori there may be more than one function h amenable to this procedure. However, if a positive regular function h governs the upper tails of  $\tau^A$ , i.e.,  $\mathbb{P}_x(\tau^A > n) \sim C_1h(x)n^{-c_2}$  as  $n \to \infty$  for some  $C_1, c_2 > 0$  for any  $x \in W^A$ , then the two above constructions lead to the same process. Indeed, by the Markov property one obtains in the limit  $m \to \infty$ 

$$\mathbb{P}_{x}(S(n) \in \mathrm{d}y \mid \tau^{\mathrm{A}} > m) = \mathbb{P}_{x}(S(n) \in \mathrm{d}y; \tau^{\mathrm{A}} > n) \frac{\mathbb{P}_{y}(\tau^{\mathrm{A}} > m - n)}{\mathbb{P}_{x}(\tau^{\mathrm{A}} > m)}$$
$$\to \mathbb{P}_{x}(S(n) \in \mathrm{d}y; \tau^{\mathrm{A}} > n) \frac{h(y)}{h(x)}.$$

Eichelsbacher and König succeeded in finding a positive regular function  $V^{A}$  which yields this coincidence:

$$V^{A}(x) = h^{A}(x) - \mathbb{E}_{x}[h^{A}(S(\tau^{A}))], \quad x \in W^{A},$$
 (1)

with  $h^{A}$  the Vandermonde determinant

$$h^{\mathcal{A}}(x) = \prod_{1 \le i < j \le k} (x_j - x_i) = \det\left((x_i^{j-1})_{i,j \in \{1,\dots,k\}}\right).$$

It should be noted that  $h^{A}$  is, up to a multiplicative positive constant, the unique harmonic function that is positive on the interior of  $W^{A}$  and vanishes on the boundary.

In potential theoretic terms, this is expressed by saying that  $h^{A}$  is the *réduite* of  $W^{A}$ . Furthermore, Eichelsbacher and König showed that the rescaled random walk weakly converges to Dyson's Brownian motion.

It is not easy to see that  $V^A$  is strictly positive on  $W^A$ , and it is surprisingly difficult to prove that  $V^A$  is well-defined, i.e., that  $h^A(S(\tau^A))$  is integrable. The approach in [10] is based on the discrete analogue of the Karlin-McGregor formula [13] for random walks and an application of a local central limit theorem. By repeated use of the Hölder inequality, Eichelsbacher and König lose track of minimal moment assumptions: they need the finiteness of moments  $\mathbb{E}[|\xi_j^{(i)}|^r]$  with  $r > ck^3$ . Denisov and Wachtel [7] improve their results by showing that the minimal moment requirement is actually r = k - 1 for k > 3. Since the k-dimensional Vandermonde determinant is a polynomial which has in each variable at most order k - 1, at least moments of order k - 1 are necessary. Hence Denisov and Wachtel's moment requirement is indeed minimal. For k = 3 they need higher moments since their approach is based on a strong coupling which will be explained later.

A conditional version on never leaving  $W^{A}$  under less integrability has not been constructed yet, and it is unclear how to do that. This is a bit surprising since the question of leaving  $W^{A}$  or not has a priori nothing to do with moments.

The case k=2 has already been extensively treated in the literature if one notes that staying in order for two walkers can be translated to staying positive for a single walker. See for example Bertoin and Doney [4].

The paper is organised as follows. We present our main results in Section 2, and the proofs are given in Section 3. In the final Section 4, we discuss an alternate h-transform for the restriction of the walk to the Weyl chamber of type C.

## 2 Extension to Weyl chambers C and D

A natural extension of the conditioned random walk setting just described is the one to different Weyl chambers. They arise in Lie theory as the set of orbits of the adjoint action on a Lie algebra or conjugation under the associated Lie group and first appeared in connection with Brownian motion in a work by Grabiner [11]. There he considered Weyl chambers of type A, B, C, D, but one should mention that the Weyl chambers B and C are actually identical. Those of type C and D are defined as follows:

$$W^{C} = \{x = (x_1, \dots, x_k) \in \mathbb{R}^k : 0 < x_1 < \dots < x_k\},$$

$$W^{D} = \{x = (x_1, \dots, x_k) \in \mathbb{R}^k : |x_1| < x_2 < \dots < x_k\}.$$

As already mentioned, the Weyl chamber of type A imposes a strict order. For type C, an additional wall at zero occurs, and for D the mirror image of  $x_1$  is incorporated into the order. It is important (see [11, Section 6] or [12, Lemma 3]) that these chambers are also equipped with respective réduites:

$$h^{D}(x) = \prod_{1 \le i < j \le k} (x_j^2 - x_i^2)$$
 and  $h^{C}(x) = h^{D}(x) \prod_{i=1}^{k} x_i$ .

As we can handle the two cases simultaneously, we will write Z for C or D. Of course we need corresponding exit times

$$\tau^{\mathbf{Z}} = \inf\{n \in \mathbb{N}_0 \colon S(n) \notin W^{\mathbf{Z}}\}.$$

One of the most important objects of this paper is the function

$$V^{Z}(x) = h^{Z}(x) - \mathbb{E}_{x}[h^{Z}(S(\tau^{Z}))], \qquad x \in W^{Z};$$
 (2)

its role will turn out to be analogous to the role of  $V^{A}$  for  $W^{A}$ .

Let us formulate our assumptions on the random walk  $(S(n))_{n\in\mathbb{N}_0}$ , which are supposed for the results of this section:

Moment Assumption (MA):  $\mathbb{E}[|\xi_j^{(i)}|^{(r^Z)}] < \infty$ , where  $r^C = 2k - 1$  and  $r^D = 2k - 2$  if  $k \geq 3$ , and  $r^C = 3$  and  $r^D > 2$  arbitrary in case k = 2.

Symmetry Assumption (SA):  $\mathbb{E}[(\xi_i^{(j)})^r] = 0$  for any odd integer  $r \leq r^Z$ .

Normalization Assumption (NA):  $\mathbb{E}[(\xi_i^{(j)})^2] = 1$ .

Like in case A, one sees from the definition of  $h^{\mathbb{Z}}$  that the moment requirements are indeed minimal for the integrability of  $h^{\mathbb{Z}}(S(n))$  in the cases  $k \geq 3$ . In the case k = 2 we need for  $\mathbb{Z} = \mathbb{D}$  some higher power for yet another application of the strong coupling, since second moments do not suffice. The assumption (SA) is crucial for the proof of the martingale property stated in Proposition 2.1. The normalization assumption is just for convenience.

First we make an interesting observation about a martingale property:

**Proposition 2.1.** The function  $h^Z$  is regular for  $(S(n))_{n \in \mathbb{N}_0}$ , i.e., for any  $x \in \mathbb{R}^k$  we have  $\mathbb{E}_x[h^Z(S(1))] = h^Z(x)$ . Thus,  $(h^Z(S(n)))_{n \in \mathbb{N}_0}$  is a martingale for any  $x \in \mathbb{R}^k$ .

The proof uses the exchangeability of the step distribution of the random walk  $(S(n))_{n\in\mathbb{N}_0}$  only, not the independence of the components. The case Z=A was treated in [14]. Two important properties of  $V^Z$  are that this function is well-defined and strictly positive on  $W^Z$ . We combine these properties with some results that are of interest in themselves:

**Proposition 2.2.** a)  $V^{\mathbb{Z}}(x) = \lim_{n \to \infty} \mathbb{E}_x[h^{\mathbb{Z}}(S(n)); \tau^{\mathbb{Z}} > n]$  for all  $x \in W^{\mathbb{Z}}$ ;

- b)  $V^{\mathbb{Z}}$  ist monotone in the sense that  $V^{\mathbb{Z}}(x) \leq V^{\mathbb{Z}}(y)$  if  $x_j x_{j-1} \leq y_j y_{j-1}$  for  $2 \leq j \leq k$  and additionally either  $x_1 \leq y_1$  (Z=C) or  $x_1 + x_2 \leq y_1 + y_2$  (Z=D);
- c)  $V^{\mathbb{Z}}(x) \sim h^{\mathbb{Z}}(x)$  in the limit  $\inf_{2 \leq j \leq k} (x_j x_{j-1}) \to \infty$  together with  $x_1 \to \infty$  (Z = C) or  $(x_1 + x_2) \to \infty$  (Z = D) respectively;
- d) there is c positive such that  $V^{Z}(x) \leq ch_{2}^{Z}(x)$  for all  $x \in W^{Z}$ , with  $h_{t}^{D}(x) = \prod_{1 \leq i < j \leq k} (t + |x_{j} x_{i}|)(t + |x_{j} + x_{i}|)$  and  $h_{t}^{C}(x) = h_{t}^{D}(x) \prod_{i=1}^{k} (t + |x_{i}|);$
- e)  $V^{\mathbf{Z}}(x) > 0$  for all  $x \in W^{\mathbf{Z}}$ .

With help of these insights we get a hold on the upper tails of the exit time:

**Theorem 2.3.** The asymptotic behavior for  $n \to \infty$  of the exit time starting from  $x \in W^Z$  is given by

$$\mathbb{P}_x(\tau^{\mathbf{Z}} > n) \sim \varkappa^{\mathbf{Z}} V^{\mathbf{Z}}(x) n^{-(\alpha^{\mathbf{Z}})/2}$$

with  $\alpha^{C} = k^2$  and  $\alpha^{D} = k^2 - k$ , and  $\varkappa^{C}$ ,  $\varkappa^{D}$  the following constants:

$$\begin{split} \varkappa^{\mathrm{D}} & = & \frac{2^{(3k^2-3k+2)/2}}{\pi^k k!} \prod_{1 \leq i < j \leq k} [(2j-1)^2 - (2i-1)^2]^{-1} \prod_{i=1}^k \left[ \Gamma\left(1+\frac{i}{2}\right) \Gamma\left(\frac{i}{2}\right) \right] \\ \varkappa^{\mathrm{C}} & = & \varkappa^{\mathrm{D}} 2^{(3k-2)/2} \prod_{i=1}^k (2k+1-2i)^{-1} \frac{\Gamma\left(\frac{k+1}{2}\right)}{\Gamma\left(\frac{1}{2}\right)}. \end{split}$$

As in [7, Example 2], we can see that our moment assumption (MA) is optimal for the validity of Theorem 2.3. Indeed, if  $\mathbb{P}(\xi_1^{(1)} > u) \sim u^{-\alpha}$  for some  $\alpha \in (r^Z - 1, r^Z)$ , their argument shows that  $\mathbb{P}_x(\tau^Z > n) \gg n^{-(\alpha^Z)/2}$ .

The next result shows that  $V^{\mathbb{Z}}$  is indeed suitable for an h-transform:

**Proposition 2.4.**  $V^{Z}$  is regular for the restriction of the transition kernel to  $W^{Z}$ .

In particular, using Theorem 2.3, the two ways of conditioning the walk to stay in  $W^{\mathbb{Z}}$  that we mentioned in the introduction coincide.

Furthermore, we prove a functional limit theorem for the conditional walk in the spirit of Donsker's theorem. Let us introduce the limit processes of the scaled random walks and state our result. For a k-dimensional Brownian motion one can make a change of measure in the sense of Doob's k-transform using the corresponding réduite:

$$\widehat{\mathbb{P}}_{x}^{(h^{\mathbf{Z}})}(B(t) \in \mathrm{d}y) = \mathbb{P}_{x}(B(t) \in \mathrm{d}y; \tau^{\mathrm{BM}, \mathbf{Z}} > t) \frac{h^{\mathbf{Z}}(y)}{h^{\mathbf{Z}}(x)}, \qquad x, y \in W^{\mathbf{Z}},$$

with  $\tau^{\mathrm{BM,Z}} = \inf\{t \geq 0 \colon B(t) \notin W^{\mathrm{Z}}\}$  denoting the exit time of the Brownian motion from the type-Z Weyl chamber when started at x. We will term the corresponding processes Dyson's Brownian motion of type Z (note that for Z = D the process that receives that name in [6] is obtained from the above one by applying an elementary bijection). It is possible to start these processes from the origin (this can be seen by the same arguments as in [16]). They can also be seen [12] as the eigenvalue processes of dynamical random matrices in the symmetry classes C and D of Altland and Zirnbauer [1], [2]. For more comments on the physical background and the corresponding stochastic differential equations see [12].

**Theorem 2.5.** For  $x \in W^{\mathbb{Z}}$ , as  $n \to \infty$ ,

$$\mathbb{P}_x\left(\frac{1}{\sqrt{n}}S(n)\in\cdot\mid\tau^{\mathbb{Z}}>n\right)\Rightarrow\mu^{\mathbb{Z}},$$

with  $\mu^Z$  the probability measure on  $W^Z$  with density proportional to  $h^Z(y) \exp(-|y|^2/2)$ . Additionally the process  $(X^n(t))_{t\geq 0} = (\frac{1}{\sqrt{n}}S([nt]))_{t\geq 0}$  under the probability measure  $\widehat{\mathbb{P}}_{x\sqrt{n}}^{(V^Z)}$  weakly converges to Dyson's Brownian motion of type Z started at x. Under  $\widehat{\mathbb{P}}_x^{(V^Z)}$ , this process converges weakly to Dyson's Brownian motion of type Z started at zero.

### 3 Proofs

First we prove the regularity of  $h^{\mathbb{Z}}$  on  $\mathbb{R}^{k}$ , which is essential for our purposes:

**Proof of Proposition 2.1.** We make an induction on the number k of components. For this we exploit the Vandermonde determinant representation and write  $h^{\mathbb{Z}}$  in the form

$$h^{\mathbf{Z}}(x) = \det \left[ (x_j^{2i-2+\gamma})_{i,j \in [k]} \right], \qquad [k] = \{1, \dots, k\},$$

where  $\gamma = 1$  for Z = C and  $\gamma = 0$  for Z = D. We dispense with another superscript as not to overburden the notation. For k = 1 the assertion trivially holds either by (SA) (Z = C) or a constant determinant (Z = D). Now fix  $k \geq 2$  and assume that our assertion is true for k - 1. For any  $x \in \mathbb{R}^k$  and  $m \in [k]$  we define

$$h_m^{\mathbf{Z}}(x) = \det \left[ (x_j^{2i-2+\gamma})_{i \in [k-1], j \in [k] \setminus \{m\}} \right],$$

which is the determinant of the matrix that we obtain by deleting the last row and the mth column. In particular, it is a (k-1)-dimensional analogue of  $h^{\mathbb{Z}}$ . Using Laplace expansion we write

$$h^{\mathbf{Z}}(x) = \sum_{m=1}^{k} (-1)^{m-1} x_m^{2k-2+\gamma} h_m^{\mathbf{Z}}(x).$$

We use this in the expectation and denote by  $\mu$  the step distribution of the random walk, to obtain

$$\mathbb{E}_x[h^{\mathbf{Z}}(S(1))] = \int_{\mathbb{R}^k} \mu(\mathrm{d}y) h^{\mathbf{Z}}(x+y) = \sum_{m=1}^k (-1)^{m-1} \int_{\mathbb{R}^k} \mu(\mathrm{d}y) (x_m + y_m)^{2k-2+\gamma} h_m^{\mathbf{Z}}(x+y).$$

We denote by  $\nu$  the m-th marginal of  $\mu$ , which does not depend on m by exchangeability, and by  $\mu_m(\mathrm{d}\tilde{y}|y_m)$  the conditional distribution of  $\mu$  given the coordinate  $y_m$ , which is exchangeable for  $\tilde{y} = (y_1, \ldots, y_{m-1}, y_{m+1}, \ldots, y_k)$ . Hence,  $\mu(\mathrm{d}y) = \nu(\mathrm{d}y_m)\mu_m(\mathrm{d}\tilde{y}|y_m)$ . By our induction hypothesis we have for any  $y_m \in \mathbb{R}$  and  $x \in \mathbb{R}^k$  that

$$\int_{\mathbb{R}^{k-1}} \mu_m(\mathrm{d}\tilde{y}|y_m) h_m^{\mathrm{Z}}(x+y) = h_m^{\mathrm{Z}}(x).$$

This allows us to complete our computation:

$$\mathbb{E}_{x}[h^{\mathbf{Z}}(S(1))] = \sum_{m=1}^{k} (-1)^{m-1} \int_{\mathbb{R}} \nu(\mathrm{d}y_{m}) (x_{m} + y_{m})^{2k-2+\gamma} h_{m}^{\mathbf{Z}}(x)$$

$$= \sum_{m=1}^{k} (-1)^{m-1} \int_{\mathbb{R}} \nu(\mathrm{d}z) \sum_{l=0}^{2k-2+\gamma} {2k-2+\gamma \choose l} x_{m}^{2k-2+\gamma-l} z^{l} h_{m}^{\mathbf{Z}}(x)$$

$$= \sum_{l=0}^{2k-2+\gamma} \int_{\mathbb{R}} \nu(\mathrm{d}z) z^{l} {2k-2+\gamma \choose l} \sum_{m=1}^{k} (-1)^{m-1} x_{m}^{2k-2+\gamma-l} h_{m}^{\mathbf{Z}}(x),$$

where we used (SA) in the third line. Now we apply the Laplace expansion to the m-sum in the last line. For  $l \geq 2$  this m-sum vanishes since its summands are equal to the determinants of matrices with two identical columns. For l = 0 it is equal to  $h^{\mathbb{Z}}(x)$ . This finishes the proof.

Now we prove regularity of  $V^{\mathbb{Z}}$  on  $W^{\mathbb{Z}}$ .

**Proof of Proposition 2.4.** For any  $x \in W^{\mathbb{Z}}$  we get by the strong Markov property and the martingale property of Proposition 2.1

$$\begin{split} &\mathbb{E}_{x}[V^{\mathbf{Z}}(S(1))1_{\{\tau^{\mathbf{Z}}>1\}}] = \\ &= \mathbb{E}_{x}[h^{\mathbf{Z}}(S(1))1_{\{\tau^{\mathbf{Z}}>1\}}] - \mathbb{E}_{x}[\mathbb{E}_{S(1)}[h^{\mathbf{Z}}(S(\tau^{\mathbf{Z}}))]1_{\{\tau^{\mathbf{Z}}>1\}}] \\ &= \mathbb{E}_{x}[h^{\mathbf{Z}}(S(1))1_{\{\tau^{\mathbf{Z}}>1\}}] - \mathbb{E}_{x}[h^{\mathbf{Z}}(S(\tau^{\mathbf{Z}}))1_{\{\tau^{\mathbf{Z}}>1\}}] \\ &= \mathbb{E}_{x}[h^{\mathbf{Z}}(S(1))1_{\{\tau^{\mathbf{Z}}>1\}}] - \mathbb{E}_{x}[h^{\mathbf{Z}}(S(\tau^{\mathbf{Z}}))] + \mathbb{E}_{x}[h^{\mathbf{Z}}(S(\tau^{\mathbf{Z}}))1_{\{\tau^{\mathbf{Z}}\leq 1\}}] \\ &= \mathbb{E}_{x}[h^{\mathbf{Z}}(S(1))] - \mathbb{E}_{x}[h^{\mathbf{Z}}(S(1))1_{\{\tau^{\mathbf{Z}}\leq 1\}}] - \mathbb{E}_{x}[h^{\mathbf{Z}}(S(\tau^{\mathbf{Z}}))] + \mathbb{E}_{x}[h^{\mathbf{Z}}(S(\tau^{\mathbf{Z}}))1_{\{\tau^{\mathbf{Z}}\leq 1\}}] \\ &= V^{\mathbf{Z}}(x). \end{split}$$

Now we turn to the proofs of the remaining results, Proposition 2.2 and Theorems 2.3 and 2.5. We decided to give a sketch only, since the methods employed by Denisov and Wachtel for the case  $W^{\rm A}$  can be straightforwardly extended. Rather than recapping all technical details in tedious length, we indicate the key steps of their proof and illuminate the differences that are necessary to adapt.

We first explain how they obtain their analogues to Proposition 2.2 and Theorems 2.3 and 2.5 for the Weyl chamber of type A, in particular the asymptotics  $\mathbb{P}_x(\tau^Z > n) \sim \varkappa^A V^A(x) n^{-k(k-1)/4}$ , with  $\varkappa^A$  a constant, and the weak convergence to Dyson's Brownian motion of type A. Their idea is to consider, additionally to  $\tau^A$ , the stopping time

$$T^{\mathcal{A}} = \inf\{n \in \mathbb{N}_0 \colon h^{\mathcal{A}}(S(n)) \le 0\}.$$

This has the advantage that the triggering of the defining condition can more easily be exploited for estimates since we have control over the sign of  $h^{A}(S(n))$ . Furthermore, obviously,  $T^{A} \geq \tau^{A}$  almost surely. Therefore, certain estimates involving  $T^{A}$  can be directly transferred to estimates involving  $\tau^{A}$ . Crucial for their approach is the fact that  $(h^{A}(S(n)))_{n \in \mathbb{N}_{0}}$  is a martingale. This yields that the sequence  $Y_{n} = h^{A}(S(n))1_{\{T^{A}>n\}}$  is a nonnegative submartingale (this is solely based on the martingale property). With this they next show that there is a universal constant C such that

$$\mathbb{E}_x[h^{\mathcal{A}}(S(n)); T^{\mathcal{A}} > n] \le Ch_2^{\mathcal{A}}(x), \qquad n \in \mathbb{N}, x \in W^{\mathcal{A}}, \tag{3}$$

where  $h_t^{A}(x) = \prod_{1 \leq i < j \leq k} (t + |x_j - x_i|)$  is defined similar to the expressions in Proposition 2.2d). Proving (3) is technical and lengthy and uses an auxiliary Weyl chamber defined by

$$W_{n,\varepsilon}^{\mathcal{A}} = \{ x \in \mathbb{R}^k : |x_j - x_i| > n^{1/2 - \varepsilon}, 1 \le i < j \le k \}, \qquad \varepsilon > 0.$$

If a point of  $W^A$  is additionally in  $W_{n,\varepsilon}^A$ , it is far away from the boundary of  $W^A$ . Furthermore,  $W_{n,\varepsilon}^A$  has the property that it is reached by the motion soon with high probability: the probability of the entrance time  $\nu_n^A = \inf\{m \in \mathbb{N}_0 \colon S(m) \in W_{n,\varepsilon}^A\}$  being bigger than  $n^{1-\varepsilon}$  decays exponentially. Indeed, we have  $\mathbb{P}_x(\nu_n^A > n^{1-\varepsilon}) \leq \exp\{-Cn^{\varepsilon}\}$ .

This can be shown by a subdivision of the trajectory into  $n^{\varepsilon}$  pieces and an application of the central limit theorem to the pieces. Also expectations of  $h^{A}(S(n))$  on the event  $\{\nu_{n}^{A} > n^{1-\varepsilon}\}$  decay exponentially, hence one can extend estimates for expectations which start from  $x \in W_{n,\varepsilon}^{A}$  to expectations with arbitrary starting points in  $W^{A}$  by the strong Markov property. For the former one can elementarily derive upper bounds with standard estimates like Doob's inequality. With this one is able to prove the bound in (3).

Now this in turn yields the integrability of  $h^{A}(S(\tau^{A}))$  by a direct application of martingale arguments. Furthermore, Denisov and Wachtel obtain from this that the function  $V^{(T^{A})}(x) = \lim_{n\to\infty} \mathbb{E}_x[Y_n]$  is well defined on the set  $\{x \colon h^{A}(x) > 0\}$ . For showing that  $V^{A}$  is strictly positive on  $W^{A}$  they use that  $(V^{(T^{A})}(S(n))1_{\{\tau^{A}>n\}})_{n\in\mathbb{N}_0}$  is a supermartingale; again this is solely based on the martingale property.

Here we terminate our survey on the proofs of the corresponding statements of d) and e) of Proposition 2.2. The corresponding results to a) and b) pop out easily from the method of proof. The proof of c) is actually part of the derivation of e). This finishes the sketch of their analogue of Proposition 2.2.

Now we turn to the sketch of the proofs of their analogues of Theorems 2.3 and 2.5. For this a coupling of random walks and Brownian motion by Major [15] is applied which has already been used in other contexts, see [3], [5].

**Lemma 3.1.** Given that  $\mathbb{E}[|\xi_j^{(i)}|^{2+\delta}] < \infty$  for some  $\delta \in (0,1)$ , a Brownian motion  $(B(t))_{t\geq 0}$  can be defined on the same probability space as the random walk  $(S(n))_{n\in\mathbb{N}_0}$  such that, for  $a\in(0,\frac{\delta}{2(2+\delta)})$ ,

$$\mathbb{P}\left(\sup_{u \le n} |S([u]) - B(u)| \ge n^{1/2 - a}\right) = o\left(n^{2a + a\delta - \delta/2}\right).$$

Other important tools for the proof of the asymptotic behavior of  $\mathbb{P}_x(\tau^A > n)$  are estimates for the upper tails of the exit time of Brownian motion from  $W^A$  due to Grabiner [11] and Varopoulos [17]. Again, the auxiliary Weyl chamber  $W_{n,\varepsilon}^A$  is used.

To attack the upper tails of  $\tau^A$ , we know from the above mentioned exponential decay of  $\mathbb{P}_x(\nu_n^A>n^{1-\varepsilon})$  that the random walk reaches  $W_{n,\varepsilon}^A$  after a short time, with high probability. Using the strong Markov property at time  $\nu_n^A$ , we only have to consider starting points  $y\in W_{n,\varepsilon}^A$  instead of x. For those, we use Lemma 3.1 with  $a=2\varepsilon$  and see that the exit times from  $W^A$  for the Brownian motion and the coupled random walk are roughly identical with high probability, since the distances between them,  $n^{1/2-2\varepsilon}$ , are negligible with respect to the typical distances  $n^{1/2-\varepsilon}$  required in  $W_{n,\varepsilon}^A$ . Hence, the upper tails of the random walk exit times can directly be related to the ones of the Brownian motion, which are well-known. After identifying the asymptotic behavior of  $\mathbb{P}_x(\tau^A>n)$ , one can use it to prove the functional limit theorem in a straightforward manner.

So, unlike in the proof of Eichelsbacher and König, there is no need to employ the discrete analogue of the Karlin-McGregor formula, or Hölder's inequality; the results are derived using a comparison to Brownian motion.

Now we argue that these proofs can be straightforwardly extended to cases C and D. This is due to several factors. First, according to Proposition 2.1,  $(h^{\mathbb{Z}}(S(n)))_{n\in\mathbb{N}_0}$  is also a martingale, and one can analogously define the corresponding sub- and supermartingales,  $(h^{\mathbb{Z}}(S(n))1_{\{T_x^{\mathbb{Z}}>n\}})_{n\in\mathbb{N}_0}$  and  $(V^{(T^{\mathbb{Z}})}(S(n))1_{\{\tau^{\mathbb{Z}}>n\}})_{n\in\mathbb{N}_0}$ . Second, as one easily sees, the inequalities  $T_x^{\mathbb{Z}} \geq \tau^{\mathbb{Z}}$  hold almost surely, too. Third, for proving the estimate in d) of Proposition 2.2, we split the functions  $h^{\mathbb{Z}}$  into  $h^{\mathbb{D}}(x) = \prod_{1 \leq i < j \leq k} (x_j - x_i)(x_j + x_i)$  (and  $h^{\mathbb{C}}$  similarly). This is a more suitable representation when used together with the corresponding auxiliary Weyl chambers defined by

$$\begin{array}{lcl} W_{n,\varepsilon}^{\mathrm{D}} &=& \{x \in \mathbb{R}^k \colon |x_j - x_i| > n^{1/2 - \varepsilon}, |x_j + x_i| > n^{1/2 - \varepsilon}, 1 \leq i < j \leq k\}, \\ W_{n,\varepsilon}^{\mathrm{C}} &=& W_{n,\varepsilon}^{\mathrm{D}} \cap \{x \in \mathbb{R}^k \colon |x_i| > n^{1/2 - \varepsilon}, 1 \leq i \leq k\}. \end{array}$$

Again the probability of the entrance time  $\nu_n^{\rm Z} = \inf\{m \in \mathbb{N}_0 \colon S(m) \in W_{n,\varepsilon}^{\rm Z}\}$  being bigger than  $n^{1-\varepsilon}$  decays exponentially, which can be proved by the same argument as for  $W_{n,\varepsilon}^{\rm A}$ .

Fourth, by the works of Grabiner and Varopoulos [11, 17], we have analogous estimates for the upper tails of the exit times from  $W^{\mathbb{Z}}$  at our disposal. Varopoulos formulated them more generally for conical regions (i.e., closed under scaling by a positive constant and addition of elements), and Grabiner formulated them directly for the Weyl chambers:

**Lemma 3.2.** a) For all  $y \in W^Z$  we have with  $\tau^{BM,Z} = \inf\{t \geq 0 : B(t) \notin W^Z\}$ ,

$$\mathbb{P}_{y}(\tau^{\mathrm{BM,Z}} > t) \le C \frac{h^{\mathrm{Z}}(y)}{t^{(\alpha^{\mathrm{Z}})/2}}, \qquad t > 0,$$

where  $\alpha^{C} = k^2$  and  $\alpha^{D} = k^2 - k$ .

b) As  $t \to \infty$ ,

$$\mathbb{P}_y(\tau^{\mathrm{BM,Z}} > t) \sim \varkappa^{\mathrm{Z}} \frac{h^{\mathrm{Z}}(y)}{t^{(\alpha^{\mathrm{Z}})/2}},$$

uniformly in  $y \in W^Z$  satisfying  $|y| \le \theta_t \sqrt{t}$  with some  $\theta_t \to 0$ .

c) For  $y \in W^{\mathbb{Z}}$ , denote by  $b_t^{\mathbb{Z}}(y,z)$  the density of  $\mathbb{P}_y(\tau^{\mathrm{BM},\mathbb{Z}} > t, y + B(t) \in \mathrm{d}z)$ . Then, as  $t \to \infty$ ,

$$b_t^{\mathbf{Z}}(y,z) \sim K^{\mathbf{Z}} t^{-\frac{k}{2}} e^{-|z|^2/(2t)} h^{\mathbf{Z}}(y) h^{\mathbf{Z}}(z) t^{-\alpha^{\mathbf{Z}}}$$

uniformly in  $y, z \in W^Z$  satisfying  $|y| \le \theta_t \sqrt{t}$  and  $|z| \le \sqrt{t/\theta_t}$  with some  $\theta_t \to 0$ , and

$$K^{\mathcal{C}} = \frac{2^k k! \varkappa^{\mathcal{C}}}{\int_{\mathbb{R}^k} e^{-|x|^2/2} |h^{\mathcal{C}}(x)| \, \mathrm{d}x}, \qquad K^{\mathcal{D}} = \frac{2^{k-1} k! \varkappa^{\mathcal{D}}}{\int_{\mathbb{R}^k} e^{-|x|^2/2} |h^{\mathcal{D}}(x)| \, \mathrm{d}x}.$$

Of course we can use the same coupling of random walks and Brownian motion as in Lemma 3.1. Using all these ingredients, we can easily adapt the strategy employed by Denisov and Wachtel to prove Proposition 2.2 and Theorems 2.3 and 2.5.

# 4 An alternate h-transform for $W^{C}$

In this section we present another function,  $\tilde{V}^{\rm C}$ , that is positive and regular on the type-C Weyl chamber,  $W^{\rm C}$ . This function is in general different from the function  $V^{\rm C}$  defined in (2), but because of its positivity and regularity on  $W^{\rm C}$ , it is amenable to an h-transform of the random walk restricted to  $W^{\rm C}$ . This illustrates our remark in the introduction: Not every h-transform of the random walk on a set W is equal to the conditional version of the walk given that it never leaves W. The point is that  $\tilde{V}^{\rm C}$  does not necessarily govern the upper tails of the exit time from  $W^{\rm C}$ , but  $V^{\rm C}$  does; see Theorem 2.3.

The idea of the construction of  $\widetilde{V}^{\mathrm{C}}$  is to first condition every component on staying positive and afterwards conditioning the resulting walk on never violating the order of the components. In other words, we first condition on never leaving  $(0,\infty)^k$  and afterwards on never leaving  $W^{\mathrm{A}}$ . Even though the intersection of these two sets is equal to  $W^{\mathrm{C}}$ , there is no reason to hope that the sequentially conditioned random walk be equal to the conditional walk constructed in Section 2; this is a general fact about conditional probabilities.

Let us now describe the construction of  $\widetilde{V}^{\mathbb{C}}$ . For  $z \in (0, \infty)$ , denote  $V(z) = z - \mathbb{E}_z[S_1(\tau^{+,(1)})]$  with  $\tau^{+,(i)} = \inf\{n \in \mathbb{N}_0 \colon S_i(n) \leq 0\}$  the exit time from  $(0,\infty)$ . When we apply the method of proof of Denisov and Wachtel to k=2 under sufficient moment assumptions, we know that V is positive and regular for the restriction of a one-dimensional symmetric random walk  $(S_1(n))_{n\in\mathbb{N}_0}$  to  $(0,\infty)$  and that it governs the upper tails of the exit time  $\tau^{+,(1)}$ . By independence,  $V^{\otimes k}$  is positive and regular for the restriction of the walk  $(S(n))_{n\in\mathbb{N}_0}$  to  $(0,\infty)^k$  and governs the upper tails of the exit time  $\tau^+ = \inf\{n \in \mathbb{N}_0 \colon S(n) \notin (0,\infty)^k\}$ . As a consequence,

$$\widehat{\mathbb{P}}_{x}^{+}(S(n) \in dy) := \lim_{m \to \infty} \mathbb{P}_{x}(S(n) \in dy \mid \tau^{+} > m)$$

$$= \mathbb{P}_{x}(S(n) \in dy; \tau^{+} > n) \frac{V^{\otimes k}(y)}{V^{\otimes k}(x)}.$$
(4)

Under  $\widehat{\mathbb{P}}_x^+$ , the walk is equal to the conditional version given that it does not leave  $(0,\infty)^k$ . Now we need the version of the function  $V^A$  defined in (1) for  $\widehat{\mathbb{P}}_x^+$ :

$$V^{+,A}(x) = h^{A}(x) - \widehat{\mathbb{E}}_{x}^{+}[h^{A}(S(\tau^{A}))], \quad x \in W^{A}.$$

**Lemma 4.1.** Assume that the step distribution of the walk is symmetric and possesses finite moments of order k-1 for  $k \geq 4$  or of some order r>2 in cases k=3 and k=2. Then the function

$$\widetilde{V}^{C} = V^{+,A} V^{\otimes k}$$

is positive in  $W^{\mathbb{C}}$  and regular for the restriction of the transition kernel to  $W^{\mathbb{C}}$ .

*Proof.* The h-transform of  $\widehat{\mathbb{P}}_x^+$  with  $V^{+,A}$  on  $W^A$  is equal to the conditional version given that the walk does not leave  $W^A$ , i.e.,

$$\lim_{m \to \infty} \widehat{\mathbb{P}}_x^+(S(n) \in \mathrm{d}y \,|\, \tau^{\mathrm{A}} > m) = \widehat{\mathbb{P}}_x^+(S(n) \in \mathrm{d}y; \tau^{\mathrm{A}} > n) \frac{V^{+,\mathrm{A}}(y)}{V^{+,\mathrm{A}}(x)}. \tag{5}$$

Using (4), we see that

$$\widehat{\mathbb{P}}_x^+(S(n) \in \mathrm{d}y; \tau^{\mathrm{A}} > n) = \mathbb{P}_x(S(n) \in \mathrm{d}y; \tau^+ > n, \tau^{\mathrm{A}} > n) \frac{V^{\otimes k}(y)}{V^{\otimes k}(x)}.$$

Using this in (5) and noting that  $\{\tau^+ > n, \tau^A > n\} = \{\tau^C > n\}$ , we arrive at

$$\lim_{m \to \infty} \widehat{\mathbb{P}}_{x}^{+}(S(n) \in \mathrm{d}y \,|\, \tau^{\mathrm{A}} > m) = \mathbb{P}_{x}(S(n) \in \mathrm{d}y; \tau^{\mathrm{C}} > n) \frac{V^{+,\mathrm{A}}(y)}{V^{+,\mathrm{A}}(x)} \frac{V^{\otimes k}(y)}{V^{\otimes k}(x)}$$
$$= \mathbb{P}_{x}(S(n) \in \mathrm{d}y; \tau^{\mathrm{C}} > n) \frac{\widetilde{V}^{\mathrm{C}}(y)}{\widetilde{V}^{\mathrm{C}}(x)}.$$

Since the left hand side is a probability measure in  $y \in W^{\mathbb{C}}$ , the right hand side is as well. This shows the regularity of  $\widetilde{V}^{\mathbb{C}}$  in  $W^{\mathbb{C}}$ . The positivity is obvious.

For the case D this approach does not work since it is not clear how to divide the condition in  $W^{D}$  into two conditions that can separately be handled with the methods presented in this paper.

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## References

- [1] ALTLAND, A. and ZIRNBAUER, M. R., Random matrix theory of a chaotic Andreev qunatom dot, *Phys. Rev. Lett.* **76**, 3420-3423 (1996)
- [2] ALTLAND, A. and ZIRNBAUER, M. R., Nonstandard symmetry classes in mesoscopic normal-superconducting hybrid structure, *Phys. Rev. B* **55**, 1142-1161 (1997)
- [3] Baik, J. and Suidan, T., Random matrix central limit theorems for non-intersecting random walks, Ann. Probab. 35, 1807-1834 (2007)
- [4] Bertoin, J. and Doney, R.A., On conditioning a random walk to stay nonnegative, Ann. Probab. 22, 2152-2167 (1994)
- [5] Bodineau, T. and Martin, J., A universality property for last-passage percolation paths close to the axis, *Electron. Comm. in Probab.* **10**, 105-112 (2005)
- [6] BORODIN, A., FERRARI, P.L., PRÄHOFER, M., SASAMOTO, T. and WARREN, J., Maximum of Dyson Brownian motion and non-colliding systems with a boundary, arXiv:0905.3989v1, 25 May 2009
- [7] Denisov, D. and Wachtel, V., Conditional limit theorems for ordered random walks, arXiv:0907.2854, 16 July 2009
- [8] Doob, J.L., Classical potential theory and its probabilistic counterpart, New York: Springer-Verlag (1984)

- [9] Dyson, F.J., A Brownian-motion model for the eigenvalues of a random matrix, J. Math. Phys. 3, 1191-1198 (1962)
- [10] EICHELSBACHER, P. and KÖNIG, W., Ordered random walks, *Electron. J. Probab.* 13, 1307-1336 (2008)
- [11] Grabiner, D.J., Brownian motion in a Weyl chamber, non-colliding particles, and random matrices, Ann. Inst. H. Poincaré Probab. Statist. 35(2), 177-204 (1999)
- [12] KATORI, M. and TANEMURA, T., Symmetry of matrix-valued stochastic processes and non-colliding diffusion particle systems, J. Math. Phys. 45, 3058-3085 (2004).
- [13] KARLIN, S.P. and McGregor, G., Coincidence probabilities, *Pacif. J. Math.* 9, 1141-1164 (1959).
- [14] KÖNIG, W., O'CONNELL, N. and ROCH, S., Non-colliding random walks, tandem queues, and discrete orthogonal polynomial ensembles, *Elec. Jour. Probab.* 7, Paper no. 5, 1-24 (2002)
- [15] Major, P., The approximation of partial sums of rv's, Z. Wahrscheinlichkeitstheorie verw. Gebiete 35, 213-220 (1976)
- [16] O'CONNELL, N. and YOR, M., A representation for non-colliding random walks, Elec. Comm. Probab. 7, 1-12 (2002)
- [17] VAROPOULOS, N.Th., Potential theory in conical domains, Math. Proc. Camb. Phil. Soc. 125, 335-384 (1999)

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