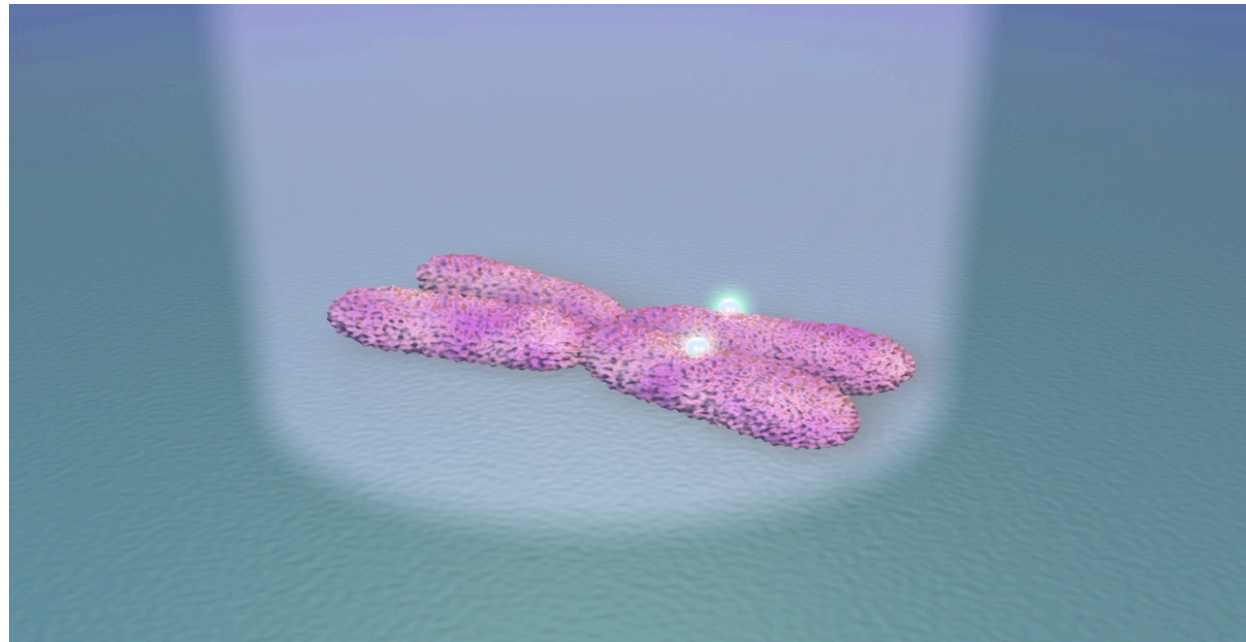


# Wechselwirkung von Licht mit metallischen Nanopartikeln in dielektrischer Umgebung



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# GLIEDERUNG

**1. TEIL:** Die Grundlagen - Plasmonen-Polaritonen Resonanz in Metall NP

**Intro:** Hochlokale und schnelle Temperaturerhöhung - Ablation der NP

**2. TEIL:** Wirkungen hoher Feldstärken in dielektrischer Umgebung von NP

**3. TEIL:** Plasmon Polariton Resonanz unterstützte **Second Harmonic Generation (SHG)**

**4. TEIL:** Terahertz Abstrahlung von metallischen Nanostrukturen

**Abspann:** Coulomb Explosion - Ablation der NP

**Summary:** **ZUSAMMENFASSUNG UND AUSBLICK**

# 1. TEIL: Plasmonen-Polaritonen Resonanz in Metall NP

(Wechsel von Transmission/Reflexion, aber auch Absorption)

## Plasmon

In Metallen stattfindende Dichteoszillationen freier Ladungsträger  $\rho$  in Ihrem eigenen Feld

$$\frac{\partial^2}{\partial t^2} \rho + g \frac{\partial}{\partial t} \rho + f_p \rho = 0 \quad f_p = \omega_p^2$$

gedämpfter Oszillator liefert  $x(t)$

$$\omega_p = 6,8 \cdot 10^{15} \text{ s}^{-1} \quad g = 2,5 \cdot 10^{14} \text{ s}^{-1}$$

$$\epsilon^*(\omega) = 1 - \frac{\omega_p^2}{\omega^2 + ig\omega}$$

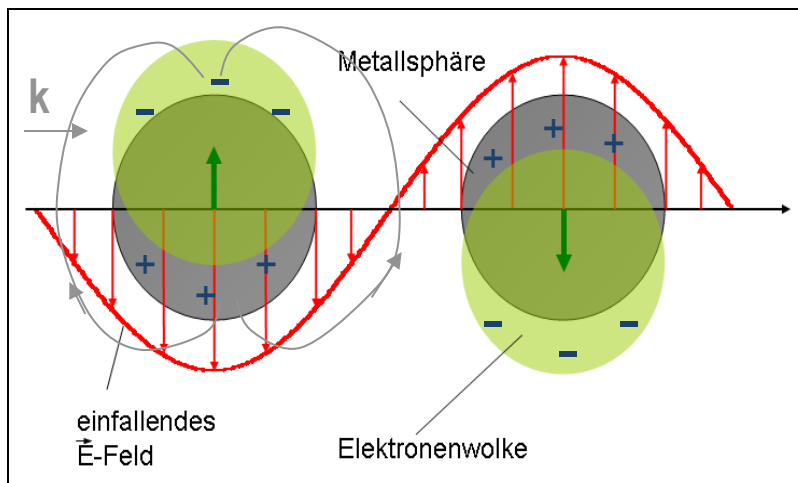
P.B. Johnson, PRB 6, 4370 (1972)

$$\epsilon' = 0.31 \quad \epsilon'' = 4.84 \quad \lambda = 530 \text{ nm}$$

## Partikel Plasmon Polariton

koppelt den Zustand von Photonen mit einer elektronischen Materialanregung

**resonantes Verhalten:** durch die Beschränkung der Elektronen auf das Partikel

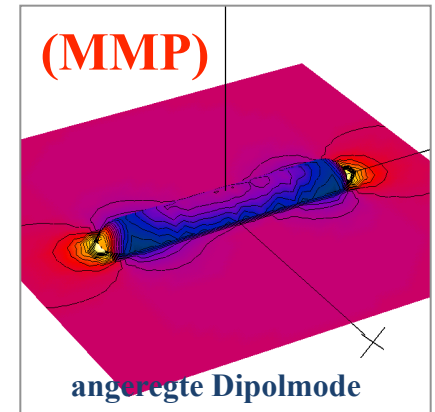
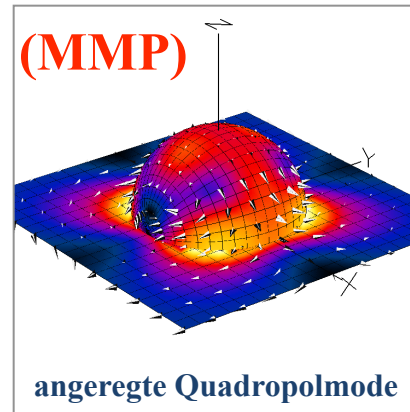
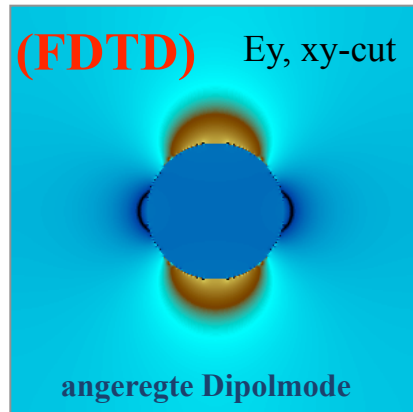


Polarisierbarkeit  $\alpha$  einer Kugel in dielektrischer Umgebung  $\epsilon_a$

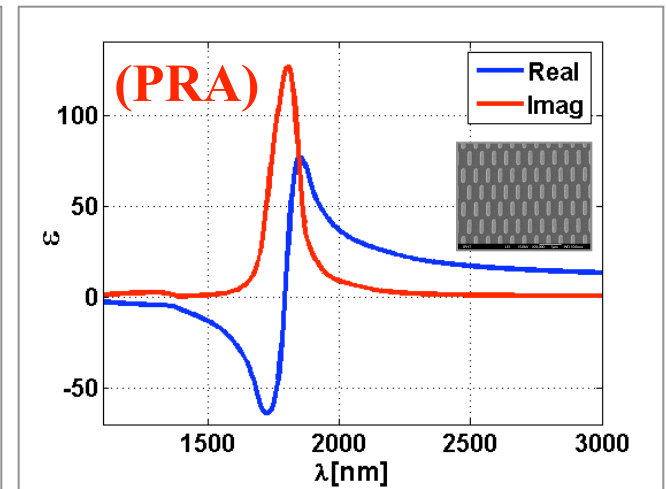
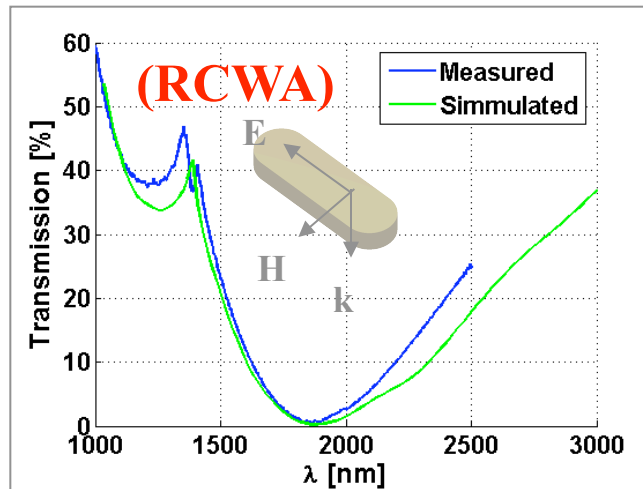
$$\alpha = 3 \frac{\epsilon_p / \epsilon_a - 1}{\epsilon_p / \epsilon_a + 2} = 3 \frac{\epsilon_p - \epsilon_a}{\epsilon_p + 2\epsilon_a}$$

in Resonanz = 0

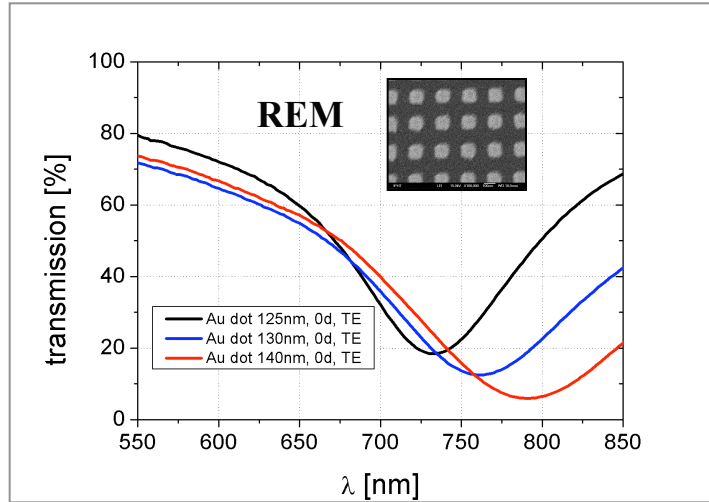
$$2\epsilon_a + \epsilon_p(\omega) = 0 \rightarrow \epsilon_p = 1 - \frac{\omega_p^2}{\omega^2} \rightarrow \omega_R = \frac{\omega_p}{\sqrt{1 - 2\epsilon_a}}$$



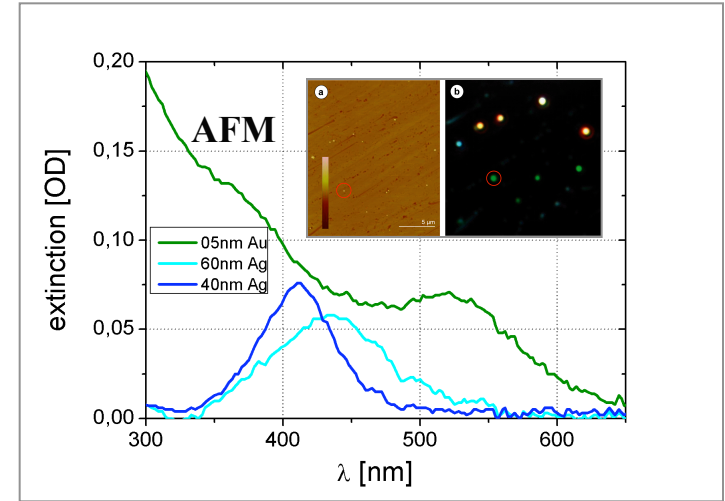
Au-Kugel  $d=30\text{nm}$ ,  $\lambda = 530\text{ nm}$  Au-Kugel  $d=60\text{nm}$ ,  $\lambda = 550\text{ nm}$  Au-Draht,  $\lambda = 1700\text{ nm}$ , E parallel



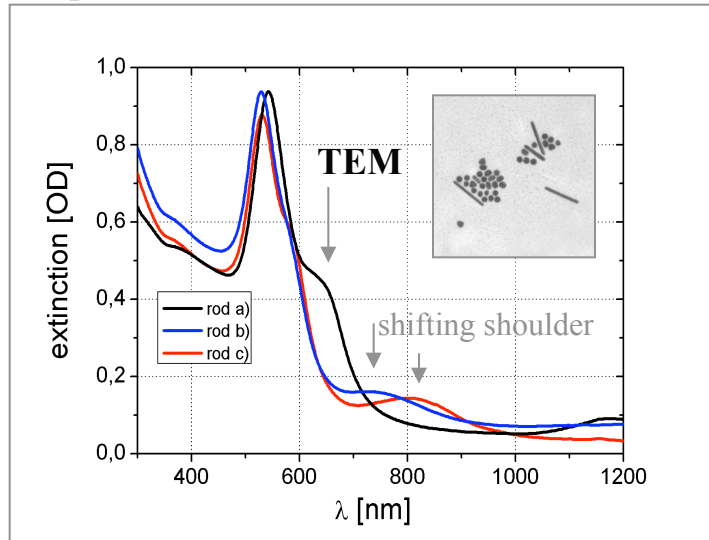
size



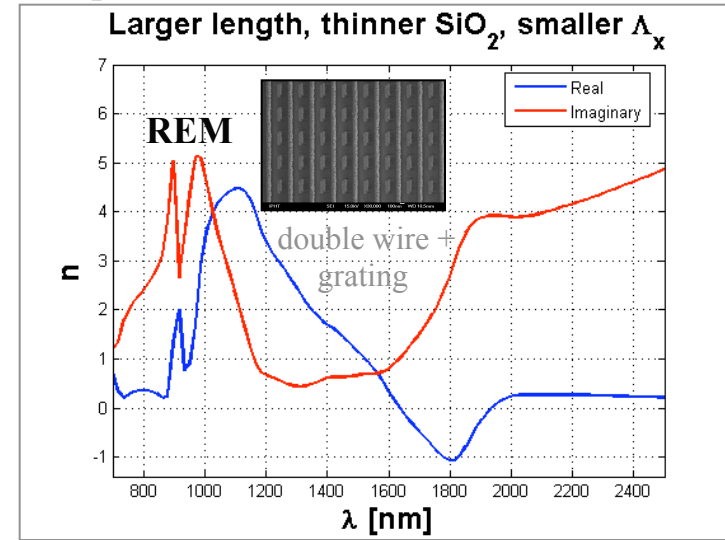
kind of metal

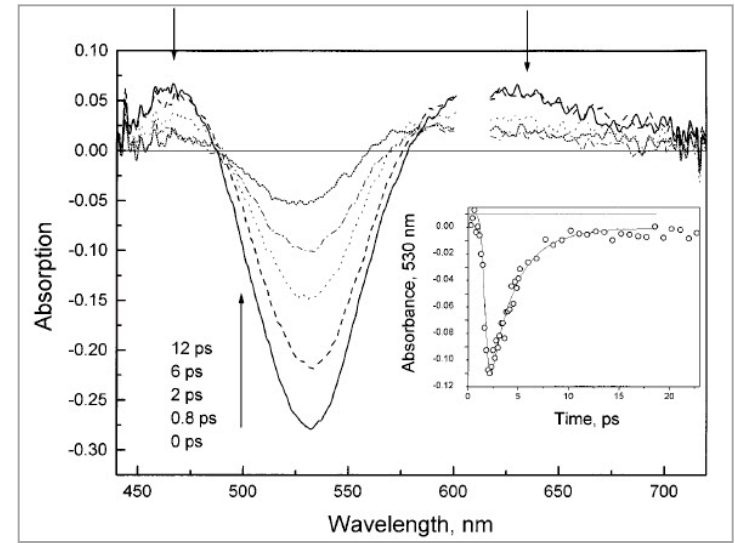
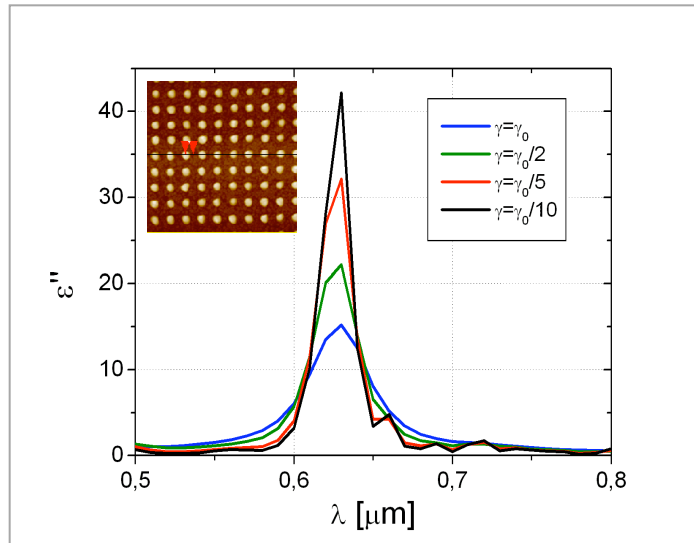


shape



composition





**Mie:** interaction of nanoparticle with light: resonance -> **problem:**

**Kramers-Kronig:** in resonance – high imaginary part of material parameters too

**Drude:** damping mechanism is loss of phase coherence of collective electron oscillation caused by  $e^- - e^-$ ,  $e^-$ -phonon-,  $e^-$ -defect- and  $e^-$ -surface scattering -> rise of electron temperature,  $e^-$ -phonon collision form new thermal equilibrium

-> **phase destruction** of amplified evanescent wave -> unfavourable for image below the Abbe limit

-> **no significant improvement** by **deep temperatures** (4K) and metallic **single crystal** structures

### How we can use a high damping mechanism for another applications?

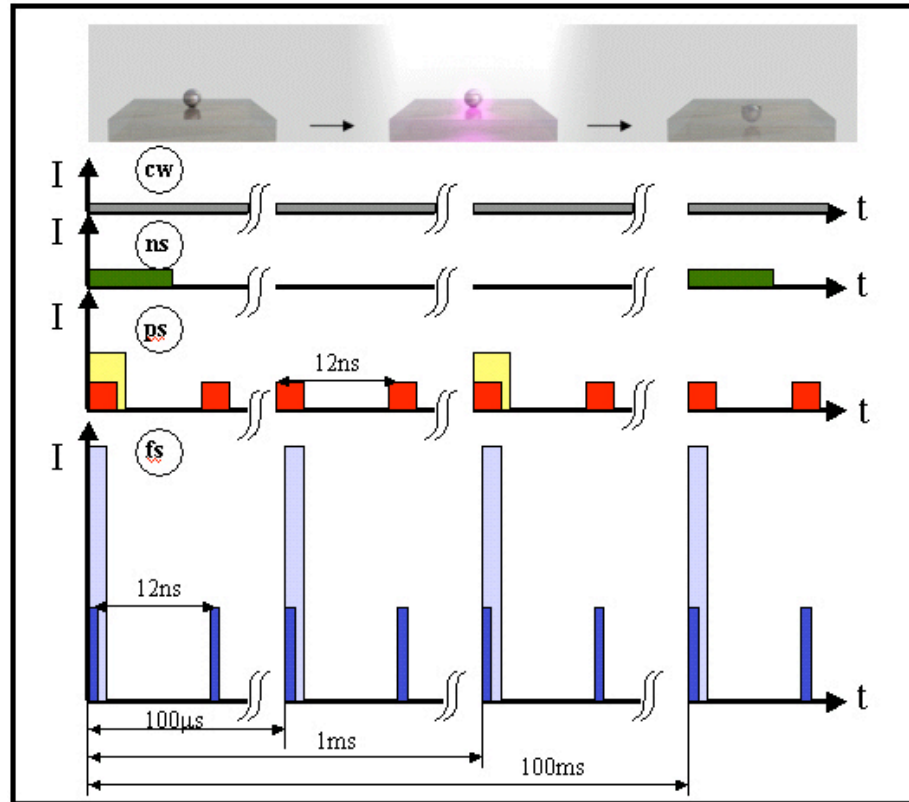
increase damping by nonlinear interaction

-> fs-laser pump-probe experiments on 30 nm gold nanoparticles show **27% transient absorption** [1]

**Intro:** hochlokale und schnelle Temperaturerhöhung – Ablation der NP

(eine Folge der hohen Absorption in Resonanz)





### ns - Laserparameter für Ablationsschwelle

$P = 0.8 \text{ W}$   
 $E_p = 80 \text{ mJ}$   
 $F = 200 \text{ mJ/cm}^2$   
 $P_p = 25 \text{ MW/cm}^2$

### ps – Laserparameter für Ablationsschwelle

$P = 70 \text{ mW}$   
 $E_p = 0.7 \text{ } \mu\text{J}$   
 $F = 100 \text{ mJ/cm}^2$   
 $P_p = 300 \text{ MW/cm}^2$

### fs - Laserparameter für Ablationsschwelle Nanopartikel

$P =$  mittlere Leistung [W]

$E_p =$  Pulsenergie =  $P/f$  [J]

$F =$  Fluenz = Energiedichte/Puls =  $P/A \cdot f$  [ $\text{J/cm}^2$ ]

$P_p =$  Leistungsdichte/Puls =  $F/t_p$  [ $\text{W/cm}^2$ ]

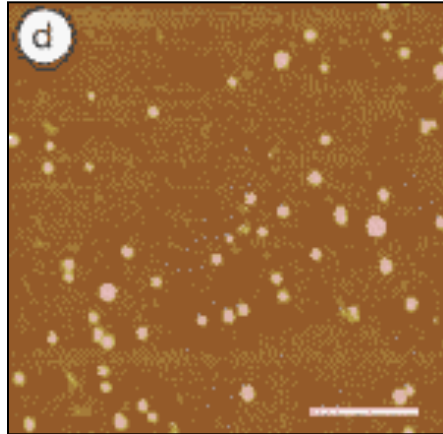
1W (cw Laser)

5 n (Pulswiederholrate)

3  $\text{mJ/cm}^2$  (Fokussierung)

30  $\text{GW/cm}^2$  (Pulslänge)

AFM picture bevor laser radiation:

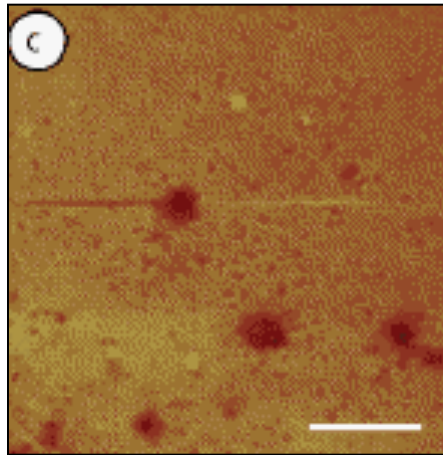


d) 30nm Au NP on top of a 50 nm e<sup>-</sup> and T sensitive PMMA layer

c) hot fragments of NP sink into PMMA

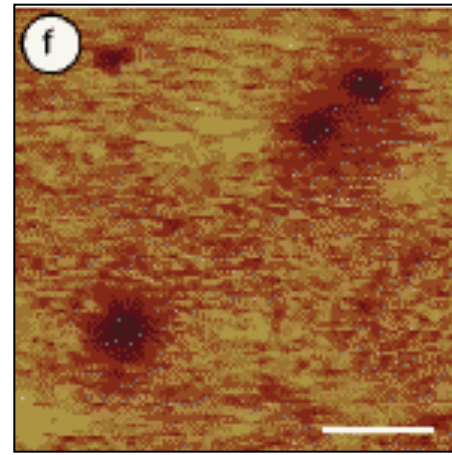
f) hot fragments of NP sink into PMMA,  
fragments smaller and nearby former NP

8 ns, 350mJ/cm<sup>2</sup>



Thermische Ablationsschwelle:  
**200mJ/cm<sup>2</sup>**

35 ps, 240mJ/cm<sup>2</sup>



Thermische Ablationsschwelle:  
**100mJ/cm<sup>2</sup>**

heat equation

$$\rho C_p \frac{\partial T}{\partial t} + \nabla \cdot (-k \nabla T) = q$$

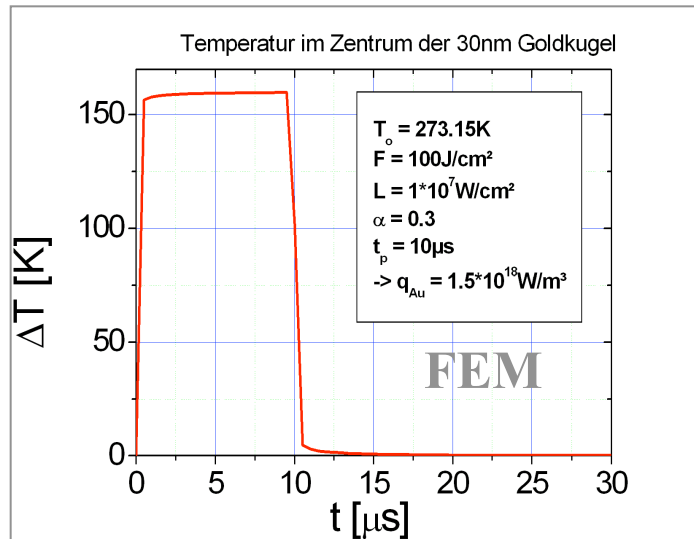
domains: Au NP (sphere with  $r=15\text{nm}$ ),  $\text{H}_2\text{O}$  around the Au NP

boundary between the two subdomains

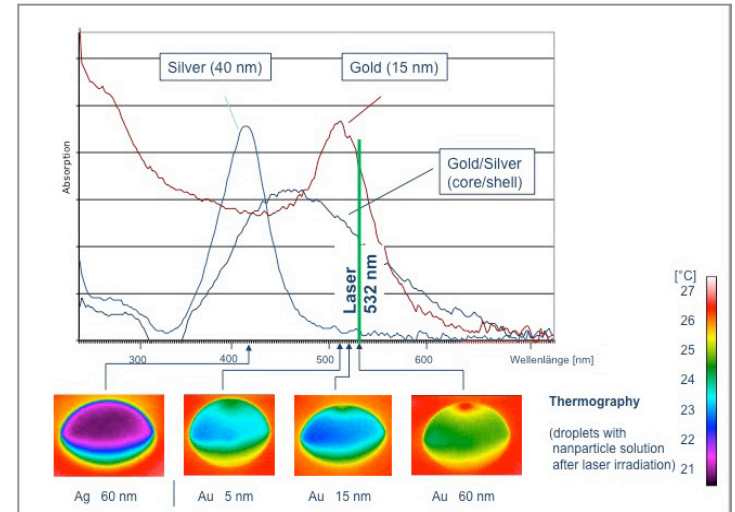
$$-\mathbf{n}_{Au} \cdot (-k_{Au} \nabla T_{Au}) - \mathbf{n}_{H_2O} \cdot (-k_{H_2O} \nabla T_{H_2O}) = 0$$

- conductive heat transfer from Au NP  $\rightarrow$   $\text{H}_2\text{O}$  starts during the laser pulse

$\rho$  - density,  $c_p$  - heat capacity,  $k$  - thermal conductivity,  $q$  - heat source

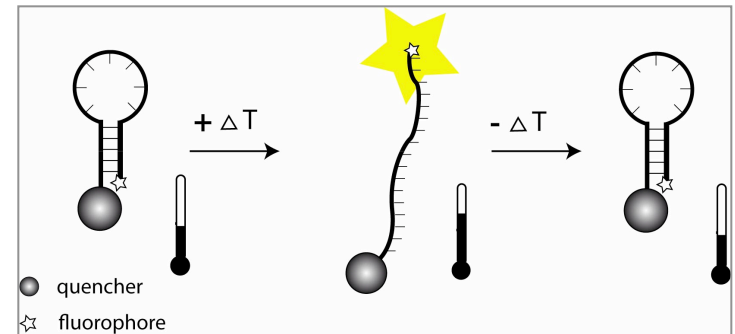


$t_p=10\mu\text{s}$ ,  $\lambda=532\text{ nm}$ ,  $F=100\text{J}/\text{cm}^2$ ,  $T_e = T_1$   
 $\rightarrow T=T_0+160\text{K}$ , 200 nm surrounding



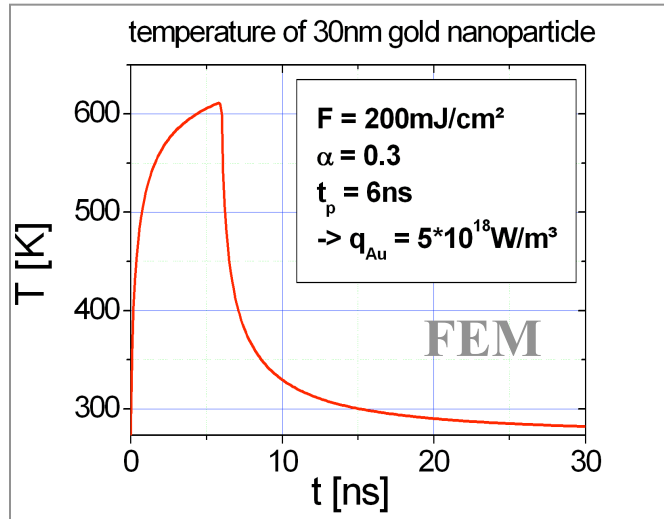
cw,  $\lambda=532\text{ nm}$ , 25mW,  $d=0.2\text{mm}$ ,  $F=830\,000\text{ J}/\text{cm}^2$

$\rightarrow T=T_0+4\text{K}$ , the complete medium in the laser focus

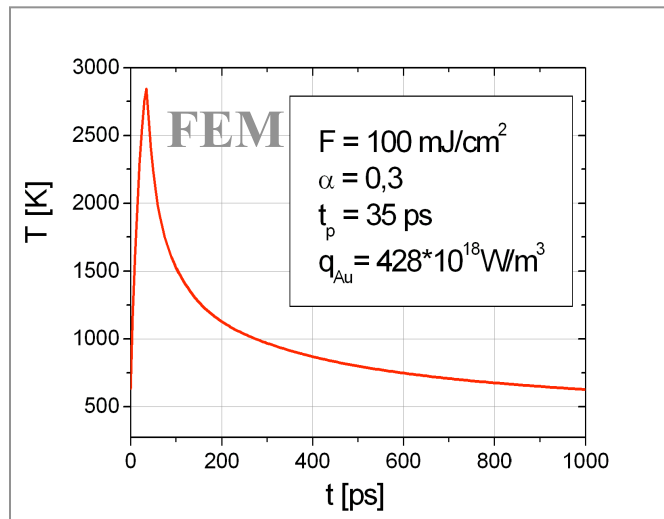


Prinzip der optisch schaltbaren Fluoreszenzsonden

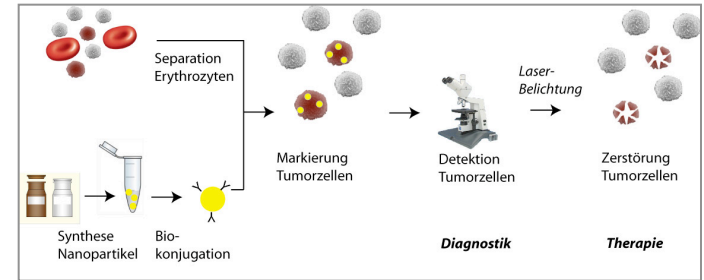
- $\rightarrow$  NP  $10\mu\text{s}$  erwärmt  $\rightarrow$  Dissoziation DNA Doppelstrang
- $\rightarrow$  gequenches Fluorophor kann fluoreszieren



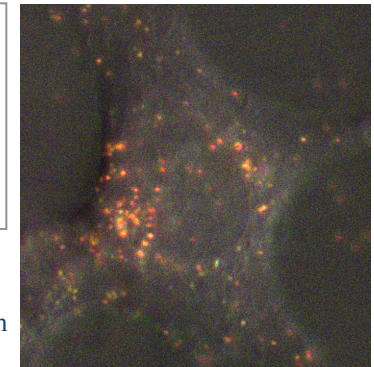
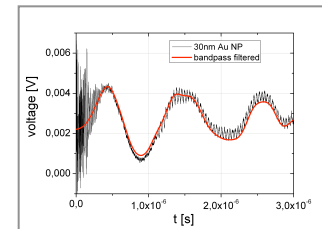
$t_p=6\text{ns}$ ,  $\lambda=532\text{ nm}$ ,  $F=200\text{mJ/cm}^2$   
 $\rightarrow T=T_0 + 330\text{K}$ , 70nm surrounding, after 100 ps heat cond.



$t_p=35\text{ps}$ ,  $\lambda=532\text{ nm}$ , dec.  $F=100\text{mJ/cm}^2$   
 $\rightarrow T=T_0+2700\text{K}$ , (to high  $\rightarrow c_e(T)!!$ ) 70 nm surrounding,  
 $\rightarrow$  during pulse no heat conductance (10nm)

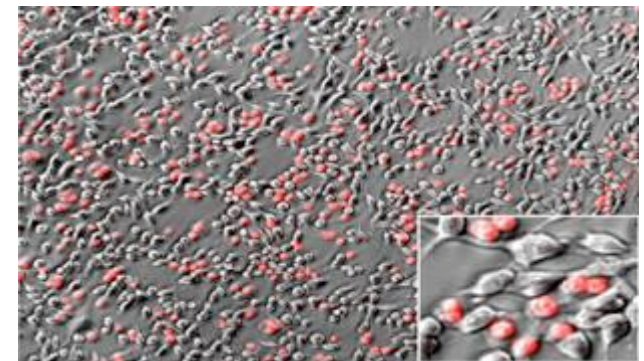


NP-basierte Therapie von Tumorzellen im Blut

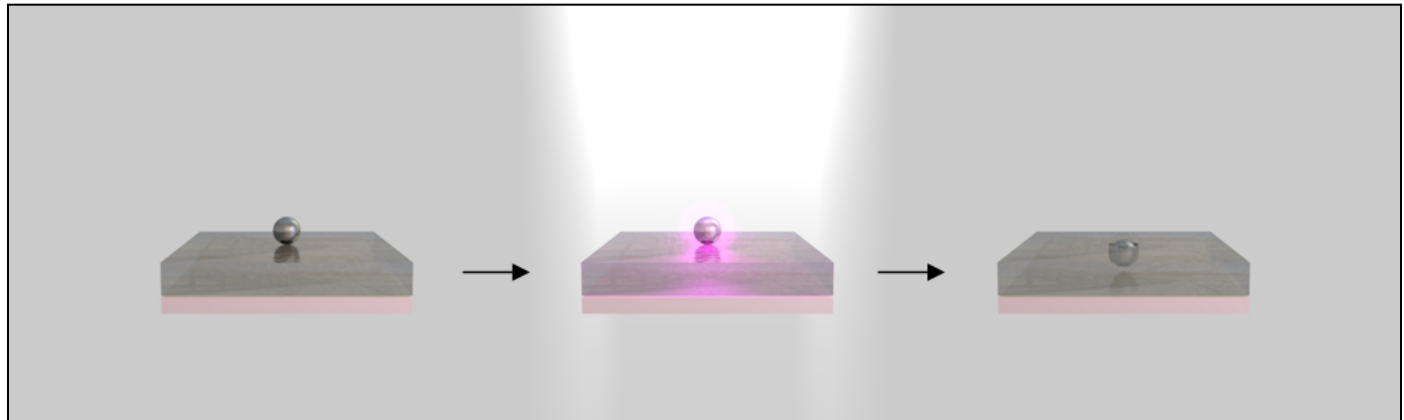


Drucksignal (PVDF Folie + Verstärker) von Au NP in H<sub>2</sub>O nach 6ns Laserpuls bei 532 nm

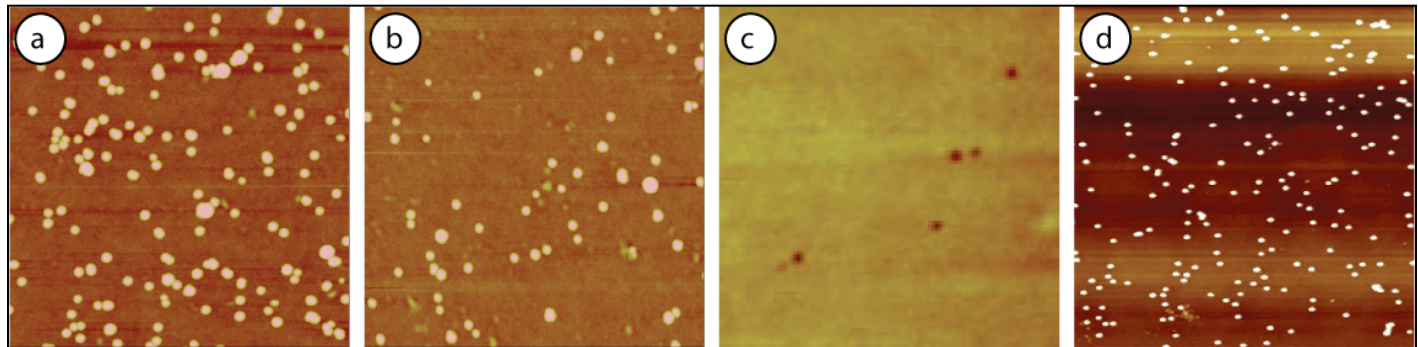
80nm Au NP auf adherenten Zellen



selektive Zerstörung von adherenten Zellen in einer Mischkultur mittels laserbestrahlter Au NP



several laser parameters on 30 nm Au NP



a  
 $F = 2,8 \text{ mJ/cm}^2$   
 $t = 100 \mu\text{s}$   
 cw  
 no sinking

b  
 $F = 2,8 \text{ mJ/cm}^2$   
 $t = 3 \text{ ps}$   
 $f = 76 \text{ MHz}$   
 no sinking

c  
 $F = 2,8 \text{ mJ/cm}^2$   
 $t = 100 \text{ fs}$   
 $f = 76 \text{ MHz}$   
**sinking** (dark regions)

d  
 $F = 2,8 \text{ mJ/cm}^2$   
 $t = 100 \text{ fs}$   
 $f = 76 \text{ MHz}$   
**no ablation** (section analysis)

**c) warming of Au NP - surrounding near to  $T_g$  so viscosity changes, particles sinking ???**

## relaxation dynamics of the electron temperature

### electron-phonon coupling model

	thermal diffusion	electron-phonon coupling	optical excitation
$c_e(T_e) \frac{\partial T_e}{\partial t} = \kappa \Delta T_e - G(T_e - T_l) + P(t)$			
$c_l \frac{\partial T_l}{\partial t} = G(T_e - T_l)$			

<u>electronic heat capacity</u>	$c_e = c_0 T_e, c_0 = 66 \text{ J/m}^3 \text{K}^2$	[1]
<u>lattice heat capacity</u>	$c_l = 3 \times 10^6 \text{ J/m}^3 \text{K}$	[2]
<u>electron-phonon coupling constant</u>	$G = 3 \times 10^4 \text{ J/psm}^3 \text{K}$	[2]
<u>transient absorption</u>	$A = 27\%$	[3]
<u>thermal conductivity losses to matrix</u>	$\kappa \Delta T_e = 0$	
<u>excitation energy density per unit time and unit volume</u>	$P(t)$	

-> absorption of photons by electrons (100 fs), electron-phonon relaxation (10 ps) and phonon-phonon relaxation (100 ps) are separated in time

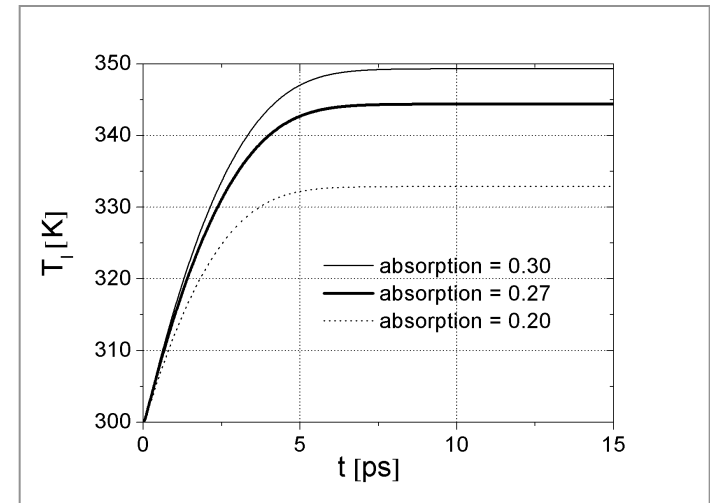
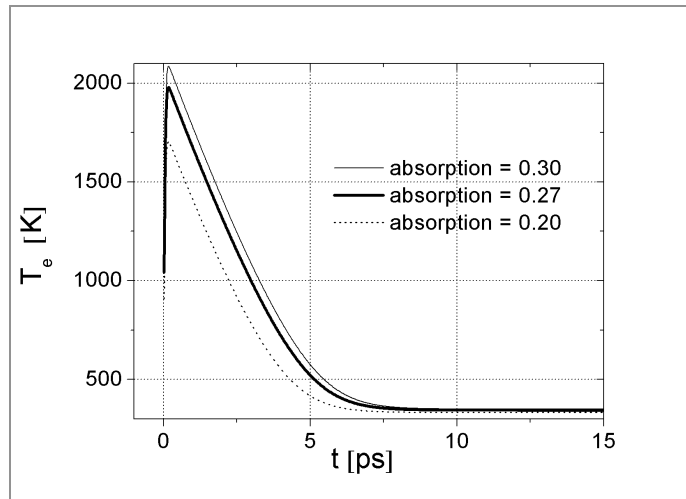
-> during fs-pulse: no time for plasma development, heating diffusion into material is negligible

[1] K.L. Kelly et al. J.Phys.Chem. B **107** p668 (2003)

[2] R. Groeneveld et al. PRB **51(17)** p11433 (1995)

[3] S.L.Logunov et al. J.Phys.Chem. B **101** p3713 (1997)

## results: electron-phonon coupling model



- electron temperature very high (2000 K, 5 ps relaxation time)
- lattice temperature increase to 344,5 K in 5 ps

- + heat of electrons in nanoparticle is fast and effective (cold ions can absorb more photons)
- reduced ablation threshold (3 mJ/cm<sup>2</sup>)

heat equation

$$\rho C_p \frac{\partial T}{\partial t} + \nabla \cdot (-k \nabla T) = q$$

domains: gold np (truncated sphere with r=15nm without a cup of 5nm)  
 PMMA matrix as cuboid of 200 by 200 by 100 nm

boundary between the two subdomains

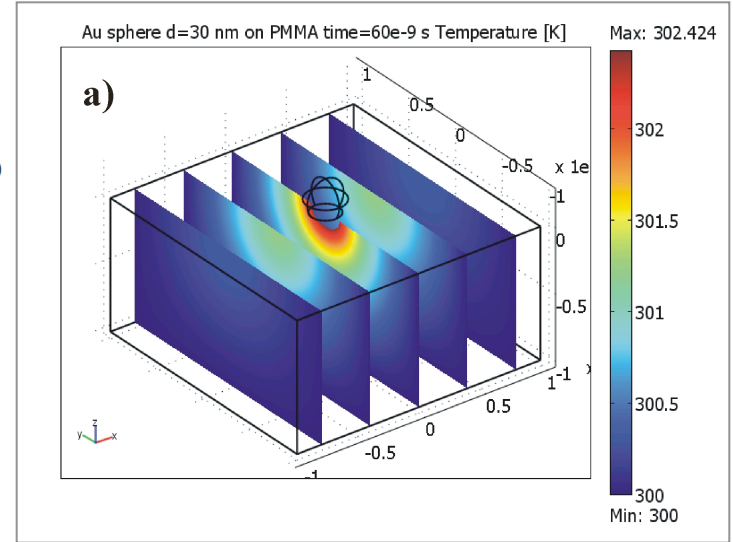
$$-\mathbf{n}_{Au} \cdot (-k_{Au} \nabla T_{Au}) - \mathbf{n}_{PMMA} \cdot (-k_{PMMA} \nabla T_{PMMA}) = 0$$

top boundaries: external natural cooling conditions

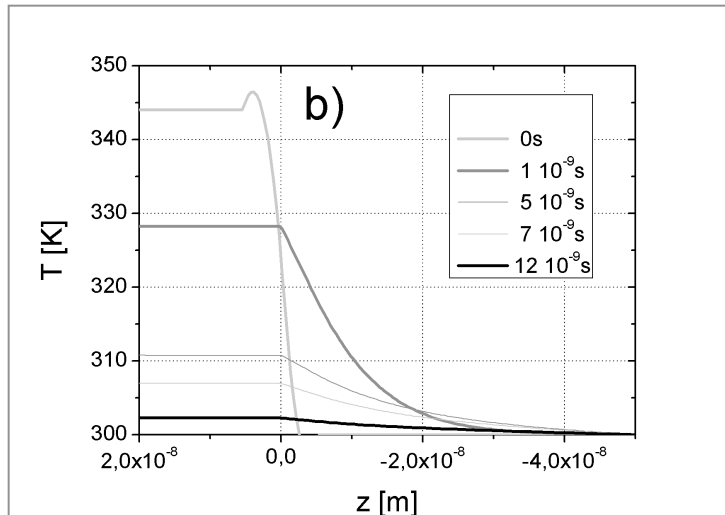
$$-\mathbf{n} \cdot (-k \nabla T) = h(T_{inf} - T)$$

$$h = 0.54 F_{lam} \left( \frac{\Delta T}{L} \right)^{0.25} \quad F_{lam} = k \left( \frac{Ra}{L^3 \Delta T} \right)^{0.25}$$

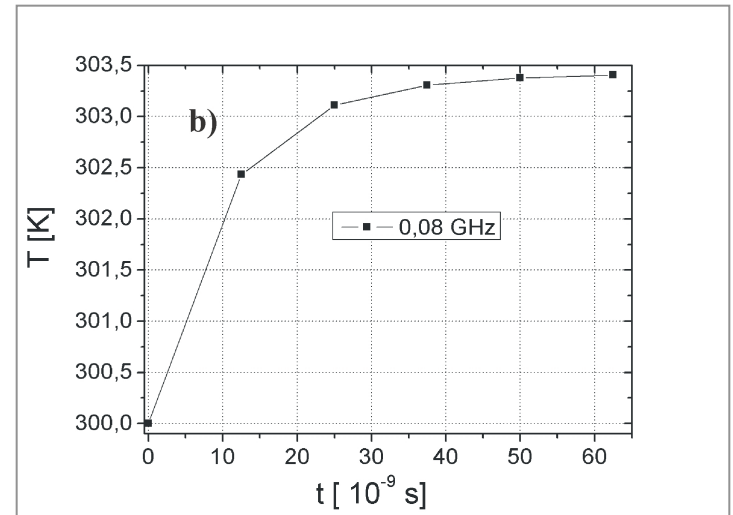
h = horizontal surface Ra = Rayleigh number L = characteristic length



conductive heat transfer from gold np -> PMMA starts after equilibrium of electron gas with lattice



temperature increase of PMMA is very local (40 nm), caused by low heat capacity of nanoparticle

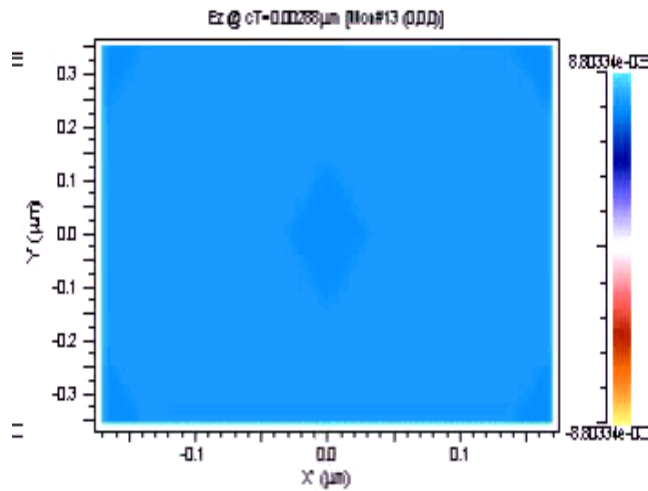
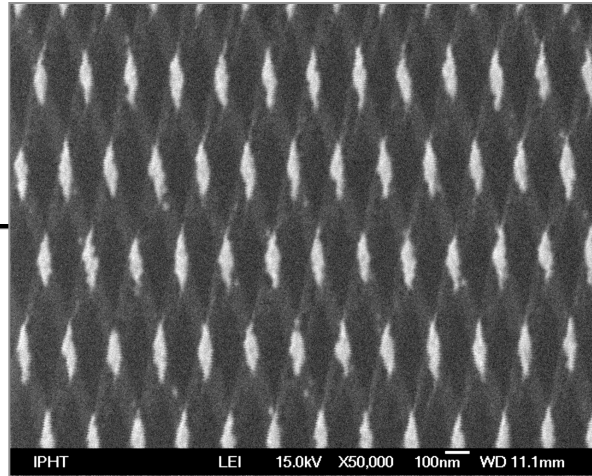


at the end of the period between two fs-laser pulses (12 ns) is a small temperature increase of PMMA and nanoparticle



### 3. TEIL: Wirkungen hoher Feldstärken in dielektrischer Umgebung von NP

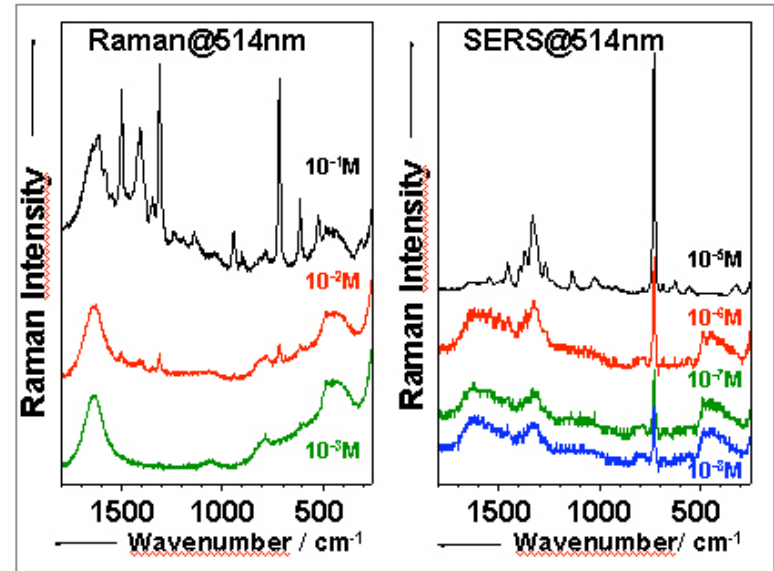
(eine Folge der Anregung von Metall NP mit fs Laserpulsen)

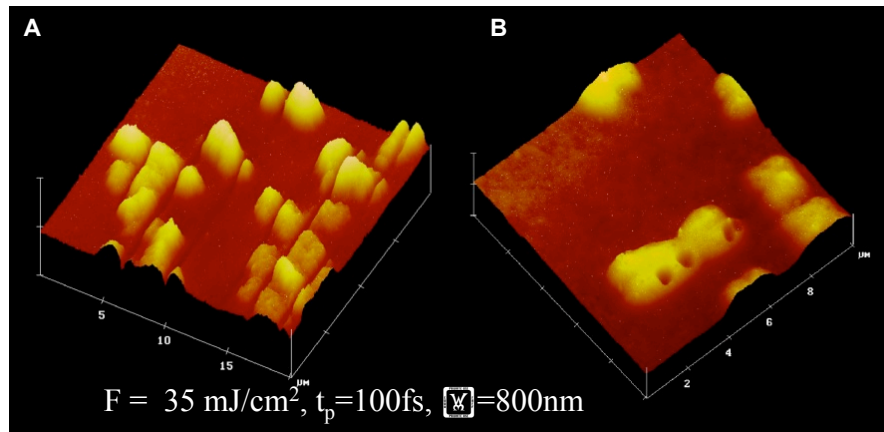


**FDTD:**

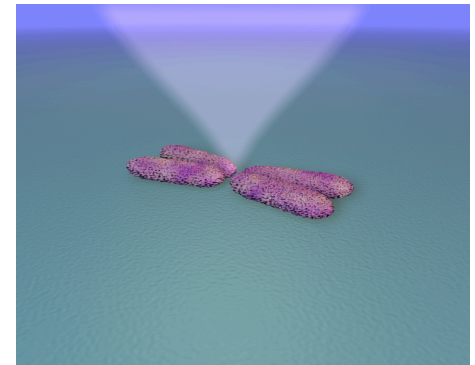
Routenstrukturen ( $r=5\text{nm}$ )

$$f = \frac{\vec{E}}{\vec{E}_0} = 13$$

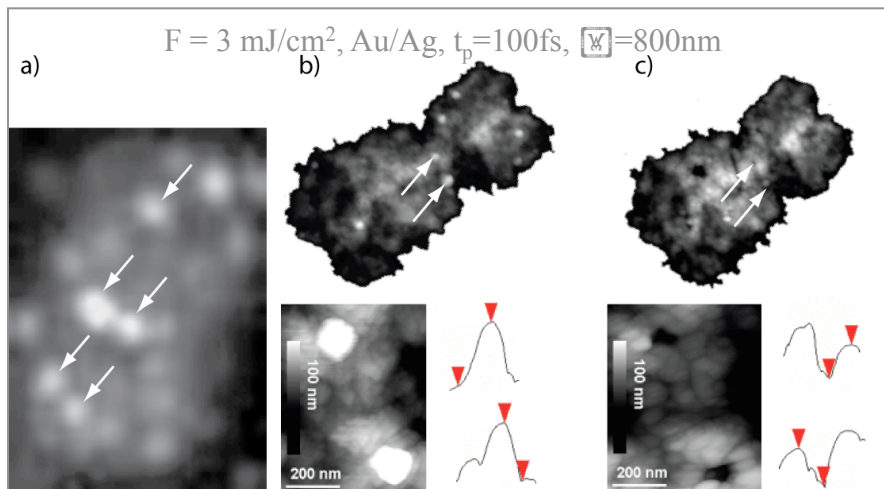




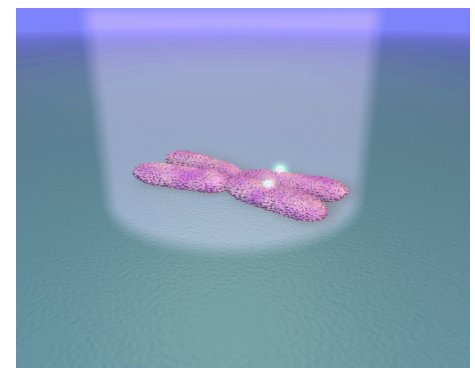
destruction of chromosomes without NP in an unspecific way



multiphoton excitation of electrons in molecules takes place, minimal destruction 100nm



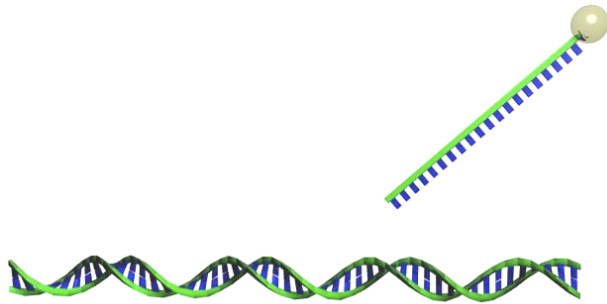
destruction of chromosomes nearby the NP in a specific way



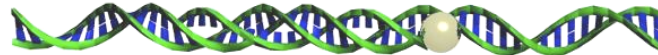
reduction of the laser pulse fluence below the unspecific destruction limit of chromosomes  
using chromosomes with specific bound NP

# ZIELSTELLUNG

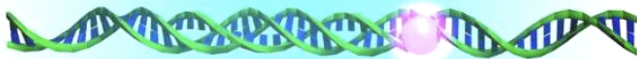
- 1 Konjugation von Nanopartikel an sequenzspezifische Oligonukleotide



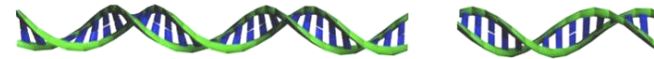
- 2 Hybridisierung der modifizierten Oligonukleotide an die Zielsequenz



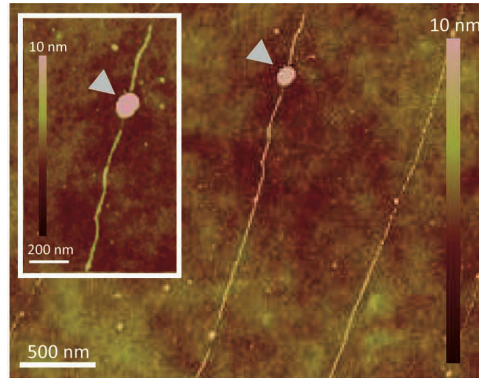
- 3 Sequenzspezifisches Schneiden der DNA an den markierten Stellen durch Laserstrahlung



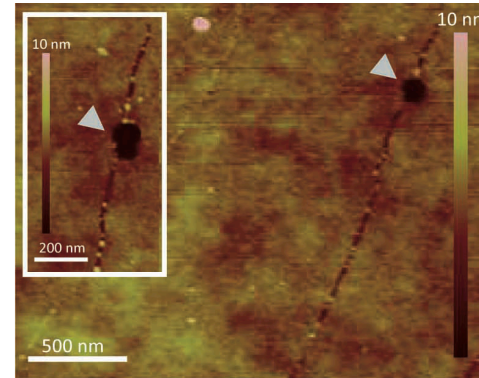
- 4 Ergebnis: Geschnittenes DNA-Molekül



Ist es möglich DNA-Moleküle spezifisch an den markierten Stellen zu schneiden?

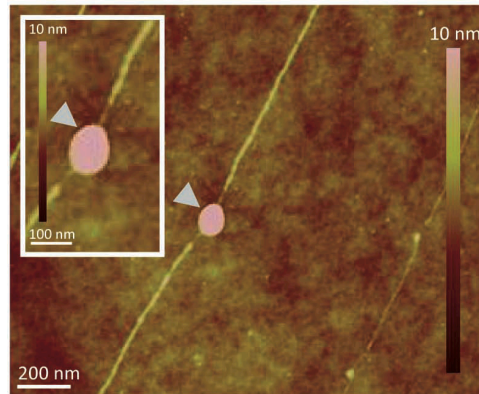


a

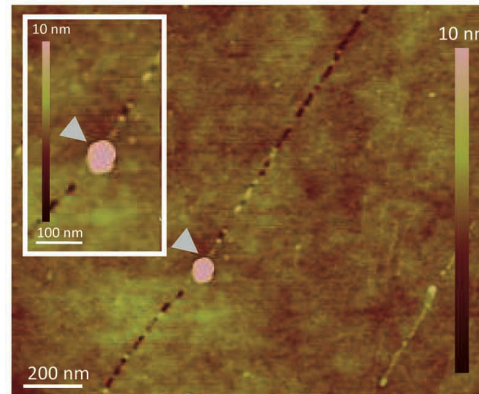


b

Zerstörung PMMA um NP  
 PMMA Zerstörung entlang DNA  
 DNA Zerstörung  $d = 90 \text{ nm}$

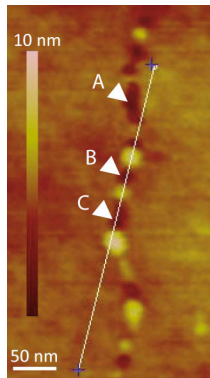


c

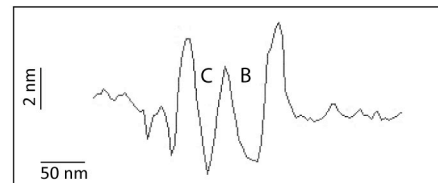


d

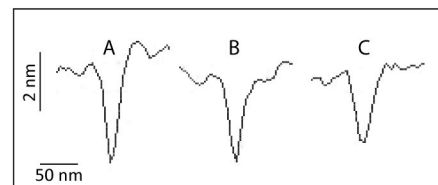
PMMA Zerstörung entlang DNA  
 DNA Zerstörung  $d = 90 \text{ nm}$



a



b

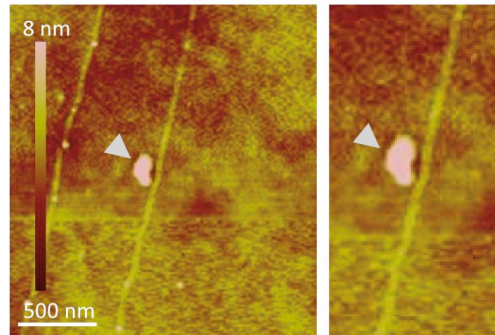


c

**DNA Bündel**  
**Tiefe 3 nm**  
**Höhe 3 nm**  
**Breite 15 nm**

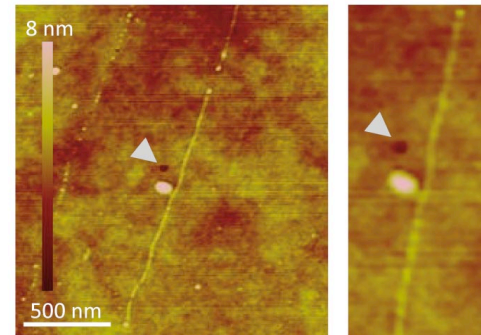
**Ergebnisse reproduzierbar !**

before laser irradiation



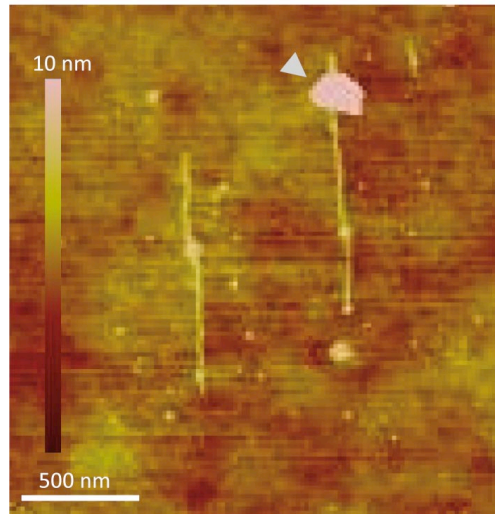
a

after laser irradiation

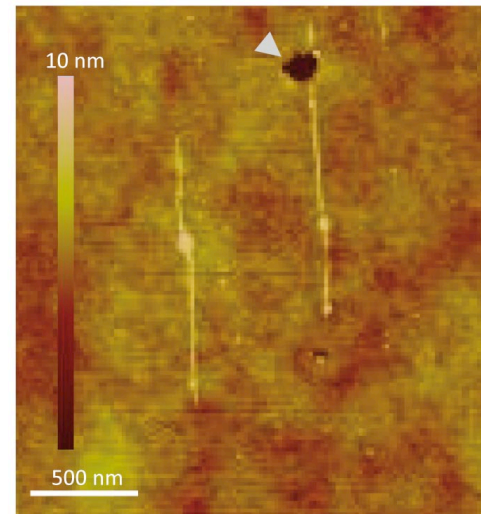


b

**D (NP-DNA) > 100 nm**



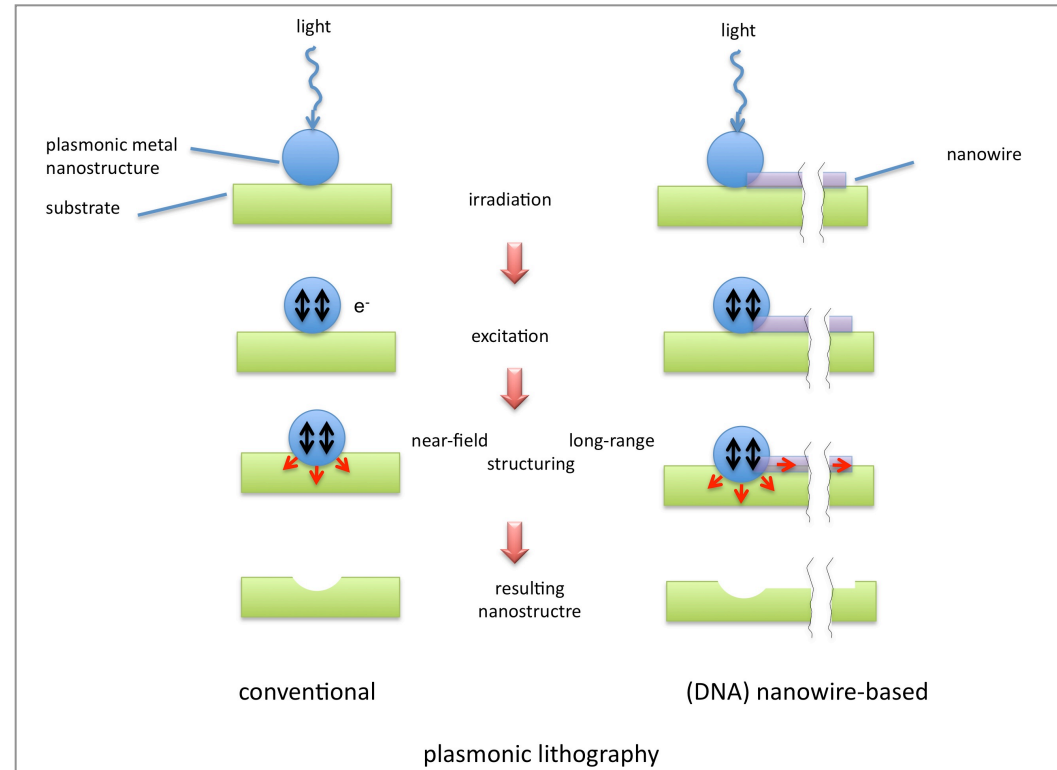
c



d

**D (NP-DNA) < 30 nm**

- > keine PMMA Zerstörung entlang der DNA
- > Kopplung an DNA sehr Abstands sensitiv



## Hypothese:

- $e^-$  des NP überwinden Austrittsarbeit oder hohen Felder an NP bewirken Zerstörung umgebender PMMA Moleküle und Übertragung  $e^-$  Impuls auf DNA Basen Paare, wo Impuls weitergeleitet wird und umgebende PMMA Moleküle ebenfalls zerstört werden
- nichtlinear angeregte  $e^-$  im NP kann man als Licht auf der Nanoskala sehen, dieses koppelt in die DNA, wo es dann propagiert. Das UV Licht zerstört PMMA Molekülbindungen

-> **neues Projekt, wo unter anderem dieses Phänomen auf der molekularen Ebene beschrieben werden soll**

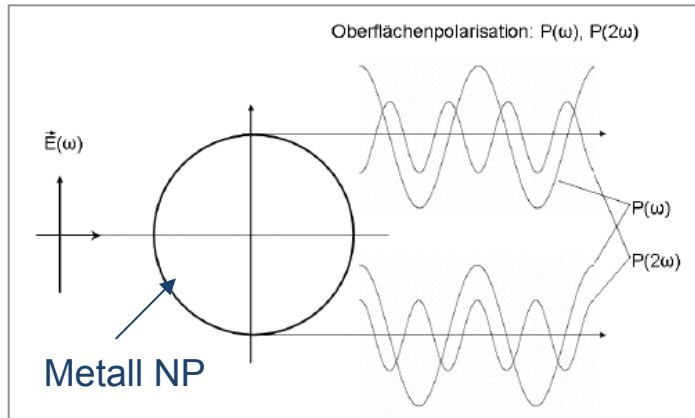
#### 4. TEIL: Plasmon Polariton Resonanz unterstützte Second Harmonic Generation (SHG)

(wenn die Laserfelder in den Bereich der Feldstärken für die Atombindungen kommen)

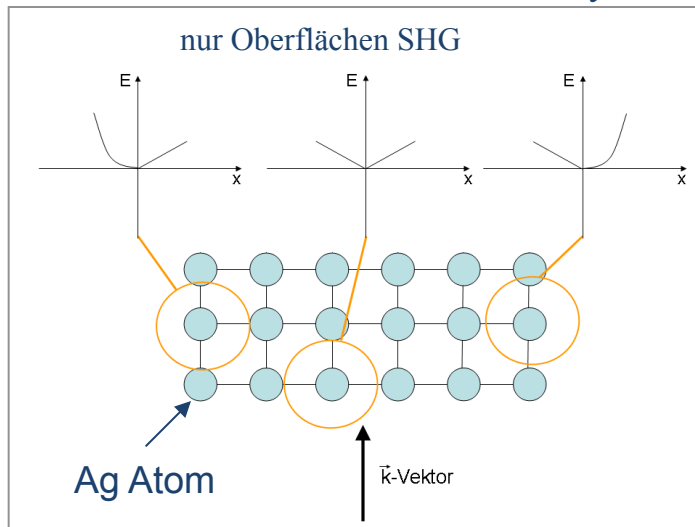


$$\vec{P} = \epsilon_0 (\chi^{(1)} \vec{E} + \chi^{(2)} \vec{E}^2 + \dots)$$

aber:  $c_{36}^{(2)KDP} = 1.1 \cdot 10^{-23} [As/V^2]$ ,  $\rightarrow E$  muss groß

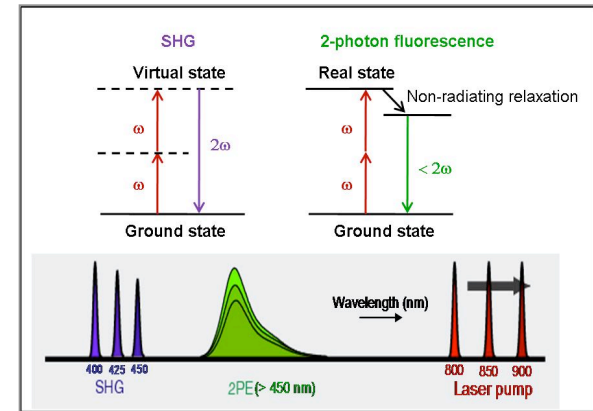


da der NP Durchmesser  $\ll \lambda \rightarrow$  asym. NP

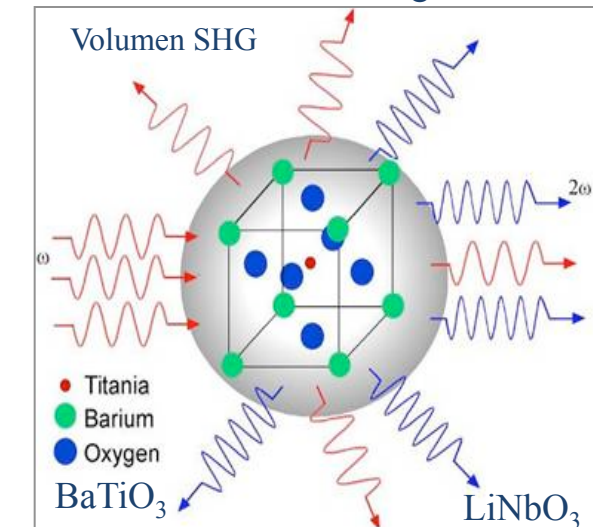


$$I(2\omega) \propto |f(2\omega)|^2 |f(\omega)|^4 I(\omega)^2 \rightarrow f = \frac{\vec{E}}{E_0}$$

aber: Plasmonenresonanzverstärkung der SHG

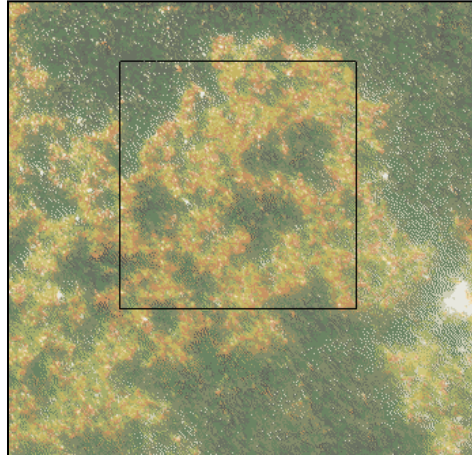


SHG und Fluoreszenz möglich

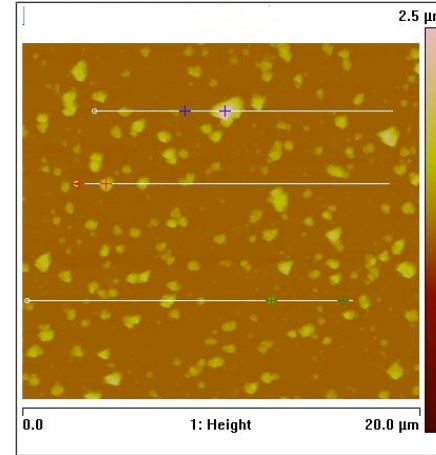


$$I(2\omega) \propto I(\omega)^2$$

aber: Volumen SHG



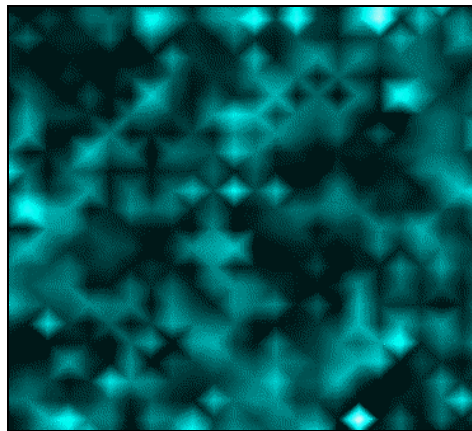
**Dunkelfeldmikroskopie:**



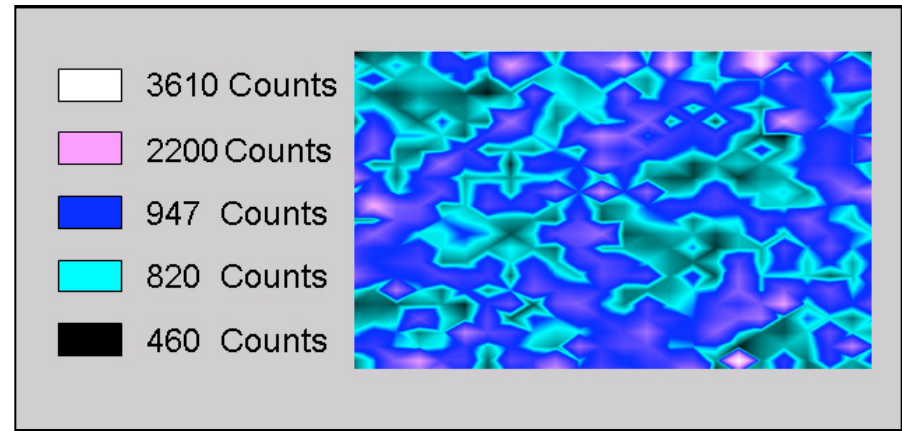
Parameter:

- $\lambda=800\text{nm}$ ,  $f=76\text{MHz}$ ,
- 40iger Objektiv  $\rightarrow d = 3 \mu\text{m}$
- $F=5,6\text{mJ/cm}^2$
- 
- 26 x 26 Pixel Tischscanning
- Interferenzfilter 400 nm
- Kurzpassfilter 800nm OD8
- 500  $\mu\text{m}$  Blende
- (konvokales Mikroskop)
- APD, SEV + Counter

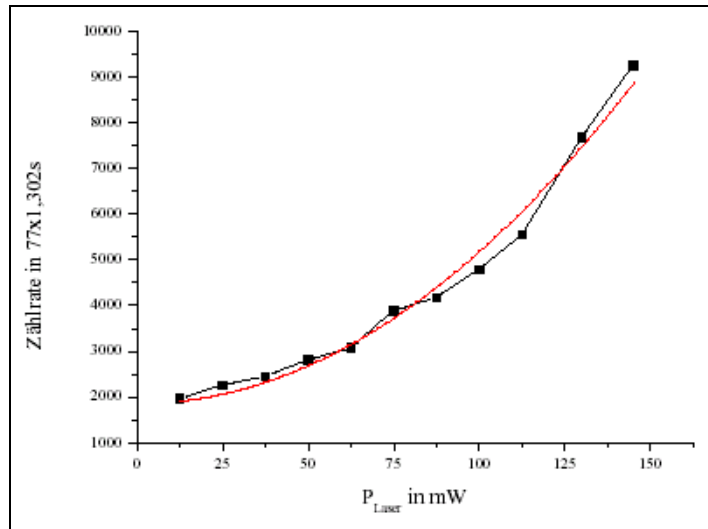
**AFM:** Nanocluster aus 25nm Ag  
Clustergröße: 100 nm – 500 nm



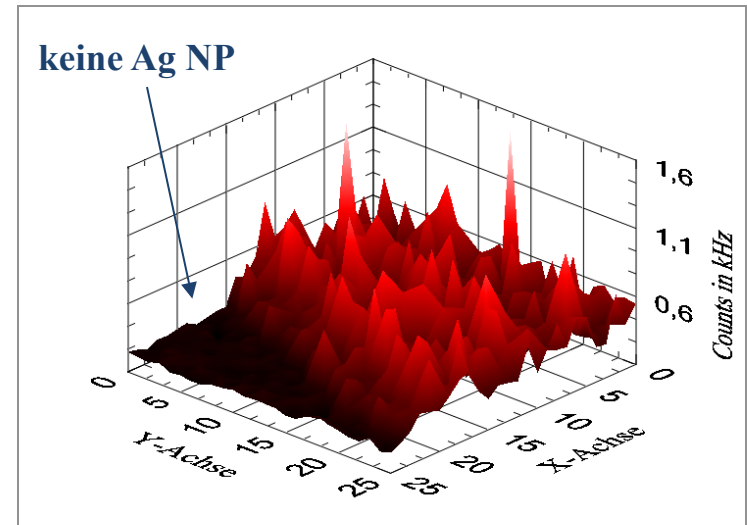
**SHG Signal** (Schwellwert 947 counts)



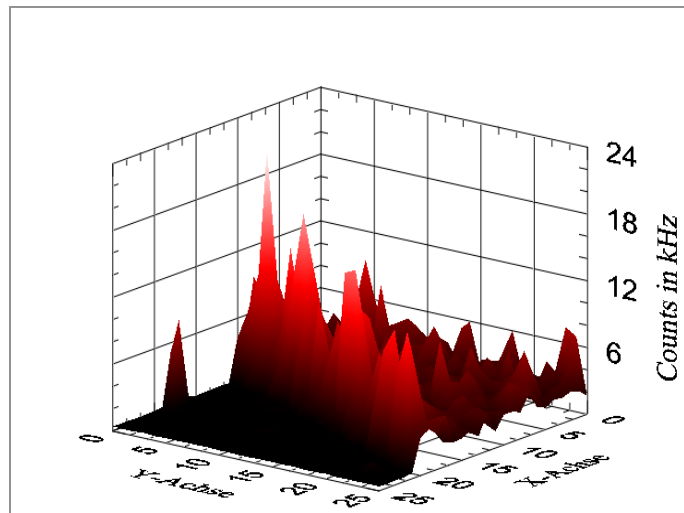
**Mehrfalschfarbendarstellung SHG Signal**



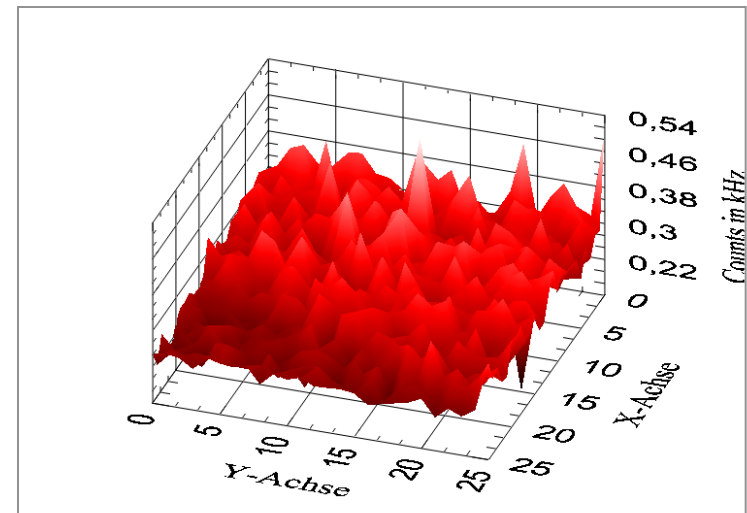
quadratische Abhängigkeit der Zählrate von Laserleistung



SHG nach der Herstellung

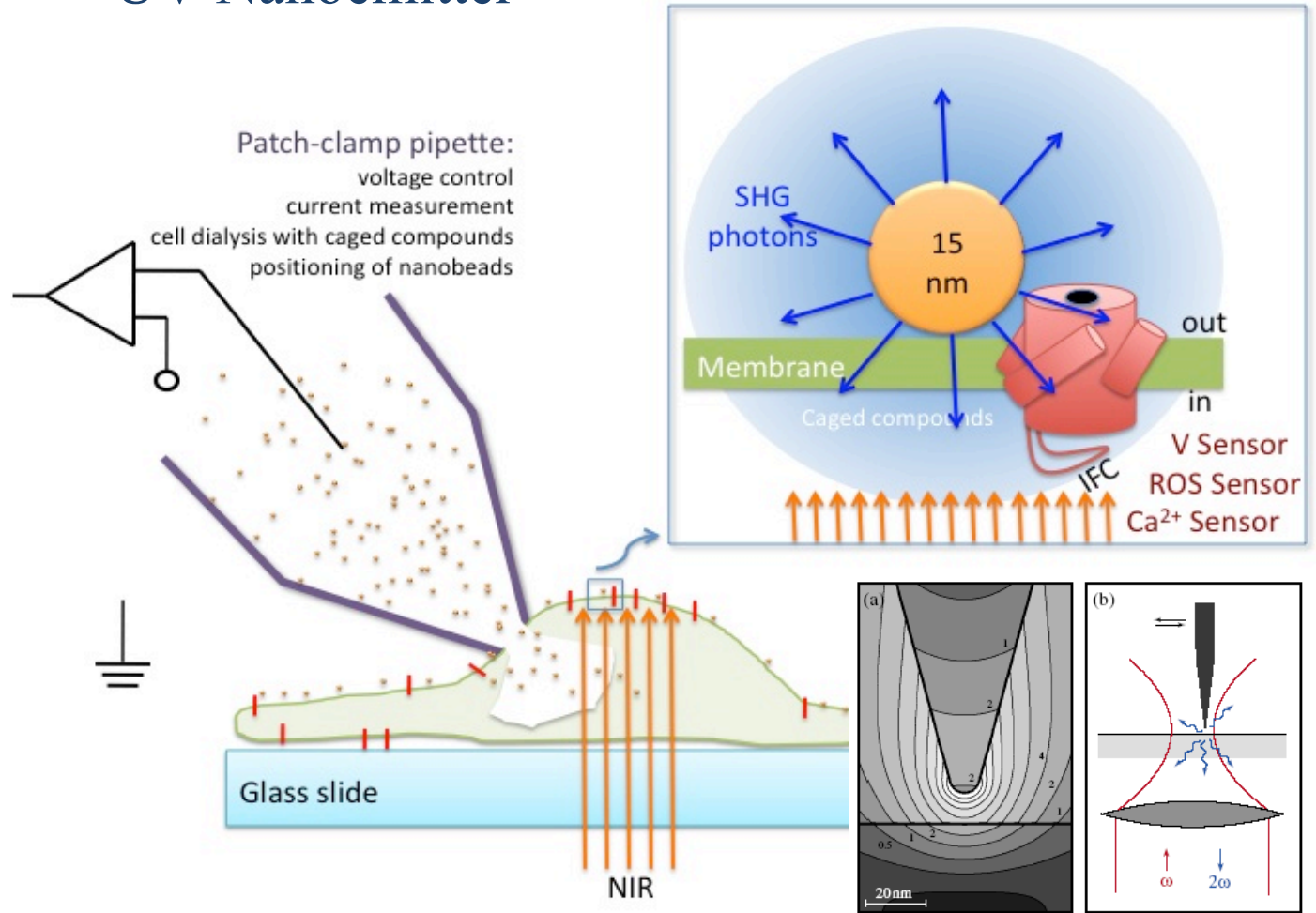


Fluoreszenz nach 24 h



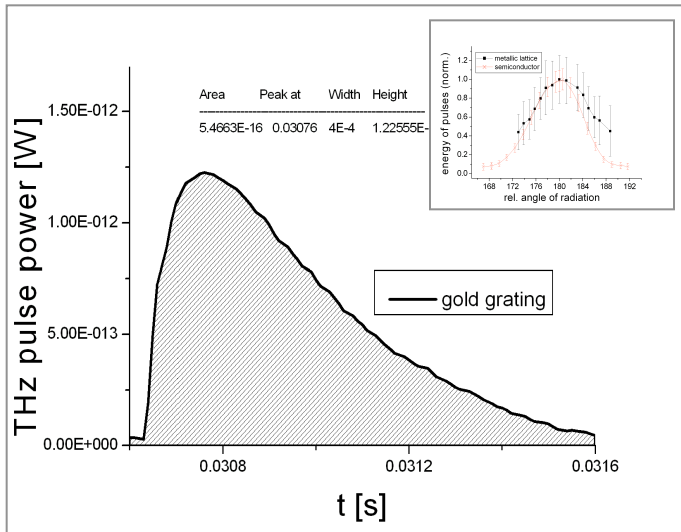
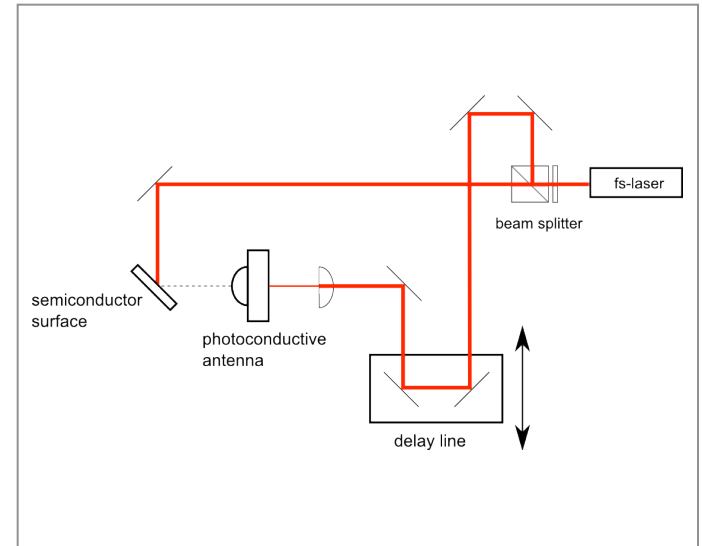
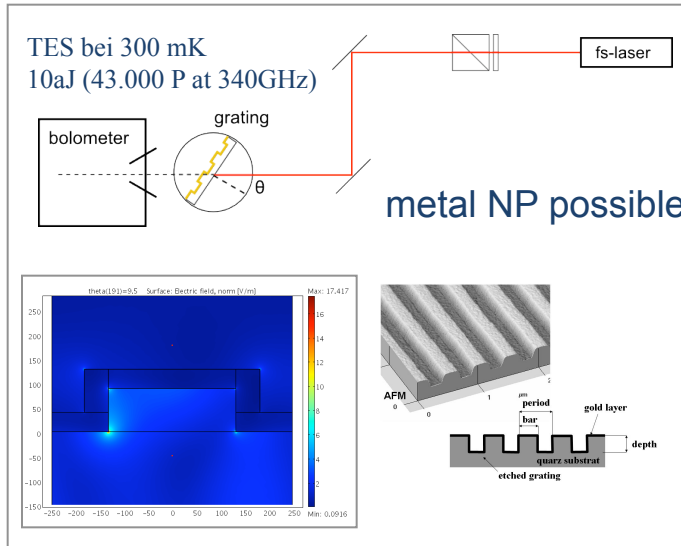
SHG nach 7 Tagen

# UV Nanoemitter



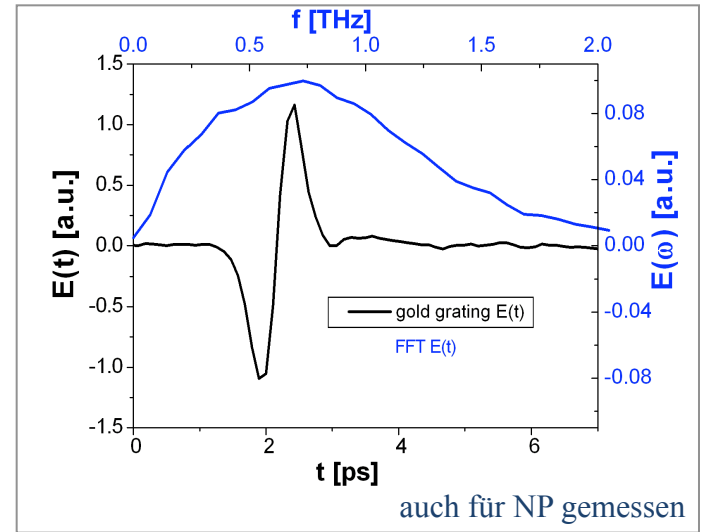
## 5. TEIL: Terahertz Abstrahlung von metallischen Nanostrukturen

(eine Folge der Beschleunigung von freien Elektronen im evaneszenten Feld, oder?)



Integration over angle:

$E_p = 1.7 * 10^{-15} \text{ J}$



Integration over spectrum:

$E_p = 170 * 10^{-15} \text{ J}$

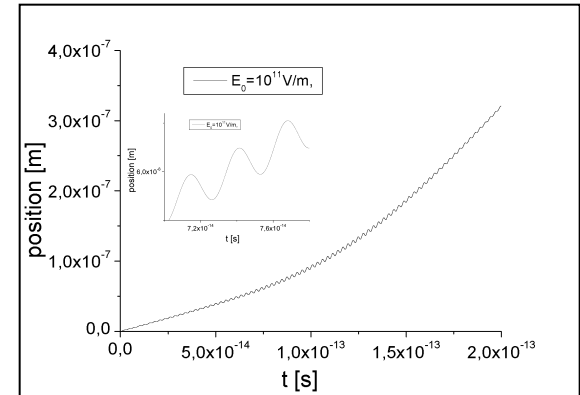
$P_p = 56 \text{ mW}$

## Modell

- freie e<sup>-</sup> über Au Oberfläche werden durch Laserfeld beschleunigt
- e<sup>-</sup> in evaneszenten Feld sieht in verschiedenen Richtungen unterschiedliche starke Felder (ponderomotive Kraft)

$$m \ddot{z} = eE_0 \cos(\omega t + \phi_0) * e^{-z(t)/\lambda_{ev}} * e^{-2 * \ln 2 \left( \frac{\log \frac{E_0}{E_S}}{\log 2} \right)^2}$$

oscillation of Laser field    evanescent field    Gauss-envelope

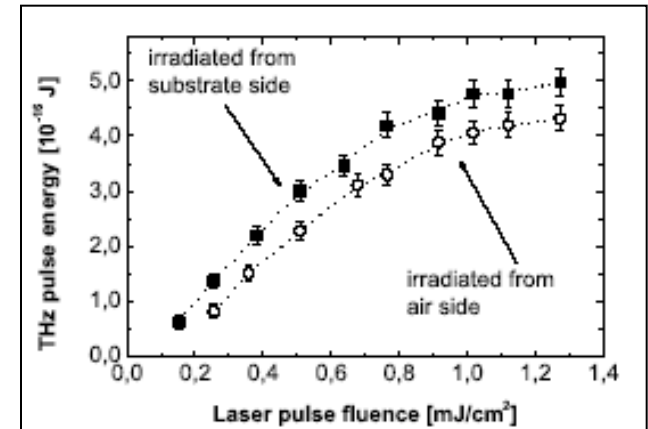
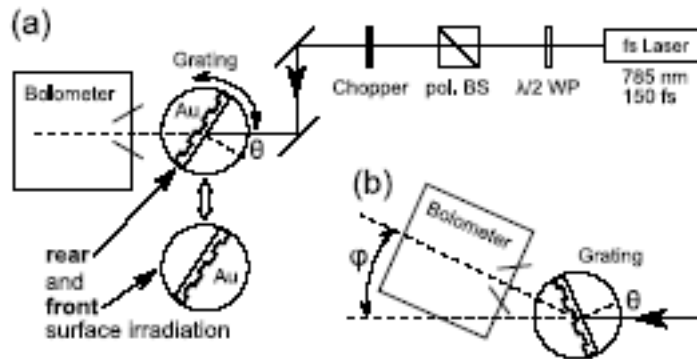


## Lösung

- 1d, klassisch, numerisch, - durch Berechnung von Position und Energie Schritt für Schritt mit Δt=1as

beschleunigte Bewegung einer Ladung strahlt ab  
 aber: **nur 250 fs**  
 el-ph Relaxation (2 ps) bestimmt THz Pulslänge  
 aber: erst ab **E = 10<sup>10</sup> V/m** e<sup>-</sup> Austritt

## Experiment



Das bisherige Modell kann diese Experimente nicht erklären -> **neues Modell notwendig**

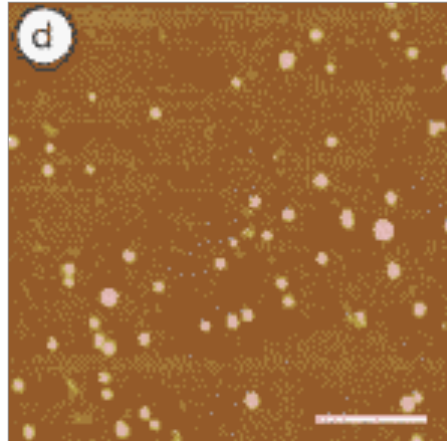
## Abspann: Coulomb Explosion – Ablation der NP

(eine Folge von zu hohen Feldern und damit Austritt von Elektronen aus der Oberfläche (Plasma))

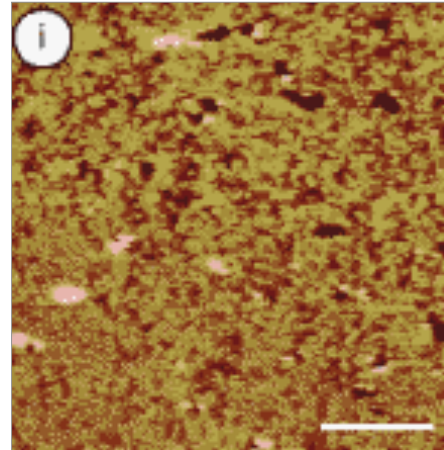
$F = 3 \cdot 10^{-7} \text{ mJ/cm}^2 = 1 \text{ GO}$  niedriger als bei dielektrischen Materialien



AFM picture bevor laser radiation:



300 fs, 90mJ/cm<sup>2</sup>



d) 30 nm Au NP on top of a 50 nm e<sup>-</sup> sensitive PMMA layer

i) cold fragments do not sink into PMMA,  
multiphoton excited e<sup>-</sup> override the work function  
NP has positive ions, which reject each other

## SUMMARY:

### lange Laserpulse auf Metall NP

- schnelle und nanoskalige Wärmequellen

(Krebstherapie/Nano Sensoren)

### hohe Laserfelder auf Metall NP

- Transport von Licht auf Nanometerskalen

(Plasmonik/Optoelektronik)

- Nano UV-Lichtquellen

(Nano Biochemie)

- Nano THz Quellen

(???)

**Hier wird professionelle theoretische Beschreibung notwendig**

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